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OBSERVATION OF THE  $\lambda$  ANOMALY IN SOLID  $D_2$  BY NUCLEAR MAGNETIC RESONANCE\*

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Several unsuccessful attempts have been made by nuclear magnetic resonance techniques<sup>1-3</sup> to observe a cooperative transition in solid  $n-D_n$ (33% para) analogous to the  $\lambda$ -type anomaly in solid  $n-H_2$  (75% ortho). In addition, Smith and Housley<sup>4</sup> found no transition even in enriched samples of  $D_2$  with para  $D_2$  concentrations up to 55%. The transition in solid  $H_2$  occurs due to the quenching of the rotational motion of the ortho H<sub>2</sub> ( $I=1$ ,  $J=1$ ) molecules implying that the transition temperature  $T_{\lambda}$  is a sensitive function of the ortho  $H_2$  concentration<sup>5,6</sup> and the density.<sup>3-7</sup> Recent heat capacity measurements by Grenier and White<sup>8</sup> demonstrated the existence of such a cooperative transition in solid  $D_2$  above 1.5°K providing the para  $D_2$   $(I = 1, J = 1)$  concentration



FIG. 1. The derivative with respect to frequency of the absorption signal from an 81%  $pD_2$  sample at 4.2°K.

was in excess of 60%. We have made nuclearresonance studies of enriched para D, samples (with concentration of about 80%) and have observed the splitting of the resonance line in D, below  $2^{\circ}$ K.

The enriched para  $D_2$  samples were prepared in a manner similar to that suggested by Cunningham, Chapin, and Johnston.<sup>9</sup> The specimen was condensed from the gas phase into a Teflon sample cell surrounded by a small amount of liquid  $H<sub>2</sub>$  and then frozen quickly by transferring liquid He on top of the  $H_2$ . Figure 1 shows a derivative of the resonance line at 4.2 K obtained with a sample containing  $81\% p-D_2$ . This line is actually a superposition of lines arising from molecules with  $I = 2$  (ortho  $D_2$ ) and  $I = 1$  (para  $D_2$ ). Since the molecules with  $I = 2$  have  $J = 0$  at these low temperatures, one does not expect that quenching of rotation will have any effect on their contribution to the resonance line, in a first approximation. This trace was made with a frequency sweep rate of 45 cps/sec at a Larmor frequency of 2.63 Mc/sec. The width of the line shown is approximately 1.<sup>0</sup> gauss as compared to 'the value of about 1.5 gauss obtained in  $n - D_2$ .<sup>3-1</sup> The asymmetry of the line becomes more pronounced as the frequency sweep rate is increased.

Reduction of the temperature to about 2.2 K revealed the existence of side peaks split apart by  $76.8 \pm 0.5$  kc/sec. Figure 2 shows a trace of the resonance line at 1.2'K with 100 kc/sec reference marks. The magnitude of the splitting calculated by Reif and Purcell' is

 $\Delta v = 3\tau d$ ,





where d for the free  $D_2$  molecule has the value<sup>11</sup> 25.24 kc/sec. For  $D_2$ , d is a function of the strength of the dipolar interaction and the magnitude of the deuteron ni clear quadrupole moment. The calculated value of  $\tau$  is 1.01,<sup>2</sup> the 1% correction term arising from the contributions of higher order spherical harmonics in the perturbation calculation. Thus the calculated value of the splitting  $(\Delta v_{\text{calc.}} = 76.5 \text{ ke/sec})$  is in excellent agreement with the experimental value.

After the splitting of the line occurred, several other effects were observed: (i) The signal-tonoise ratio deteriorated rather badly; (ii) the line became so asymmetric that the recorder trace bore little resemblance to the conventional derivative curve; (iii) the remnant central line became noticeably broader, with a width of  $5-8$ gauss. All attempts to recover the symmetry of the line (such as reduction of the sweep rate and lowering the radiofrequency power) failed. Investigations into the origin of the asymmetry and

the apparent loss of signal after the  $\lambda$  transition are in progress.

Due to the relatively low signal-to-noise ratio, no attempt was made at this time to determine the transition temperature from the nmr data. The value of the transition temperature for 81 $%$  $p$ -D<sub>2</sub> derived from the heat capacity data is about 2.8'K.

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<sup>1</sup>J. Hatton and B. V. Rollin, Proc. Roy. Soc. (London) A199, 222 (1949).

 ${}^{2}$ F. Reif and E. M. Purcell, Phys. Rev. 91, 631 (1953).

<sup>3</sup>S. A. Dickson and H. Meyer, Proceedings of the International Conference on Low Temperature Physics, to be published.

 ${}^4G$ . W. Smith and R. M. Housley, Phys. Rev. 117, 732 (1960).

<sup>5</sup>T. Sugawara, Y. Masuda, T. Kanda, and E. Kanda, Sci. Repts. Research Insts., Tohoku Univ., Ser. A 7, <sup>67</sup> (1955).

 ${}^6G$ . W. Smith and C. F. Squire, Phys. Rev.  $111$ , 188 (1958).

 $N$ W. D. McCormick and W. M. Fairbank, Bull. Am. Phys. Soc. 3, 166 (1958).

 ${}^{8}G$ . Grenier and D. White, J. Chem. Phys. 40, 3015 (1964).

<sup>9</sup>C. C. Cunningham, D. Chapin, and H. L. Johnston, J. Am. Chem. Soc. 80, <sup>2382</sup> (1958).

<sup>10</sup>T. Sugawara, Sci. Repts. Research Insts., Tohoku Univ., Ser. A 8, 95 (1956).

 $N.$  F. Ramsey, Phys. Rev.  $85, 60$  (1952).

## NONLINEAR ANISOTROPIC TERMS IN THE FREE-CARRIER FARADAY ROTATION IN CUBIC SEMICONDUCTORS

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The nonlinearity of the magnetic field dependence of free-carrier Faraday rotation in cubic semiconductors has been investigated theoretically and experimentally in the low-field region  $(\mu B \ll 1$ , where  $\mu$  and B are the mobility and the magnetic field, respectively). The effect arises due to  $B^3$  terms in the expansion of the highfrequency electric current in the presence of the magnetic field. Unlike the first-order term of the Faraday rotation, the nonlinearity is anisotropic for nonspherical energy bands. The effect is particularly interesting in the lowfrequency range ( $\omega \tau \ll 1$ ,  $\omega$  and  $\tau$  being the frequency and the collision time, respectively), where it provides a direct measure of the scattering average  $S = \langle \tau^4 \rangle / (\langle \tau^2 \rangle \langle \tau \rangle^2)$ . Measurements of the nonlinear terms were carried out at room temperature at 35 Gc/sec on a series of  $n$ -type silicon crystals. The observed anisotropy of the effect agreed well with the theoretical analy-