²H. Kleinpoppen, H. Krueger, and R. Ulmer, Phys. Letters 2, 78 (1962).

3D. Hils, H. Kleinpoppen, and H. Koschmieder, Proc. Phys. Soc. (London) 89, 35 (1966).

 $⁴H$. Hafner and H. Kleinpoppen, Z. Physik 198, 315</sup> (1967).

 5 A similar decrease, however not as strong, is shown in the L_{α} polarization as reported by W. O. Ott, W. E. Kauppila, and W. L. Fite, in the Fifth International Conference on the Physics of Electronic and Atomic Collisions, Leningrad, USSR, July, 1967 (to

be published), p. 534, and Phys. Rev. Letters 19, 1361 (1967).

 $6W$. L. Fite and R. T. Brackmann, Phys. Rev. 113,

1151 (1956); W. L. Fite, R. F. Stebbings, and R. T.

Brackmann, Phys. Rev. 116, 356 (1959).

 ${}^{7}D.$ J. T. Morrison and M.R. H. Rudge, Proc. Phys. Soc. (London) 89, 44 (1966).

 8 M. R. H. Rudge, Proc. Phys. Soc. (London) 86, 763 (1965).

 $^{9}P. G.$ Burke, S. Ormonde, and W. Whitaker, Proc. Boo. (London) 92, 919 (1967).

DIRECT MEASUREMENTS OF PREDISSOCIATION PROBABILITIES IN CH AND CD MOLECULES*

James E. Hesser and Barry L. Lutz

Department of Astrophysical Sciences, Princeton University, Princeton, New Jersey (Received 22 January 1968)

Quantitative measurements of nonradiative decay rates in the $C^2\Sigma^+$ state of CH and CD molecules are presented via absolute phase shift versus frequency data and are interpreted as direct measures of predissociation probabilities. The C-state decay rates measured in CH are approximately three times larger than those in CD and comparison with absorption oscillator strength measurements shows that the predissociation probabilities of the rotational levels observed in CH range from approximately 3×10^7 to 1.6 $\times10^8$ sec⁻¹, i.e., approximately 3 to 18 times the radiative transition probability.

The visible and near ultraviolet emission spectra of CH and CD have been known for some $time₁$ and the transitions in CH have been identified in solar, stellar, and interstellar absorption.² Einstein transition rates, A , have recently been measured for two of the three known emission systems in CH, viz., for $A^2\Delta - X^2I$ and $B^2\Sigma^-\rightarrow X^2\Pi.^3,4$ Herbig⁵ has deduced an approximate absorption oscillator strength for the $C^2\Sigma^+$ -X²II transition at 3144 Å relative to the $A - X$ and $B - X$ transitions from the observed equivalent widths of the interstellar CH lines in the spectrum of ζ Ophiuchi. Linevsky⁶ has recently reported the first laboratory measurement of a relative f value for the $C-X$ system, using a King furnace; his value of $f(C-X)/f(A-X)$ $= 1.26$, if put on an absolute scale using Fink and Welge's⁴ A -X transition probability, leads to a radiative lifetime of 115 nsec for the C state. The corresponding value derived from Herbig's estimate would be ⁴³⁰ nsec.'

In recent flash-photolysis experiments with CH,N, and CD,N, Herzberg' noted that although $C-X$ fluorescence is observed in the spectrum of the deuterated compound, it is absent or weaker by a factor of at least 5 in the normal compound. He has interpreted this intensity difference between the two isotopes as evidence for a weak predissociation affecting the C state

of CH even though the $C-X$ transition of CH is regularly seen in emission in discharges and flames, where it is, however, much weaker than the $A - X$ and $B - X$ emissions in both CH and CD. No anomalies have been reported in the rotational intensity distribution of the C- X emission.^{1,9}

We have measured the decay rate of the $C^2\Sigma^+$, $v' = 0$ level in CH and in CD at low pressures using the phase shift apparatus of Lawrence^{10,11} in which emission spectra are excited by an rf-modulated beam of low-energy electrons. We have found a substantially more rapid decay than is consistent with the lifetime for purely radiative decay expected from the absorption measurements of either Linevsky or Herbig.

 $CH₄$ of high chemical purity and $CD₄$ of high isotopic purity¹² were used for electron-beam excitation of the C-X emission of CH and CD and the phase shifts were determined using the Ne-II multiplet at 1908-1935 Å as an abso
lute phase reference.¹¹ Electron energies of lute phase reference.¹¹ Electron energies of 100 V were used for these measurements, after careful investigation had shown no dependence of the phase shifts on electron energy in the range from 90 to 200 V and no dependence on pressure in the excitation region over a workable range from about 1 to 11 μ . With equal pressures of the parent molecules in the

excitation region, the emission intensity of CH $C-X$, excited in CH_4 , is three to four times weaker than the corresponding CD $C-X$ emission, excited in CD_4 . The phase versus frequency curves of the $C-X$ transitions in CH and CD are shown in Fig. 1 along with theoretical phase versus frequency curves computed from $tan \varphi$ = $2\pi f \tau$ (f= modulation frequency, τ = radiative lifetime), for pure exponential decays with lifetimes of 6.⁵ and 55 nsec. Apart from some distortion of the measured phase curves due to the presence ot a range of decay rates (see below), it is immediately evident that: (1) The lifetimes present in CH and CD differ markedly from each other, with CH showing the shorter lifetimes; and (2) that all lifetimes throughout the range covered by either isotope are markedly shorter than the radiative lifetime corresponding to Linevsky's f value. These results immediately confirm Herzberg's interpretation of the comparative weakness of the CH emission as evidence for weak predissociation of the C state. They further show that an even weaker predissociation affects the C state of CD.

