

favor a 1^+ D -wave assignment for the A_1 . However, in this analysis, a background amplitude having the angular dependence given by one-pion exchange has been subtracted.

¹¹There are also a variety of theoretical predictions for the properties of the A_1 . The literature can be traced from the following papers: S. Weinberg, in Proceedings of the Fourteenth International Conference on High Energy Physics, Vienna, Austria, September, 1968 (CERN Scientific Information Service, Geneva, Switzerland, 1968), and Phys. Rev. Letters **22**, 1023 (1969); P. Brooker and J. M. Kosterlitz, "Hard pion results from the Veneziano model" (to be published).

¹²I. T. Drummond, P. V. Landshoff, and W. J. Zakrewski, Cambridge University Report No. DAMTP 69/7 (to be published).

¹³M. Gell-Mann, M. L. Goldberger, F. E. Low, E. Marx,

and F. Zachariasen, Phys. Rev. **133**, B145 (1964).

¹⁴C. D. Froggatt and G. Ranft, to be published.

¹⁵K. Watson, Phys. Rev. **88**, 1163 (1952); J. S. Ball, W. R. Frazer, and M. Nauenberg Phys. Rev. **128**, 478 (1962); J. Gillespie, Final State Interactions (Holden-Day Publishers, San Francisco, Calif., 1964).

¹⁶J. D. Jackson, Nuovo Cimento **34**, 1644 (1964).

¹⁷The partial-wave amplitudes A_L^{JM} are of course not bounded by the unitarity circle familiar in elastic scattering phase shift analysis.

¹⁸If all the $J^P=1^+$ partial waves are dominated by resonance production then factorization implies that the ratio of S -wave to D -wave amplitudes should be the same for $M=0$ and $M=1$. Since an appreciable effect occurs only in the $M=0$ S -wave amplitude, this cannot be used as a practical test of local duality for the double-Regge amplitude.

ERRATA

INTERPRETATION OF RECENT EXPERIMENTAL TESTS OF VECTOR-MESON DOMINANCE. W. Schmidt and D. R. Yennie [Phys. Rev. Letters **23**, 623 (1969)].

The last line of Table I ($t = -10m_\pi^2$) for C^{\parallel} should read 0.35, 0.38, 0.39 at $E_\gamma = 4, 8, \text{ and } 16$ GeV.

MAGNETOELECTRIC EVIDENCE FOR THE ATTAINABILITY OF TIME-REVERSED ANTIFERROMAGNETIC CONFIGURATIONS BY METAMAGNETIC TRANSITIONS IN DyPO_4 . George T. Rado [Phys. Rev. Letters **23**, 644 (1969)].

On p. 646, column 1, the tenth line above Eq. (3) should read "...probably not a single..."

OMEGA PRODUCTION IN $\pi^+d \rightarrow \pi^+\pi^-\pi^0pp$ AT 4.19 GeV/c. G. S. Abrams, B. Eisenstein, and H. Gordon [Phys. Rev. Letters **23**, 673 (1969)].

(1) Page 674, column 2, line 6 now reads "The best fit, shown in Fig. 1(d), ..."; it should be "The best fit, shown in Fig. 1(a), ..."

(2) Page 675, column 2, the equation $\pi^+n \rightarrow \omega^0N^{*++}(1236)$ should be $\pi^+p \rightarrow \omega^0N^{*++}(1236)$.

(3) Page 676, footnote 12, should be:

$$f_{++}^1 = S^J \left[\alpha_+ \alpha_\rho + b_+ \left(\frac{t}{4m_\pi^2} - 1 \right) \right] \frac{m_\pi \tau_{\pi\omega}}{2s_0} \left(\frac{s-u}{2s_0} \right)^{\alpha_\rho - 1} \sin\theta_t \left(\frac{\alpha_\rho}{\Gamma(\alpha_\rho + 1)} \right),$$

$$f_{+-}^1 = S^J \alpha_+ t^{1/2} (1 + \cos\theta_t) \frac{\tau_{\pi\omega}}{2s_0} \left(\frac{s-u}{2s_0} \right)^{\alpha_\rho - 1} \frac{\alpha_\rho^2}{2\Gamma(\alpha_\rho + 1)},$$

$$f_{-+}^1 = S^J \alpha_+ t^{1/2} (1 - \cos\theta_t) \frac{\tau_{\pi\omega}}{2s_0} \left(\frac{s-u}{2s_0} \right)^{\alpha_\rho - 1} \frac{\alpha_\rho^2}{2\Gamma(\alpha_\rho + 1)},$$

$$f_{++}^1 = S^J \alpha_- \left(\frac{\tau_{pp}}{2s_0} \right) \left(\frac{s-u}{2s_0} \right)^{\alpha_B - 1} \frac{\alpha_B}{2\sqrt{2}\Gamma(\alpha_B + 1)} \sin\theta_t,$$

$$f_{++}^0 = S^J \left[\alpha_B \alpha_- \frac{t + m_\omega^2 - m_p^2}{2m_\omega} + b_1 (1 - \alpha_B) \tau_{\pi\omega} \right] \frac{\tau_{pp} t^{-1/2}}{4s_0 \Gamma(\alpha_B + 1)} \left(\frac{s-u}{2s_0} \right)^{\alpha_B - 1} \cos\theta_t,$$

where ...