perature display a first-order transition in  $UO_2$ , but in the lattice properties, in contrast to Blume's hypothesis. This result would suggest a re-examination of the transition mechanism in  $UO_2$ . Part of the data on elastic constant versus temperature is in accord with the idea of a first-order phase change, but  $C_{44}$  versus temperature would seem to be quite anomalous and indicative of a considerable interaction between lattice and spins well above the transition.

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<sup>1</sup>M. Blume, Phys. Rev. 141, 517 (1966).

<sup>2</sup>B. T. Willis and R. I. Taylor, Phys. Letters <u>17</u>, 188 (1965).

<sup>3</sup>B. C. Frazer, G. Shirane, D. E. Cox, and C. E. Olsen, Phys. Rev. <u>140</u>, A1448 (1965).

<sup>4</sup>M. R. Daniel, Phys. Letters <u>22</u>, 131 (1966).

<sup>5</sup>S. J. Allen, J. Appl. Phys. (to be published).

<sup>6</sup>K. B. Aring and A. J. Sievers, J. Appl. Phys. (to be published).

<sup>7</sup>O. G. Brandt and C. T. Walker, to be published.

<sup>8</sup>C. P. Bean and D. S. Rodbell, Phys. Rev. <u>126</u>, 104 (1962).

<sup>9</sup>G. Dolling and R. A. Cowley, Phys. Rev. Letters <u>16</u>, 683 (1966).

<sup>10</sup>C. W. Garland and C. F. Yarnell, J. Chem. Phys. <u>44</u>, 1112 (1966); C. W. Garland and R. Renard, <u>ibid</u>. <u>44</u>, 1120 (1966); R. Renard and C. W. Garland, <u>ibid</u>. <u>44</u>, 1125 (1966); C. W. Garland and R. Renard, <u>ibid</u>. <u>44</u>, 1130 (1966); and C. W. Garland and C. F. Yarnell, <u>ibid</u>. <u>44</u>, 3678 (1966).

## $\frac{3}{2}$ - GROUND STATE OF V<sup>47</sup> †

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Of those nuclei in the  $1f_{7/2}$  shell which are readily accessible to experimental investigation, V<sup>47</sup> is one of the least well studied. Thus, although in studies of the reactions  $Cr^{50}(p, \alpha)$ ,<sup>1</sup>  $Ti^{46}(p, \gamma)$ ,<sup>2</sup> and  $Ti^{47}(p, n)$ ,<sup>3</sup> energy levels have been observed up to about 3 MeV in excitation, in no cases have spin-parity assignments been made. In particular, the assumed<sup>4</sup>  $\frac{5}{2}^{-}$  character for the ground state of V<sup>47</sup> has been based largely on its observed  $\beta^+$  decay.

On the other hand, there are now extensive theoretical predictions concerning V<sup>47</sup> as well as a wide range of other  $1f_{7/2}$  nuclei. For example, within the framework of a pure  $1f_{7/2}$ shell-model configuration, McCullen, Bayman, and Zamick<sup>5</sup> and Ginnochio<sup>6</sup> have successfully reproduced the excitation energies of the low-lying  $\frac{7}{2}$  and  $\frac{5}{2}$  states for many odd-A nuclei. However, their calculations generally predict the  $\frac{3}{2}$  states arising from this configuration to be as much as 1 MeV higher in excitation than their observed locations. More recently, Federman and Talmi<sup>7</sup> have obtained good agreement for the low-lying states in the Ca isotopes by assuming a deformation of the  $1f_{7/2}$  orbital. Malik and Scholz<sup>8</sup> have also shown that the level schemes for a number of odd-A nuclei in this shell can be rather well described in terms of the Nilsson model with Coriolis coupling. According to this model, the relative locations of the  $\frac{7}{2}$ ,  $\frac{5}{2}$ , and  $\frac{3}{2}$  states may be a sensitive measure of the deformations in these nuclei. In this communication we wish to report on a study of the low-lying  $\frac{3}{2}$ ,  $\frac{5}{2}$ , and  $\frac{7}{2}$  states in V<sup>47</sup> arising from the  $(1f_{7/2})^3$  proton configuration. Information concerning the locations of hole states is also presented.

