

Professor Bareyre for his generous cooperation through private correspondence in supplying us with information on the errors.

This detailed comparison of our results and those of Refs. 2 and 3 shows that indeed there is excellent agreement between us as far as the $H_{3,9}$ parameters are concerned. For the $H_{3,11}$ parameters, the δ 's also agree very well, but there is a discrepancy between the η 's. It should be kept in mind, however, that the method of Refs. 2 and 3 are an energy-dependent, overall fit to data in a large range, of which the compared energy is the upper limit, and hence some disagreement with our single-energy method is not altogether surprising.

¹C. Lovelace, in Proceedings of the International Conference on Elementary Particles, Heidelberg, Germany, 1967, edited by H. Filthuth (North-Holland, Amsterdam, 1968), p. 79.

²R. Ayed, P. Bareyre, and G. Villet, Phys. Lett. 31B, 598 (1970).

³R. Ayed, P. Bareyre, and G. Villet, private communication.

OPTICALLY MODULATED X-RAY DIFFRACTION. Isaac Freund and B. F. Levine [Phys. Rev. Lett. 25, 1241 (1970)].

(1) In Eq. (13) replace $ie/2m\omega_x^2$ by $ie/2m\omega_x$ and $[\hat{a}_s \cdot \vec{Q}(hkl)]$ by $(c/\omega_x)[\hat{a}_s \cdot \vec{Q}(hkl)]$; this makes $\theta_{psl}(hkl)$ dimensionless in accord with previous usage.

(2) In line four following Eq. (17) replace $V - \nabla(0)$ by $V = \nabla(0)$.

(3) In the discussion preceding Eq. (22), where $\mathfrak{D}_x^{\omega_i}(\vec{r})$ is defined, note that $\omega_i = \omega_1, \omega_2,$ or ω_3 is an optical frequency and that E_x refers to the x component of the optical field $\vec{E}(\omega_i)$.

COLLISIONAL EFFECTS ON INDUCED EMISSION AND ABSORPTION TRANSITION PROBABILITIES IN ATOMIC SYSTEMS. Chung-Nan Chang and Sotiris Koutsoyannis [Phys. Rev. Lett. 25, 1399 (1970)].

Equation (12) of the text should read

$$N \left\langle \frac{\exp(\pm i\vec{k}_l \cdot \vec{R})}{R^3} \right\rangle_{\text{ens}} = \frac{N}{V} \int_{\tau} \frac{\exp(\pm i\vec{k}_l \cdot \vec{R})}{R^3} d\tau = nC,$$

where C is the real quantity

$$C = -4\pi \int_{R^*}^{\infty} \frac{\sin k_l x}{k_l x^2} dx.$$

Upon taking the square of the absolute value of the matrix elements it is found that the transition probabilities are indeed functions of the density n but both are modified by the same factor $(1 + \alpha n)^2$ making their ratio independent of the density to this approximation.

COUPLED-CHANNEL BORN-APPROXIMATION CALCULATION OF TWO-NUCLEON TRANSFER REACTIONS IN DEFORMED NUCLEI. T. Tamura, D. R. Bes, R. A. Broglia, and S. Landowne [Phys. Rev. Lett. 25, 1507 (1970)].

The factor i^l in Eq. (4) was not included in the computation. Therefore, the sign of F_2 and, consequently, the theoretical curves in Figs. 1(a)-1(c) were erroneous. Corrected results that replace those in the old Fig. 1(a) are presented in the new figure given here. The calculations were made with $0^+ - 2^+ - 4^+$ coupling in both incident and exit channels, everything else being the same as in previous calculations. Note the improved agreement with experiment, in particular concerning the 2^+ cross section. Qualitative statements concerning Figs. 1(b) and 1(c), in particular those concerning the importance of the multi-step processes, remain correct. We are very much indebted to Dr. T. Udagawa and Dr. B. Sorensen for their kind cooperation in finding and confirming the above error.

