Professor Bareyre for his generous cooperation through private correspondence in supplying us with information on the errors.

This detailed comparison of our results and those of Refs. 2 and 3 shows that indeed there is excellent agreement between us as far as the H_{39} parameters are concerned. For the H_{311} parameters, the δ's also agree very well, but there is a discrepancy between the η's. It should be kept in mind, however, that the method of Refs. 2 and 3 are an energy-dependent, overall fit to data in a large range, of which the compared energy is the upper limit, and hence some disagreement with our single-energy method is not altogether surprising.

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[1]C. Lovelace, in Proceedings of the International Conference on Elementary Particles, Heidelberg, Germany, 1967, edited by H. Filthuth (North-Holland, Amsterdam, 1968), p. 79.

[2]R. Ayed, P. Bareyre, and G. Villet, Phys. Lett. 31B, 598 (1970).

[3]R. Ayed, P. Bareyre, and G. Villet, private communication.

OPTICALLY MODULATED X-RAY DIFFRACTION. Isaac Freund and B. F. Levine [Phys. Rev. Lett. 25, 1241 (1970)].

(1) In Eq. (13) replace $ie/2m\omega_x^2$ by $ie/2mc\omega_x$ and $[\hat{u}_s\cdot\vec{Q}(hkl)]$ by $(c/\omega_x)[\hat{u}_s\cdot\vec{Q}(hkl)]$; this makes $\vec{\theta}_{psi}(hkl)$ dimensionless in accord with previous usage.

(2) In line four following Eq. (17) replace $V-\mathfrak{V}(0)$ by $V=\mathfrak{V}(0)$.

(3) In the discussion preceding Eq. (22), where $\mathfrak{D}_x^{\omega_i}(\vec{r})$ is defined, note that $\omega_i=\omega_1$, ω_2, or ω_3 is an optical frequency and that E_x refers to the x component of the optical field $\vec{E}(\omega_i)$.

COLLISIONAL EFFECTS ON INDUCED EMISSION AND ABSORPTION TRANSITION PROBABILITIES IN ATOMIC SYSTEMS. Chung-Nan Chang and Sotiris Koutsoyannis [Phys. Rev. Lett. 25, 1399 (1970)].

Equation (12) of the text should read

$$N\left\langle\frac{\exp(\pm i\vec{k}_l\cdot\vec{R})}{R^3}\right\rangle_{\rm ens}=\frac{N}{V}\int_\tau\frac{\exp(\pm i\vec{k}_l\cdot\vec{R})}{R^3}d\tau=nC,$$

where C is the real quantity

$$C=-4\pi\int_{R^*}^\infty\frac{\sin k_lx}{k_lx^2}dx.$$

Upon taking the square of the absolute value of the matrix elements it is found that the transition probabilities are indeed functions of the density n but both are modified by the same factor $(1+\alpha n)^2$ making their ratio independent of the density to this approximation.

COUPLED-CHANNEL BORN-APPROXIMATION CALCULATION OF TWO-NUCLEON TRANSFER REACTIONS IN DEFORMED NUCLEI. T. Tamura, D. R. Bes, R. A. Broglia, and S. Landowne [Phys. Rev. Lett. 25, 1507 (1970)].

The factor i^l in Eq. (4) was not included in the computation. Therefore, the sign of F_2 and, consequently, the theoretical curves in Figs. 1(a)-1(c) were erroneous. Corrected results that replace those in the old Fig. 1(a) are presented in the new figure given here. The calculations were made with 0^+-2^+-4^+ coupling in both incident and exit channels, everything else being the same as in previous calculations. Note the improved agreement with experiment, in particular concerning the 2^+ cross section. Qualitative statements concerning Figs. 1(b) and 1(c), in particular those concerning the importance of the multistep processes, remain correct. We are very much indebted to Dr. T. Udagawa and Dr. B. Sorensen for their kind cooperation in finding and confirming the above error.