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## NUCLEAR-MAGNETIC-RESONANCE RELAXATION BY SPIN-ROTATIONAL INTERACTIONS IN THE SOLID

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The interaction of the nuclear magnetic dipole moment with the magnetic field produced by rotation of the molecule containing the nucleus is known to operate in gases and the modulation of this interaction by molecular collisions is in many cases rate determining for nmr relaxation. Until a few years ago, however, it was generally assumed<sup>1</sup> that anisotropic interactions completely quench the molecular rotational magnetic moment in the condensed phases and that spin-rotational interactions therefore vanish both in liquids and in solids. Gutowsky<sup>2</sup> was the first to point out that this is not always the case for liquids<sup>3,4</sup> and that spin-rotational interactions in some fluoromethanes significantly contribute to the nuclear magnetic relaxation. In the present Letter we report what we believe is the first experimental evidence that spin-rotational interactions do operate in some molecular solids and represent in a certain temperature range the dominant nmr spin-lattice relaxation mechanism.

The temperature dependences of the  $F^{19}$  spin-lattice relaxation times ( $T_1$ ) of  $SF_6$ ,  $SeF_6$ , and  $TeF_6$ , observed at 9.2, 15.2, and 23.1 Mc/sec by pulse methods, are presented in Fig. 1 and show, except at very low temperatures, the same anomalous behavior as found by Gutowsky<sup>4</sup> and Hubbard<sup>5</sup> to be characteristic for the presence of spin-rotational interactions. Three different  $T_1$  regions are apparent: the liquid (above  $-51^\circ C$  for  $SF_6$ ,  $-35^\circ C$  for  $SeF_6$ , and

$-38^\circ C$  for  $TeF_6$ ), the solid phase I (between the melting point and  $-182^\circ C$  for  $SF_6$ ,  $-146^\circ C$  for  $SeF_6$ , and  $-72^\circ C$  for  $TeF_6$ ), and the low-temperature solid phase II. Second-moment data demonstrate a high degree of reorientational freedom in the solid phase I: The intra-molecular dipolar interactions are averaged out and only the intermolecular interactions contribute to the observed  $F^{19}$  second moments. Linewidth transitions in  $SeF_6$  and  $TeF_6$  at  $-150^\circ C$  and  $-110^\circ C$ , respectively, demonstrate the freezing in of this motion in solid phase II.

In liquid  $SF_6$  the  $T_1$  results agree with the ones of Hubbard and Hackleman<sup>6</sup>: Here as well as in  $SeF_6$  and  $TeF_6$  the relaxation times are frequency independent and slowly increase with decreasing  $T$ , demonstrating that spin-rotational interactions are dominant.<sup>6</sup> The interesting thing is, however, that this trend is still preserved in the solid phase I even more than  $100^\circ$  below the melting point. After a small initial drop on going from the liquid to the solid, the relaxation times increase with decreasing temperature until in  $SeF_6$  and  $SF_6$  a flat maximum is reached. In  $SeF_6$  the flat plateau extends till the solid-solid transition point dividing phase I from phase II, whereas in  $SF_6$  the  $T_1$  starts to decrease with decreasing  $T$  in the region between  $-150$  and  $-180^\circ C$ . In  $TeF_6$  the solid-solid phase transition apparently occurs before the maximum in  $T_1$  is reached.

