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¹²For nonzero external four-momentum at least a logarithmic divergence presumably develops, as in the $\pi^+-\pi^0$ mass shift. However, in principle, these weak-electromagnetic vertices could be less singular. We adopt a skeptical view of the finite result in Eq. (8) similar to that which one can take with respect to the $\pi^+-\pi^0$ mass difference in Eq. (1)—this in addition to the skepticism accompanying the hypothetical basis of the calculation.

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EXTERIOR THREE-PARTICLE WAVE FUNCTION. H. Pierre Noyes [Phys. Rev. Letters **23**, 1201 (1969)].

The equation for $\cos \mu_{i_s}$ in Eq. (3) should read

$$\cos \mu_{i_s} = \left[\frac{m_i m_s}{(m_i + m_s)(m_s + m_{s'})} \right]^{1/2}.$$

In order for the complete set $u_p^{j_i}(x) f_\lambda(qy)$ to have eigenvalue z in the three-particle Hilbert space, the first factor should be restricted to $p^2 = z - q^2$. This can be accomplished by writing the representation given in the line before Eq. (6) as

$$\begin{aligned} u_p^{j_i}(x) W^i(x) \\ = -(2/\pi) p \exp(-i\delta_\rho) \int_0^\infty dk k t_j^i(k, p; z - q^2). \end{aligned}$$

Since in the one-variable equation subsequently derived only the half-off-shell two-particle t matrix occurs, nothing in this paper is affected by this change. The bracket in the sentence following Eq. (10) should read: “[which lies in the strip bounded by $y_s = (x_s \cos \mu_{i_s} \pm R) / \sin \mu_{i_s}$].”