†National Science Foundation Cooperative Fellow, 1963-1964.

<sup>1</sup>The report of a recent conference on this subject is given in <u>Metallic Solid Solutions</u>, edited by J. Friedel and A. Guinier (W. A. Benjamin, Inc., New York, 1963).

<sup>2</sup>For a modern discussion of alloys and additional references, see C. Kittel, <u>Quantum Theory of Solids</u>

(John Wiley & Sons, Inc., New York, 1963), Chap. 18. <sup>3</sup>J. S. Langer and S. H. Vosko, J. Phys. Chem. Solids <u>12</u>, 196 (1959). Also, see Kittel, Ref. 2, Chap. 6.

<sup>4</sup>For a recent review, see F. J. Blatt, in <u>Solid State</u> <u>Physics</u>, edited by F. Seitz and D. Turnbull (Academic Press, Inc., New York, 1957), Vol. 4.

<sup>5</sup>R. L. Odle and C. P. Flynn, to be published.

<sup>6</sup>N. F. Mott, Proc. Cambridge Phil. Soc. <u>32</u>, 281 (1936).

<sup>7</sup>A. Blandin and J. Friedel, J. Phys. Radium <u>21</u>, 689 (1960).

<sup>8</sup>We are aware of two previous measurements: N. S. Garif'ianov and M. A. Starikov, Zh. Eksperim. i Teor. Fiz. <u>35</u>, 798 (1958) [translation: Soviet Phys.-JETP <u>8</u>, 553 (1959)]; G. D. Wignall, J. E. Enderby, C. E. W. Hahn, and J. M. Titman, Phil. Mag. 12, 433 (1965).

<sup>9</sup>For the role of spin-orbit coupling in CESR relaxation, see R. J. Elliott, Phys. Rev. <u>96</u>, 266 (1954).

<sup>10</sup>See M. Hansen, <u>The Constitution of Binary Alloys</u> (McGraw-Hill Book Company, Inc., New York, 1958), 2nd ed.

<sup>11</sup>For a recent review of CESR spin-lattice relaxation, see Y. Yafet, in <u>Solid State Physics</u>, edited by F. Seitz

and D. Turnbull (Academic Press, Inc., New York, 1963), Vol. 14.

<sup>12</sup>For the proof of this, see Yafet, Ref. 11, and also T. R. Carver and C. P. Slichter, Phys. Rev. <u>102</u>, 975 (1956).

<sup>13</sup>CESR has been observed in the following metals: Li, Na, K, Rb, Cs, and Cu.

<sup>14</sup>The spectrometer was designed by Dr. C. H. Henry.
<sup>15</sup>J. L. Hall and R. T. Schumacher, Phys. Rev. <u>127</u>, 1892 (1962).

<sup>16</sup>F. J. Dyson, Phys. Rev. <u>98</u>, 349 (1955); also see G. Feher and A. F. Kip, Phys. Rev. <u>98</u>, 337 (1955).

 $^{17}$ For an elementary calculation of this contribution, see D. Pines and C. P. Slichter, Phys. Rev. <u>100</u>, 1014 (1955).

<sup>18</sup>We have studied the effect of having a distribution of particle sizes on the relation between  $\Delta H$  and  $T_1$  using Dyson's theory and find that  $\gamma \Delta H T_1 = 1.1$  holds to within ±10% for our particle distributions and linewidths.

<sup>19</sup>See, for example, J. M. Ziman, <u>Principles of the</u> <u>Theory of Solids</u> (Cambridge University Press, Cambridge, England, 1964).

<sup>20</sup>F. Herman and S. Skillman, <u>Atomic Structure Calcu-</u> <u>lations</u> (Prentice Hall, Inc., Englewood Cliffs, New Jersey, 1963).

<sup>21</sup>For the role of spin-orbit coupling and spin-flip scattering in the Knight shift in superconductors, see J. Appel, Phys. Rev. <u>139</u>, A1536 (1965).

<sup>22</sup>N. F. Mott and H. S. W. Massey, <u>The Theory of Atomic</u> <u>Collisions</u> (Clarendon Press, Oxford, England, 1965), 3rd ed., p. 257.

