

EXCITED ROTATIONAL BAND IN "SPHERICAL"  $^{150}\text{Sm}$  †

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Evidence is presented from the reaction  $^{152}\text{Sm}(p,t)^{150}\text{Sm}$  for the existence of an excited rotational band based on the 1.256-MeV  $0^+$  state of  $^{150}\text{Sm}$ , with  $2^+$  and  $4^+$  members at 1.417 and 1.819 MeV.

The low-lying states of nuclei around  $N=88$  undergo a transition in character from harmonic vibrational ( $N<88$ ) to rotational ( $N>88$ ). This can be interpreted as indicating a shape change from spherical to deformed. However, these nuclei, especially those close to  $N=88$ , are quite soft to shape variation and probably have zero-point vibrations whose amplitude is comparable with their average deformation.<sup>1</sup> In the following we will refer to "spherical" and "deformed" states, realizing that this describes only in some average way the nuclear shape.

$^{150}\text{Sm}$  ( $N=88$ ) and  $^{152}\text{Sm}$  ( $N=90$ ) are spherical and deformed, respectively, in their ground states, but have excited  $0^+$  states which are believed to have shapes quite unlike the ground states. The 1.091-MeV  $0^+$  state of  $^{152}\text{Sm}$  is strongly populated by the  $(t,p)$  reaction from the spherical ground state of  $^{150}\text{Sm}$ ,<sup>2,3</sup> in contrast to the usual weak population of excited  $0^+$  states by the  $(t,p)$  and  $(p,t)$  reactions. Hinds *et al.*<sup>2</sup> interpret this as indicating a spherical character for the state. Further support for this picture comes from a recent  $(p,t)$  experiment<sup>4</sup> in which this state is not seen in pickup from deformed  $^{154}\text{Gd}$ .

Likewise, the unusual strength with which the  $(p,t)$  reaction connects the deformed ground state of  $^{152}\text{Sm}$  with the 1.256-MeV  $0^+$  state of  $^{150}\text{Sm}$ <sup>5</sup> indicates that the latter is a deformed state of a nucleus whose ground state is spherical. If the 1.256-MeV state is indeed deformed, it must serve as the basis for a rotational band. The resolution of the earlier  $(p,t)$  experiment left unanswered the question of whether such a rotational band exists.

We are currently studying the  $(p,t)$  reaction at 19 MeV for a range of rare-earth nuclei. Targets of  $^{152}\text{Sm}$  on C backing have been bombarded with 19-MeV protons from the John H. Williams Laboratory MP tandem. The resultant tritons were analyzed with a split-pole magnetic spectrometer and were recorded on nuclear emulsion. The overall energy resolution was approximately 10 keV. The spectrum of tritons recorded at  $\theta=25^\circ$  is displayed in Fig. 1. Table I gives the excitation energies, angular momenta, and

relative cross sections for the states observed in  $^{150}\text{Sm}$ .

The 1.256-MeV  $0^+$  state is strongly excited, as in the earlier experiment,<sup>5</sup> and candidates for the  $2^+$  and  $4^+$  members of its rotational band appear at 1.417 and 1.819 MeV, respectively. The 1.417-MeV state has been assigned  $J^\pi=2^+$  by Lure, Peker, and Prokofév<sup>6</sup> and by Smither and Buss.<sup>7</sup> Lure, Peker, and Prokofév suggested that the state is a rotation of the 1.256-MeV  $0^+$  state. The 1.819-MeV state has recently been assigned  $J^\pi=4^+$  by Smither and Buss<sup>7</sup> from resonance capture in the reaction  $^{149}\text{Sm}(n,\gamma)^{150}\text{Sm}$ . Thus, the angular momenta of these states are proper for a  $K=0$  rotational band.