Levels in V<sup>47</sup> have been studied<sup>9</sup> by means of the reaction Ti<sup>46</sup>(He<sup>3</sup>, d) using a 16.5-MeV He<sup>3</sup> beam from the University of Pennsylvania tandem accelerator. Deuteron spectra were recorded at angles ranging from 7° to 40° with a broad-range magnetic spectrograph and with an over-all resolution <20 keV. Figure 1 shows a partial deuteron spectrum measured at 30° for the transitions leading to the ground and first three excited states at 0.089, 0.147, and 0.258 MeV. The impurity group corresponding to the 1.663-MeV level in V<sup>49</sup> arose because of the presence of 14.5% Ti<sup>48</sup> in the target.

In Fig. 2, deuteron angular distributions corresponding to the ground and second excited states are shown. The curves shown were calculated from distorted-wave theory using the code JULIE. Of particular interest here is the unambiguous  $l_p = 1$  assignment to the ground-



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DISTANCE ALONG PLATE (cm)

FIG. 1. Deuteron spectrum from the reaction  $\text{Ti}^{46}(\text{He}^3, d) \text{V}^{47}$  measured at 30°. The numbers in parentheses are the values of the orbital angular momenta transferred in the reaction.



FIG. 2. Deuteron angular distributions corresponding to the ground state and 0.147-MeV state in V<sup>47</sup>. The curves are the results of distorted-wave calculations.

state transition. This clearly rules out the earlier  $\frac{5}{2}^-$  assignment for the ground state and limits the spin-parity to either  $\frac{1}{2}^-$  or  $\frac{3}{2}^-$ . However, of these possibilities, only  $\frac{3}{2}^-$  is consistent with the measured  $\log ft = 4.9$  in the decay  $V^{47}(\beta^+) Ti^{47}$ .<sup>4</sup>

The intense (S=0.73)  $l_p = 3$  transition suggests a  $\frac{7}{2}^-$  assignment for the level at 0.147 MeV. The  $\frac{5}{2}^-$  state arising from the  $(1f_{7/2})^3$  configuration is expected to lie close to this  $\frac{7}{2}^-$  level<sup>6,8</sup> and should be only weakly excited since it cannot be formed by a first-order stripping mechanism. It is, therefore, reasonable to identify this state with the level at 0.089 MeV to which a relatively weak, nonstripping transition was observed. The level at 0.258 MeV was seen to be excited by a weak  $l_p = 2$  transition, and this level can be identified as a  $\frac{3}{2}^+$ hole state. Only one other low-lying even-parity state was seen in this study. This is the 1.664-MeV state to which a  $l_p = 0$  transition was observed and which can, therefore, be identified as a  $\frac{1}{2}^+$  hole state.

It is interesting to note that the Coriolis coupling model employed by Malik and Scholz<sup>8</sup> can successfully account for the  $\frac{3}{2}$  ground state of V<sup>47</sup> only if a deformation as large as either  $\beta \approx -0.5$  or  $\beta \approx +0.6$  is assumed. However, if the over-all agreement with the level schemes of neighboring  $(1f_{7/2})^3$  nuclei is taken into account, their calculations tend to favor a negative deformation.

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<sup>6</sup>J. N. Ginnochio, Phys. Rev. <u>144</u>, 952 (1966).

<sup>7</sup>P. Federman and I. Talmi, Phys. Letters <u>22</u>, 469 (1966).

<sup>8</sup>F. B. Malik and W. Scholz, Phys. Rev. <u>150</u>, 919 (1966).
<sup>9</sup>B. Rosner and D. J. Pullen, Bull. Am. Phys. Soc.

<u>11</u>, 850 (1966).

<sup>&</sup>lt;sup>†</sup>Work supported by the National Science Foundation. <sup>1</sup>G. Brown, A. MacGregor, and R. Middleton, Nucl. Phys. <u>77</u>, 385 (1966).

 $<sup>^{2}</sup>$ H. Albinsson and J. Dubois, Phys. Letters <u>15</u>, 260 (1965).

<sup>&</sup>lt;sup>3</sup>G. J. McCallum, A. T. G. Ferguson, and G. S. Mani, Nucl. Phys. <u>17</u>, 116 (1960).

<sup>&</sup>lt;sup>4</sup><u>Nuclear Data Sheets</u>, compiled by K. Way <u>et al</u>. (Printing and Publishing Office, National Academy of Sciences-National Research Council, Washington 25, D. C.), p. 60-2-41.

<sup>&</sup>lt;sup>5</sup>J. D. McCullen, B. F. Bayman, and L. Zamik, Phys. Rev. <u>134</u>, B515 (1964).