## $M_A$ CENTERS IN ADDITIVELY COLORED KCl

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Strong evidence now exists that an  $F_A$  center is simply an F center with an impurity ion at a nearest-neighbor site.<sup>1-4</sup> While similarly perturbed centers involving F aggregates would likewise be expected, their existence has not been clearly demonstrated. The purpose of the present note is to report evidence in KCl of the simplest such association, an M center with a Na<sup>+</sup> ion neighbor, i.e., an  $M_A$  center.

The usual *M* center consists of two nearestneighbor *F* centers along a  $\langle 110 \rangle$  lattice direction.<sup>5-6</sup> Its main transition, the *M* band, lies on the low-energy side of the *F* band. Figure 1 shows the absorption in this spectral region for three crystals which differ primarily in their Na<sup>+</sup> concentration. Curve *a* shows the *M* band in an undoped crystal with a natural content of 0.003 mole% NaCl. Compared to this spectrum, curve b reveals a weak absorption at about 820 m $\mu$  in a crystal with 0.04 mole% NaCl. The relative size of this shoulder absorption increases for crystals with increasing NaCl content until it predominates. This is evidently the case for a crystal containing 0.35 mole% NaCl, as shown in curve c. Similarly, the size of the  $F_A$  band (at about 580 m $\mu$ ) relative to that of the F band is very small for the crystal of Fig. 1, curve a, but is comparatively larger for crystals with increasing Na<sup>+</sup> concentration.

The *M* center has three symmetry axes along which its optical dipole moments lie. One is the  $\langle 110 \rangle$  vacancy axis, which is the direction for the *M* band. The other two are the  $\langle 110 \rangle$ and  $\langle 100 \rangle$  directions perpendicular to this axis. *M*-center transitions with moments along these normally overlap the *F* band. A common pro-



FIG. 1. Absorption spectra measured at liquid-helium temperature for three crystals differing in Na<sup>+</sup> content. Curve *a* refers to an undoped sample (actually containing 0.003 mole% NaCl); curve *b* refers to a sample with 0.04 mole% NaCl; curve *c* refers to a crystal with 0.35 mole% NaCl.

cedure for determining these is to measure the *F* absorption with polarized light after reorienting *M* centers with polarized *F* light at low temperatures.<sup>7-10</sup> It turns out that both the  $\langle 110 \rangle$  and  $\langle 100 \rangle$  transitions in the *F* region are effective in causing reorientations.

Specifically, to determine the  $\langle 110 \rangle M$ -center absorptions one might expose a crystal containing F and M centers to F light propagating in the [100] direction with electric vector oriented in the  $[0\bar{1}1]$  direction. The dichroic spectrum, which is the absorption difference, [011]  $-[0\bar{1}1]$ , for light again propagating in the [100] direction, contains only the  $\langle 110 \rangle M$ -center transitions. The reason for this is that the subtraction effectively removes absorptions due to isotropic centers like the F center, and also the  $\langle 100 \rangle M$ -center bands, since dichroism is not induced in these using this experimental arrangement.

The  $\langle 100 \rangle M$  transitions can be determined from  $\langle 100 \rangle$  dichroic spectrum using the same arrangement after orienting with light polarized in the [010] direction. This situation, however, is complicated by the fact that [010] light induces dichroism in both sets of transitions. However, the actual  $\langle 100 \rangle$  transitions are easily determined as follows: It can be shown that when the  $\langle 100 \rangle$  and  $\langle 110 \rangle$  dichroic spectra are scaled for equal *M*-band dichroism, then the contributions of the  $\langle 110 \rangle$  transition in the vicinity of the *F* band have the same magnitude but opposite sign. Thus, addition of the  $\langle 100 \rangle$  transition.

Figure 2 shows dichroic spectra for two crys-

tals in which the centers giving the absorption in the *M*-spectral region were oriented first with  $[0\overline{1}1]F$  light and then with [010]F light. The solid curves relate to a crystal with dichroism due almost entirely to the centers giving the  $820-m\mu$  absorption. This was accomplished by preparing a crystal with spectrum similar to Fig. 1, curve *c*, and exposing it to polarized *F* light at liquid-helium temperature for a relatively short period of time. A separation is possible since these centers reorient much more rapidly than *M* centers. One also avoids  $F_A$ center reorientations at this low temperature.<sup>3</sup> The dashed curves shown in Fig. 2 refer to a crystal with oriented *M* centers alone.