Angular distributions of  $^{152}\text{Sm}(p,t)^{150}\text{Sm}$  with  $L=0$  leading to  $0^+$  final states are shown in Fig. 2. They have nearly identical shapes, while the various  $L=2$  angular distributions shown in Fig. 3 have dissimilar shapes. This dissimilarity of  $L=2$  angular distributions is reminiscent of that noticed by McLatchie *et al.* in the reaction  $^{154}\text{Sm}(p,t)^{152}\text{Sm}$ .<sup>4</sup> We have demonstrated that it is not a  $Q$ -value effect by raising the proton energy to 20 MeV and observing no change in the relative shapes of the  $L=2$  angular distributions. The  $L=2$  angular distributions are not sufficiently characteristic to serve as a basis for assigning angular momenta to the final states. Nevertheless,

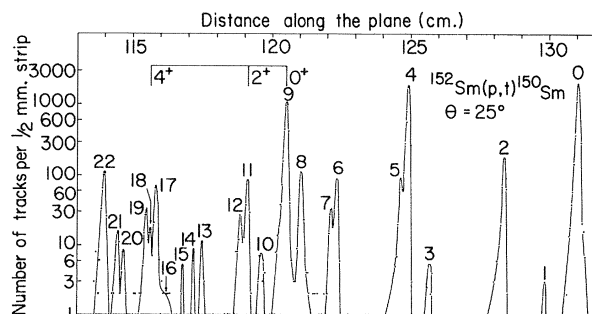


FIG. 1. Spectrum of tritons observed at  $\theta=25^\circ$  from the bombardment of a  $^{152}\text{Sm}$  target with 10-MeV protons. Levels in the residual nucleus  $^{150}\text{Sm}$  and contaminants are identified in Table I. The  $0^+$ ,  $2^+$ , and  $4^+$  members of the excited  $^{150}\text{Sm}$  rotational band are indicated.

Table I. Excitation energies,  $J^\pi$  values, and relative cross sections of the levels numbered in Fig. 1. States in  $^{150}\text{Sm}$  which have been observed previously are assigned the previously measured energies (without errors) and  $J^\pi$  values.<sup>a-e</sup> States not previously observed have been assigned excitation energies (with errors) by interpolation from known excitation energies.

| Level no. | $E_x$ (MeV)                   | $J^\pi$        | $\sigma_{\text{lab.}}$ ( $25^\circ$ ) (arbitrary units) |
|-----------|-------------------------------|----------------|---|
| 0         | 0.0                           | $0^+$          | $5537 \pm 250$  |
|           | $^{152}\text{Sm}$ g. s.       | $0^+$          | $< 66$  |
|           | $^{148}\text{Sm}$ g. s.       | $0^+$          | $< 31$  |
| 1         | $^{147}\text{Sm}$ g. s.       | $7/2^-$        | 7   |
| 2         | .334                          | $2^+$          | $527 \pm 23$  |
| 3         | $^{152}\text{Sm}$ (.685 MeV)  | $0^+$          | 17  |
| 4         | .741                          | $0^+$          | $4718 \pm 200$  |
| 5         | .774                          | $4^+$          | $185 \pm 25$  |
| 6         | 1.046                         | $2^+$          | $241 \pm 18$  |
| 7         | 1.072                         | $3^-$          | $78 \pm 15$   |
| 8         | 1.165                         | $2^+$          | $15 \pm 8$  |
|           | 1.194                         | $2^+$          | $346 \pm 20$  |
| 9         | 1.256                         | $0^+$          | $3172 \pm 100$  |
|           | 1.279                         | $3^+, (6^+)$   | $< 50$  |
| 10        | 1.357                         | $3^-$          | $30 \pm 10$   |
| 11        | 1.417                         | $2^+$          | $231 \pm 25$  |
| 12        | 1.449                         | $4^+$          | $89 \pm 20$   |
| 13        | $1.604 \pm .010$              | -              | $22 \pm 7$  |
| 14        | 1.642                         | $4^+$          | $16 \pm 6$  |
| 15        | 1.684                         | $3^-$          | $8 \pm 4$   |
| 16        | 1.759                         | $2^-$ or $3^-$ | $7 \pm 4$   |
|           | $^{152}\text{Sm}$ (1.775 MeV) | $3^-$ or $4^-$ | $< 7$   |
| 17        | 1.794                         | $2^+$          | $192 \pm 18$  |
| 18        | 1.819                         | $4^+$          | $41 \pm 9$  |
| 19        | 1.833                         | $2^+$ or $5^+$ | $91 \pm 12$   |
| 20        | $1.925 \pm .010$              | -              | $23 \pm 7$  |
| 21        | 1.949                         | -              | $46 \pm 8$  |
| 22        | $2.006 \pm .010$              | -              | $352 \pm 20$  |

<sup>a</sup>See Ref. 3.

