GADOLINIUM HYPERFINE FIELDS

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The Gd hyperfine field, in the magnetically ordered state of several GdX $_2$ cubic Lavesphase compounds, shows a large variation in magnitude and sign. A comparison of our results with those on other rare-earth systems supports the supposition of a constant $4f$ core-polarization field in the rare earths.

We have measured the gadolinium hyperfine field in the magnetically ordered state at 4.2° K for several GdX₂ (X=A1, Rh, Pt, Fe, Mn) intermetallic Laves-phase compounds by nuclear-magnetic-resonance spin-echo techniques. Our results show a large variation in the magnitude of the Gd hyperfine field throughout this series. This is in sharp contrast to the essentially constant value of the Gd hyperfine field previously reported for Gd nuclei in a variety tially constant value of the Gd hyperfine field
previously reported for Gd nuclei in a variety
of environments, namely, Gd metal,^{1,2} GdN,³ previously reported for Gd nuclei in a varie
of environments, namely, Gd metal,^{1,2} GdN
and Gd³⁺ ions.^{4,5,6} As a result of the earlie: investigation, which yielded a field of approximately -340 kOe (the negative sign implying that H_{eff} is directed opposite to the local magnetization), it was believed that the major contribution to the Gd hyperfine field comes from core polarization by the unfilled 4f shell. Our results point out the importance of other contributions to the hyperfine field which we will call "valence" polarization.

Our experimental results on the Gd hyperfine fields in the cubic Laves-phase intermetallic compounds and in GdN are summarized in Table I. All of these data were obtained on magnetically ordered samples at 4.2 K in zero applied magnetic field. In these systems, where

Table I. Gadolinium hyperfine fields.

	H_{eff} at 4.2°K
Compound	$(kOe)^a$
GdAl ₂	-170
GdRh ₂	-103
GdPt,	180 ^b
GdFe ₂	$+453$
GdMn ₂	$285 \pm 30^{\circ}$
GdN	370 ^b

 a_{μ} (¹⁵⁷Gd) = 0.32 μ _N.

bThe signs of these fields are uncertain but are believed to be negative.

the Gd sites have cubic symmetry, the spinecho spectra showed two maxima of approximately equal amplitude which occur at frequencies in the ratio of the 155 Gd and 157 Gd nuclearly magnetic moments. We note that the magnitude of the Gd hyperfine fields in these samples varies from 103 kOe in GdRh, to 453 kOe in GdFe, . Furthermore, the sign of the field at Gd in $GdFe₂$, as determined by the frequency shift produced by a static applied magnetic field, is positive with respect to the local magnetization direction. In the light of previous data giving negative fields for Gd^{3+} , this positive field establishes the presence of a large positive contribution in GdFe, which is presumably due to a significant polarization of the valence electrons. Such a large positive contribution needed to explain our results in GdFe, has also been found necessary to explain Mössbauereffect data in other RFe , $(R = Tb, Dy, Er, Tm)$ compounds. In contrast with this latter series of rare-earth compounds, Gd is expected to have a negligible orbital contribution to the hyperfine field.

In the absence of any orbital contribution to the hyperfine field at a Gd nucleus, the effective field can be written as the sum of two terms:

$$
H_{\text{eff}} = H_{\text{core}} + H_{\text{val}},
$$

where H_{core} is the contribution due to polarization of the core electrons by the unfilled $4f$ shell and H_{val} is the contribution due to the valence electrons. In writing H_{eff} this way we are assuming that the various mechanisms responsible for valence-electron polarization can be collected into a single term and merely superimposed on a constant core-electron contribution. On the basis of the reported field at the Gd nuclei in Gd metal, GdN, and the Gd^{3+} ion, we put $H_{\text{core}} = -340$ kOe. Since $H_{\text{eff}} = +453$

kOe in GdFe₂, H_{val} must be of the order of 800 kOe in that compound. It is significant that the hyperfine fields at Tb, Dy, Tm, and Er, in TbFe₂,⁷ DyFe₂,⁸ TmFe₂,⁹ and ErFe₂,¹⁰ respectively, are also about 800 kOe higher than in tively, are also about 800 kOe higher than in
the corresponding pure metals Tb ,⁷ Dy,⁸ Tm,¹¹ and $Er.^{12}$ Since the fields in these metals are re:
12 already large and positive because of the orbital contribution, there is no change in the sign of the hyperfine field in going from the pure metals to the $RFe₂$ systems as occurs in going from Gd metal to $GdFe_2$. It is interesting to note that the numerical difference between the metal and $RFe₂$ hyperfine fields for these systems is essentially independent of the large variation of the orbital contribution to the hyperfine field. This suggests that the polarization mechanism responsible for this difference is confined to the valence $(5d,$ 6s) electrons of these rare earths, and that the outer electrons are weakly coupled to the $4f$ electrons. It is probable that the mechanism responsible for this difference in hyperfine fields is the same as that responsible for the high Curie temperatures of all of these $RFe₂$ compounds as compared with the metals, and could be caused by overlap of d -character electron wave functions localized about the rare earth and iron sites, respectively. The importance of the d character of the valence electrons in Gd metal has already been demonstrated in the recent band-structure calculation of Dimin the recent band-structure calculation of Di
mock and Freeman,¹³ and we would expect the d character of the conduction band to be especially important in the $RFe₂$ systems.

In Fig. 1 we compare the hyperfine field at Dy nuclei as determined by Mössbauer experiments¹⁴ with our Gd nmr results in the rareearth Laves-phase systems as a function of $R-R$ nearest neighbor $(n.n.)$ distance $(normal-)$ ized to the Gd-Gd n.n. distance in $GdAl₂$). The similarity in the variation of the rare-earth hyperfine fields in these two series of Lavesphase compounds suggests that (1) the 4f orbital contribution (which is the dominant contribution in the Dy compounds and is negligible in the Gd compounds) is essentially unchanged throughout each series and is largely determined by the crystal symmetry, and therefore (2) the $4f$ core-polarization field H_{core} is also constant throughout each series. This similar variation of the Gd and Dy fields in the presence of the large difference in the absolute magnitude of their respective hyperfine fields

FIG. 1. Dy and Gd hyperfine fields as a function of the rare-earth separation, normalized to that in GdAl₂. The asterisks indicate that the signs of the GdPt, and GdMn, fields are uncertain but are believed to be negative.

(6.² MOe in Dy compounds and 0.34 MOe in the Gd compounds) strongly supports the constancy of the 4f core-polarization field in the rare earths and emphasizes the importance of "valence" electron polarization.

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