FIG. 2. Normalized dichroic spectra measured at liquid-helium temperature for two crystals; one containing only oriented  $M_A$  centers (the solid curves), the other containing oriented M centers alone (the dashed curves).  $M_A$  centers were oriented at liquid-helium temperature; M centers, at about 200°K. (a) The difference in absorption, [011]-[011], after orienting with [011] F light. These give the  $\langle 110 \rangle$  transitions. (b) The difference in absorption, [001]-[010], after orienting with [010] F light. (c) The sums ([001] - [010] + ([011]-[011]). These give the  $\langle 100 \rangle$  transitions.

It is evident from these curves that the two centers are similar in several important respects. Their absorptions in the *M* region exhibit both  $\langle 110 \rangle$  and  $\langle 100 \rangle$  dichroism, consistent for a  $\langle 110 \rangle$ -type center. Each has similar transitions in the *F* region with correspondingly oriented dipole moments. By contrast, the *M*-center's  $\langle 100 \rangle$  absorption bands in the *F* region lie at slightly longer wavelengths than those of the new center. Moreover, the new center has an absorption at 610 m $\mu$ , on the longwavelength side of the *F* band, not present in the *M*-center spectrum while the *M* center has a weak absorption at 634 m $\mu$  absent in the new center.

The over-all similarity in both the behavior and the absorption of the two centers leads one to conclude that the new center is simply a perturbed M center. From its dependence on impurity concentration it is very likely perturbed by a neighboring Na<sup>+</sup> ion.<sup>11</sup>

Unlike the F center, the M center has several nonequivalent nearest-neighbor sites. If these were equally likely, one would possibly observe several overlapping bands in the M region. The fact that one transition is observed predominantly is an indication that one configuration is most probable. From the over-all differences in the M and  $M_A$  spectra, it is likely that the Na<sup>+</sup> ion lies somewhere in the (100) plane containing the M center. The substitutional site along the  $\langle 110 \rangle$  direction bisecting the center would give a configuration with the highest coordination and symmetry, and is considered the most reasonable alternative.

Na<sup>+</sup> has always been a rather common impurity found in KCl crystals, particularly a number of years ago. Thus,  $M_A$  centers could have formed in earlier studies along with the usual M centers during the F-aggregation process. This is suggested by the variation in reported M-center properties such as the M-band peak value.<sup>12</sup> It is hoped that the results of the present investigation will provide a new basis for possibly understanding these former anomalous effects.

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<sup>1</sup>I. Masakazu and E. Sugioka, Phys. Rev. <u>110</u>, 1070 (1958).

<sup>2</sup>K. Kojima, N. Nishimaki, and T. Kojima, J. Phys. Soc. Japan 16, 2033 (1961).

<sup>3</sup>F. Lüty, Z. Physik <u>165</u>, 17 (1961); <u>177</u>, 369 (1964).
 <sup>4</sup>R. L. Mieher, Phys. Rev. Letters <u>8</u>, 362 (1962).

<sup>5</sup>C. Z. van Doorn, Philips Res. Rept. Suppl. No. 4, 89 (1962).

<sup>6</sup>For a discussion of earlier evidence of the  $F_2$  model, see W. D. Compton and H. Rabin, <u>Solid State Physics</u>, edited by F. Seitz and D. Turnbull (Academic Press, Inc., New York, 1964), Vol. 16. For conclusive resonance data see H. Seidel, Phys. Letters 7, 27 (1963).

<sup>7</sup>C. Z. van Doorn and Y. Haven, Phys. Rev. <u>100</u>, 753 (1955).

<sup>8</sup>F. Okamoto, Phys. Rev. <u>124</u>, 1090 (1961).

<sup>9</sup>C. J. Delbecq. Z. Physik <u>171</u>, 560 (1963).

<sup>10</sup>T. J. Turner, R. DeBatist, and Y. Haven, Phys. Status Solidi <u>11</u>, 535 (1965).

<sup>11</sup>The present discussion is confined to the relatively weakly doped samples in the range of Fig. 1. For more heavily doped samples, the M absorption becomes broader and diffuse, indicating the presence of more complicated structures.

<sup>12</sup>For a tabulation of these values, see the review article in Ref. 6.