
 E R R A T A

EXISTENCE OF PIONS WITH SPIN. G. Rinaudo, A. Marzari-Chiesa, G. Gidal, and A. E. Werbrouck [Phys. Rev. Letters 14, 761 (1965)].

The last line of the paragraph before acknowledgments should contain -3 instead of -4 as exponents. We sincerely thank Dr. S. Fung and Dr. S. Meyer of Rutgers University for pointing out this error.

OBSERVATION OF JAHN-TELLER TUNNELING BY ACOUSTIC LOSS. E. M. Gyorgy, M. D. Sturge, D. B. Fraser, and R. C. LeCraw [Phys. Rev. Letters 15, 19 (1965)].

The equation for τ^{-1} should be multiplied throughout by a factor $(kT)^{-2}$. We are grateful to Dr. P. L. Scott and others for drawing our attention to this error in transcription.

CALCULATIONS OF SOUND VELOCITIES IN CRYSTALLINE HELIUM AT ZERO TEMPERATURE. L. H. Nosanow and N. R. Werthamer [Phys. Rev. Letters 15, 618 (1965)].

In Table I, the transverse sound velocities in the 100 and 110 directions are incorrectly labeled. At molar volumes of 19.88 and 21.63 cc, the 100 sound velocities are, respectively, 343 and 335 meters per second; both are two-fold degenerate. The 110 sound velocities at 19.88 cc should be 174 and 343 meters per second; at 21.63 cc they should be 152 and 335 meters per second. All other values in this table are correct.

The transverse sound velocities in Fig. 2 are also incorrectly labeled. The curve which is now labeled "100" should be labeled "100(2) & 110." The curve which is now labeled "111(2) & 100" should be labeled only "111(2)." The

curve which is labeled "110(2)" should be labeled only "110."

APPLICATION OF CURRENT COMMUTATION RULES TO NONLEPTONIC DECAY OF HYPERONS. Hirotaka Sugawara [Phys. Rev. Letters 15, 870 (1965)].

The last equations at the bottom of page 871 should read

$$\begin{aligned} \langle B^{(\nu_1)} | T_{\nu}^{(27)} | B^{(\nu_2)} \rangle &= \begin{pmatrix} 8 & 27 & 8 \\ \nu_2 & \nu & \nu_1 \end{pmatrix} \langle \langle 8 | T^{(27)} | 8 \rangle \rangle, \\ \langle B^{(\nu_1)} | T_{\nu}^{(8)} | B^{(\nu_2)} \rangle &= \begin{pmatrix} 8 & 8 & 8s \\ \nu_2 & \nu & \nu_1 \end{pmatrix} \langle \langle 8 | T^{(8)} | 8 \rangle \rangle_s \\ &\quad + \begin{pmatrix} 8 & 8 & 8a \\ \nu_2 & \nu & \nu_1 \end{pmatrix} \langle \langle 8 | T^{(8)} | 8 \rangle \rangle_a. \end{aligned}$$

Added note.—Professor Y. Nambu has informed me that he and a collaborator¹¹ had used the same method in studying soft-pion emission. He also remarked that our Eqs. (A) and (B) are just the $\Delta I = \frac{1}{2}$ rule. This point was also pointed out by Suzuki.¹² He independently obtained the same results as ours and also observed that our equations are valid only for the *s* wave because of the *C* invariance in the strong interaction (the contribution to the *p* wave is identically zero) making our Eq. (D) consistent with experiment. In that case the *p* wave comes from the SU(3) broken amplitude and from the subtraction term. We cannot neglect contributions from the infinite past also.

¹¹Y. Nambu and E. Shrauner, Phys. Rev. 128, 862 (1962).

¹²M. Suzuki, this issue [Phys. Rev. Letters 15, 986 (1965)].