<sup>b</sup>See Ref. 6.

<sup>c</sup>See Ref. 7

<sup>d</sup>R. K. Smither, Phys. Rev. 150, 964 (1966).

<sup>e</sup>C. M. Lederer, J. M. Hollander, and I. Perlman, *Table of Isotopes* (Wiley, New York, 1967), 6th ed.

the angular distribution for the 1.417-MeV rotational state is at least compatible with the assignment of  $J^\pi = 2^+$  to the state. Thus, the  $(p, t)$  angular distributions for the excited rotational states support the angular momentum assignment for the  $0^+$  member and are consistent with the assignment for the  $2^+$  member. The angular distribution of the 1.819-MeV state shown in Fig. 3 cannot be used to confirm the  $4^+$  assignment.

Further evidence for the rotational nature of these states comes from their excitation energies. Figure 4 shows the energy spacings of the ground-state rotational bands in deformed even-even samarium nuclides and of the proposed ex-

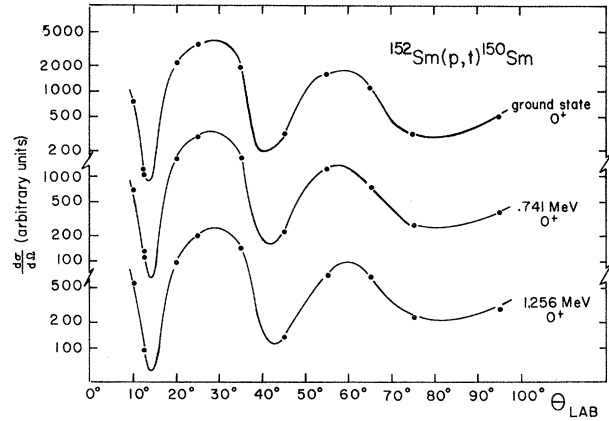


FIG. 2.  $L=0$  angular distributions of tritons from the reaction  $^{152}\text{Sm}(p,t)^{150}\text{Sm}$ . The cross sections for three different  $0^+$  final states in  $^{150}\text{Sm}$  are given in the same arbitrary units.

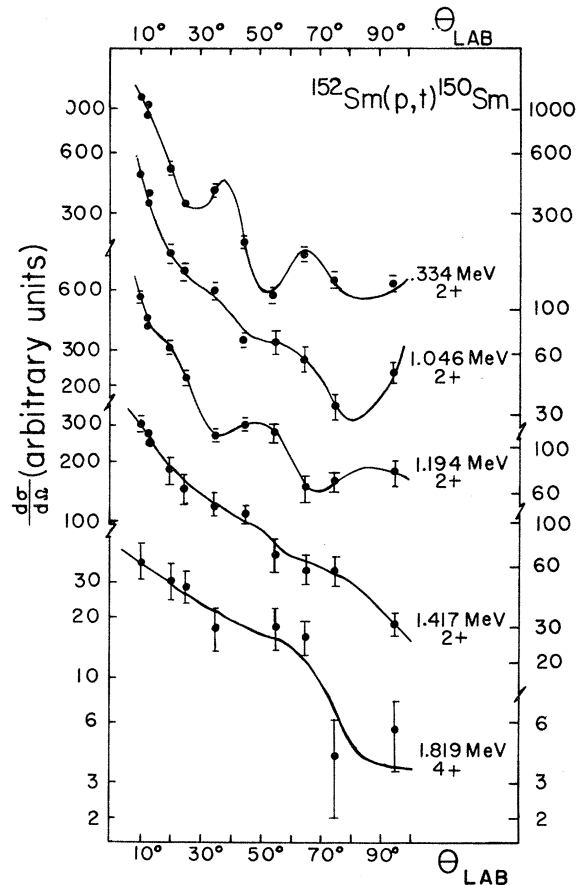


FIG. 3. Triton angular distributions from  $^{152}\text{Sm}(p,t)^{150}\text{Sm}$  leading to some final states in  $^{150}\text{Sm}$  with  $J^\pi = 2^+$  ( $L=2$ ) and  $J^\pi = 4^+$  ( $L=4$ ). The  $L=4$  angular distributions are all poorly defined due to the proximity of the  $4^+$  states to strong neighboring states. The angular distribution shown for the 1.819-MeV  $4^+$  state is typical in this regard.