In all three compounds  $T_1$  does not depend

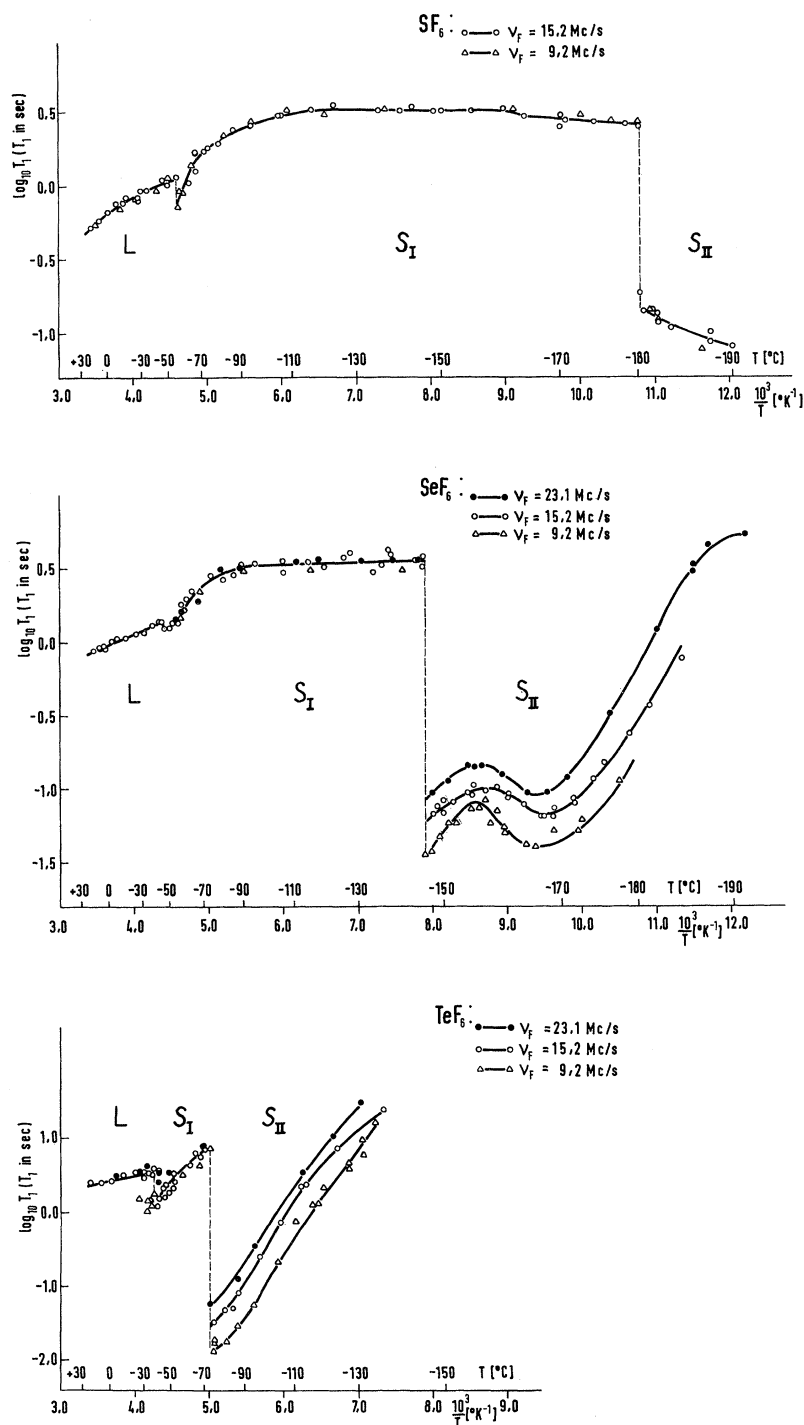


FIG. 1. Temperature dependences of the  $\text{F}^{19}$  spin-lattice relaxation times in  $\text{SF}_6$ ,  $\text{SeF}_6$ , and  $\text{TeF}_6$ .

on frequency in this temperature region so that  $\omega^2\tau_d^2 \ll 1$  (where  $\tau_d$  is the correlation time governing the exponential time decay of the autocorrelation functions for molecular orientation and position). If dipolar interactions

were to dominate the relaxation process, the condition  $\omega^2\tau_d^2 \ll 1$  would lead to

$$(1/T_1)_d = C_1 \tau_d; \tag{1}$$

and since

$$\tau_d = \tau_d^0 \exp(E_d/kT), \quad (2)$$

the relaxation time should decrease with decreasing temperature, contrary to observation. As will be shown later, the anisotropy of the chemical shift tensor  $\Delta\delta$  is too small to contribute significantly to relaxation and since the longest  $T_1$  observed in  $\text{TeF}_6$  is 32 sec, whereas the maximum  $T_1$  value in the solid phase I of this compound is 5 sec, paramagnetic impurities likewise cannot be responsible for the observed anomaly. The only plausible mechanism which remains appears to be the spin-rotation interaction.

Though no theory of spin-rotational relaxation exists as yet for solids, the similarity of the liquid and the solid-phase-I  $T_1$  behavior suggests that, as in case of the BPP theory,<sup>7</sup> some features of the models originally developed for liquids can be transferred to disordered solids. The "transient rotation" model<sup>4</sup> of Gutowsky, where spin-rotational interactions are assumed to operate only when the molecule acquires enough energy to overcome the potential barrier and is actually rotating, as well as the "rotational diffusion" model of Hubbard,<sup>5</sup> both predict for  $\omega^2\tau_d^2 \ll 1$  and spherical molecules a spin-rotational contribution to the relaxation time of the form

$$(1/T_1)_{\text{SR}} = C_2/\tau_d, \quad (3)$$

so that the resulting expression for  $T_1$ ,

$$T_1 = \frac{\tau_d}{C_2 + C_1(\tau_d)^2}, \quad (4)$$

as obtained from  $(1/T_1) = (1/T_1)_d + (1/T_1)_{\text{SR}}$ , qualitatively reproduces the experimental plots of  $T_1$  vs  $10^3/T$  in solid phase I except at lowest temperatures where  $C_2$  seems to become

temperature dependent. Though one should not ascribe too much significance to the numerical values of the constants occurring in Eq. (4), it is perhaps worthwhile to mention that a typical value of  $\tau_d$  at the  $T_1$  maximum is  $10^{-10}$  sec and of  $C_2$ ,  $10^{-11}$ . On going to solid phase II,  $\tau_d$  discontinuously increases and dipolar interactions become rate determining for spin-lattice relaxation, as can be seen from the frequency dependence of  $T_1$  and the occurrence of dipolar BPP minima. From the plots of  $1/T_1$  versus the inverse square of the Larmor frequency, the anisotropy contribution to the relaxation rate can be determined.  $\Delta\delta$  is found to be so small that it cannot influence relaxation in the  $T$  and frequency range considered.

The above results seem to throw new light on molecular reorientations and intermolecular interactions in disordered solids and open the possibility of studying correlation functions for the angular velocity of molecules not only in gases and liquids but also in the solid phase. Work along these lines is in progress.

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