port No. UCRL-8030, 1969 (unpublished), Pt. I, Footnote (n).
[8]B. French, in Proceedings of the Fourteenth International Conference on High Energy Physics, Vienna, Austria, September, 1968 (CERN Scientific Information Service, Geneva, Switzerland, 1968), p. 95; M. Aguilar-Benitez et al., CERN Report No. D. Ph. II/PHYS 69-3, 1969 (unpublished).

ERRATA

ALIGNMENT-INVERSION WALLS IN NEMATIC LIQUID CRYSTALS IN THE PRESENCE OF A MAGNETIC FIELD. W. Helfrich [Phys. Rev. Letters 21, 1518 (1968)].

It was wrongly assumed that $s_1s_2+t_1t_2=0$ if no surface torques are in operation. However, the equation does hold if the alignment inversion wall is plane, as are the walls treated quantitatively in the text.

GENERAL PARAMETRIZATION OF TRAJECTORY AND RESIDUE FUNCTIONS FOR DAUGHTER REGGE POLES. Loyal Durand, III, Paul M. Fishbane, and L. M. Simmons, Jr. [Phys. Rev. Letters 22, 261 (1969)].

The result for the daughter residue functions given in Eq. (4) of this paper is the most general possible only for the case of unequal-mass scattering $m_1\neq m_2$, $m_3\neq m_4$. For full generality, one must use an alternative expression,

$$\bar{\beta}_n(s)=(-1)^n\left[\frac{\Gamma(n+1)\Gamma(2\bar{\alpha}_n+2)}{\Gamma(2\bar{\alpha}_n+2-n)}\right]\left[\frac{\Gamma(\bar{\alpha}_n-n+1)}{\Gamma(\bar{\alpha}_n+1)}\right]^2 B_n$$
$$\times\left\{\sum_{j=0}^{n}\frac{\Gamma(2\bar{\alpha}_n+2-n)}{\Gamma(2\bar{\alpha}_n+2-n-j)}(\sqrt{s})^j b_j^{(1,2)}C_{n-j}^{(\bar{\alpha}_n-n+1)}(\cosh\beta_1)\right\}\{(1,2)\rightarrow(3,4)\}. \quad (4')$$

The two expressions are completely equivalent for unequal-mass scattering processes in the sense that they lead to identical restrictions on the first $n-1$ derivatives of $\bar{\beta}_n(s)$ at $s=0$. [The higher derivatives of $\bar{\beta}_n(0)$ are not restricted.] On the other hand, the equal-mass limits of (4) and (4′) are quite different. In particular, the equal-mass limit of Eq. (4) contains fewer free parameters than are allowed by general analyticity arguments, a point which we had overlooked. The expression for $\bar{\beta}_n$ given above does not suffer from this defect. We would like to thank Dr. J. H. Weis for pointing out the problem with the published result.

COHERENT MICROWAVE RADIATION FROM BiSb ALLOYS. C. A. Nanney and E. V. George [Phys. Rev. Letters 22, 1062 (1969)].

In two places typographical errors have occurred when the units microseconds were substituted for nanoseconds. These are as follows:

Page 1063, column 2, the sentence beginning in line 10 should read, "At frequencies away from the apex only short (~250 nsec) rf bursts are observed because the current pulse is not quite flat and, unlike at the apex, the frequency tunes rapidly (see Fig. 1) with field."

Page 1064, column 1, the passage beginning 14 lines from the bottom should read, "However, the BiSb stops emitting completely in a time ~200 nsec after the current pulse is turned off. Furthermore, the narrow (250 nsec) rf pulses slightly away from the apex should have approximately the same relaxation time as at the apex, but their decay is nearly two orders of magnitude faster than suggested by the spontaneous emission relaxation time of ~3 μsec."