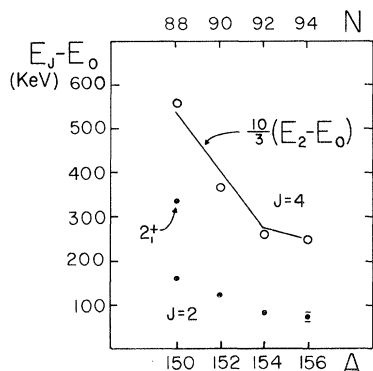


FIG. 4. Energy spacings of  $K=0$  rotational bands in even-even samarium nuclides. The energy of the state  $J$  relative to the  $0^+$  band head is shown for  $J=2$  (closed circles) and for  $J=4$  (open circles). The solid line connects the energies which are predicted from a  $J(J+1)$  formula. The spacings shown are for the ground-state bands of  $^{152,154,156}\text{Sm}$  ( $E_0=0$ ) and the excited band of  $^{150}\text{Sm}$  ( $E_0=1.256$  MeV). For  $^{150}\text{Sm}$ , the energy of the first-excited  $2^+$  state ( $2_1^+$ ) relative to the ground state is shown for comparison.

cited rotational band in  $^{150}\text{Sm}$ . The adherence of these spacings to a  $J(J+1)$  dependence supports the rotational picture. The systematic increase in the rotational energies with decreasing neutron number reflects a decrease in the moments of inertia. This systematic  $N$  dependence of the moment of inertia indicates that the intrinsic shape of the excited  $^{150}\text{Sm}$  rotational band is smoothly related to the ground-state shapes of the deformed samarium nuclides.

The strengths with which the  $(p, t)$  reaction populates the  $2^+$  and  $4^+$  members of the excited  $^{150}\text{Sm}$  rotational band, relative to the  $0^+$  member, are  $0.20 \pm 0.04$  and  $0.03 \pm 0.01$ , respectively. The numbers are the relative differential cross sections summed over a distance set of reaction an-

gles. We have measured the corresponding strengths for the ground-state rotational band of  $^{152}\text{Sm}$  in the reaction  $^{154}\text{Sm}(p, t)^{152}\text{Sm}$ ,  $E_p=19$  MeV. They are  $0.272 \pm 0.015$  and  $0.045 \pm 0.005$ . Broglia, Riedel, and Udagawa<sup>8</sup> have shown that the relative strength to the state  $J$  observed in two-neutron-transfer reactions on deformed nuclei measures the  $J$ th multipole component of the deformation carried by the transferred pair. Thus, the similarity of the relative strengths points to a similarity in intrinsic shape for the ground-state  $^{152}\text{Sm}$  and excited  $^{150}\text{Sm}$  rotational bands.

In summary, the levels at 1.256, 1.417, and 1.819 MeV in  $^{150}\text{Sm}$  have angular momenta, energy spacings, and two-neutron pickup intensities which identify them as members of a rotational band with properties similar to the ground bands of the neighboring samarium nuclides.

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<sup>2</sup>S. Hinds, J. H. Bjerregaard, O. Hansen, and O. Nathan, Phys. Lett. **14**, 48 (1965).

<sup>3</sup>J. H. Bjerregaard, O. Hansen, O. Nathan, and S. Hinds, Nucl. Phys. **86**, 145 (1966).

<sup>4</sup>W. McLatchie, J. E. Kitching, and W. Darcey, Phys. Lett. **36B**, 529 (1969).

<sup>5</sup>J. R. Maxwell, G. M. Reynolds, and N. M. Hintz, Phys. Rev. **151**, 1000 (1966).

<sup>6</sup>E. Ya. Lure, L. K. Peker, and P. T. Prokof'ev, Izv. Akad. Nauk SSSR Ser. Fiz. **32**, 74 (1968) [Bull. Acad. Sci. USSR, Phys. Ser. **32**, 74 (1968)].

<sup>7</sup>R. K. Smither and D. J. Buss, Bull. Amer. Phys. Soc. **15**, 86 (1970).

<sup>8</sup>R. A. Broglia, C. Riedel, and T. Udagawa, Nucl. Phys. **A135**, 561 (1969).