The observed phase curves in CH and CD (Fig. 1) cannot be represented by pure exponential decays with a unique lifetime for each isotope. The curves could in principle be interpreted in terms of unique C-state lifetimes combined with radiative cascading from longer lived, higher lying, excited states. However, no higher lying states have ever been observed in emission despite careful searches. The observations can be plausibly assigned to spectrally unresolved overlapping of emissions from Cstate rotational levels covering a range of predissociation probabilities. A numerical analysis of the phase data of Fig. 1 was carried out based on the simplifying assumption of overlapping of emission from states of only two distinct lifetimes. Good fits to the CH and CD phase shifts at all frequencies were obtained on that basis with lifetimes of 6 and 26 nsec in CH and of 18 and 55 nsec in CD. These values indicate the approximate ranges over which the lifetimes of the individual rotational levels of CH and CD must vary to produce the observed phase curves. The levels for which these find-.
ings apply are those contributing to an 8-Å wide region around the peak of the $C-X$ $(0,0)$ emission band, corresponding to the band selected by the monochromator for the phase measurements. The probabilities for predissociation

FIG. 1. Phase shift versus frequency data obtained for the $C^2\Sigma^+$ states of CH and CD. By comparison with the solid curves corresponding to exponential decays of 6.5 and 55 nsec, one may readily see the two effects reported in this work, namely the short mean lives of the $2\Sigma^+$ states in both CH and CD and the change in mean life in CD compared with CH.

corresponding to the observed spread of lifetimes are 3×10^7 to 1.6×10^8 sec⁻¹ in CH and 9×10^6 to 4.7×10^7 sec⁻¹ in CD, after subtract ing the radiative contribution of 9×10^6 sec⁻¹ as determined from Linevsky's⁶ and Fink and Welge's⁴ measurements. This is probably the weakest predissociation that has been recognized and positively identified in a case where an entire emission band is observed at low pressure over the complete range of expected J values without any obvious anomalies in the rotational intensity distribution. The results reported above also constitute the first direct, quantitative measurements of probabilities for predissociation through determination of decay constants.

The only dissociation limit lying below the v' =0 level of the $C^2\Sigma^+$ state is the limit corresponding to atomic C and H in their normal states, ${}^{3}P+{}^{2}S$. None of the molecular states correlating with that limit: $X^2\Pi$, $B^2\Sigma^-$, and the two predicted states ${}^{4}\Sigma^{-}$ and ${}^{4}\Pi$, would lead to an allowed predissociation of the $C^2\Sigma^+$ state (the ground state, $X^2\Pi$, can be excluded because

of the Franck-Condon principle). The variation of predissociation with rotation could be caused either by a modification of a rotationdependent, approximate selection rule (e.g., a change of coupling case in the 4II state with increasing rotation of the molecule), or by a variation of a tunneling probability due to the dependence of effective potential curves on rotational energy. Attempts will be made to determine whether the predissociation probability increases or decreases with J value in order to eliminate some of these possibilities.

We wish to express our gratitude to Dr. G. Herzberg of the National Research Council of Canada for many helpful comments and suggestions and to Dr. G. Herbig, Lick Observatory, Santa Cruz, California, for communicating his detailed results prior to publication. We are also indebted to Dr. K. Dressier for critical discussions throughout this work.

2See, for instance, T. Dunham, Jr., Publ. Am. Astron. Soc. 10, 123 (1940); L. Spitzer, Jr., and G. Field, Astrophys. J. 121, ³⁰⁰ (1955); P. Thaddeus and J. F. Clauser, Phys. Rev. Letters 16, 819 (1966); C. E. Moore and H. P. Broida, J. Res. Nat. Bur. Std., A 63

Moore and H. P. Broida, J. Res. Nat. Bur. Std. , ^A 63, 19, 279 (1959).

 3 R. G. Bennett and F. W. Dalby, J. Chem. Phys. 32, 1716 {1960).

 ${}^{4}E$. H. Fink and K. H. Welge, J. Chem. Phys. 46. 4315 (1967); we have repeated Fink and Welge's lifetime measurements of the A and B states of CH and have confirmed their results.

⁵G. Herbig, paper presented at International Astronomical Union Commission 14, Prague, 1967 (unpublished).

6M. J. Linevsky, J. Chem. Phys. 47, ³⁴⁸⁵ (1967).

 7 Herbig's estimate of the f value is, within the experimental error of the method, in agreement with the laboratory measurements of Linevsky.

 ${}^{8}G$. Herzberg, paper presented at International Astronomical Union Commission 14, Prague, 1967 (unpublished).

⁹A. M. Bass and H. P. Broida, A Spectrophotometric Atlas of the Spectrum of CH from 3000A to 5000A, U. S. National Bureau of Standards Monograph No. 24

{U.S. Government Printing Office, Washington, D.C., 1961).

 $10G$. M. Lawrence, J. Quant. Spectry. Radiative Transfer 5, 359 (1965); Q. M. Lawrence and B. D. Savage, Phys. Rev. 141, 67 (1966).

¹¹B. D. Savage and G. M. Lawrence, Astrophys. J. 146, 940 (1966).

The mass spectroscopic analysis of the 99.95% pure $CH₄$ provided by the Matheson Company was $CO₂ < 10$ ppm, $O_2 = 13$ ppm, $N_2 = 45$ ppm, $C_2H_6 = 12$ ppm, and C_3H_8 < 5 ppm. The isotopic purity of the Merck, Sharp and Dohme CD_4 was 99.13% while chemical impurities due to air were $<0.1\%$.

^{*}This research is supported by National Aeronautics and Space Administration Grant No. NsQ-414.

 1 T. Shidei, Japan J. Phys. 11, 23 (1936); L. Gerö, Z. Physik 117, 709 (1941), and 118, 27 (1941).