

$$f_V(\Xi^- \Lambda) = (\frac{2}{3})^{1/2}, f_V(\Xi^0 \Sigma^+) = 1, f_V(\Xi^- \Sigma^0) = 1/\sqrt{2};$$

where  $f_V(\Sigma^- n)$  is the vector coupling for  $\Sigma^- \rightarrow n + \text{leptons}$ . From  $\beta$  decay of  $O^{14}$  and  $Al^{26}$  we obtain  $\cos\theta = 0.980$  or  $\sin\theta = 0.20$  [see J. Sakurai, Phys. Rev. Letters 12, 79 (1964)]. From the recent values  $R = (1.07 \pm 0.13) \times 10^{-3}$  for the branching ratio and  $f_A/f_V = 1.03$  for  $\Lambda \beta$  decay (V. G. Lind, T. O. Binford, M. L. Good, and D. Stern, to be published), we find  $|f_V(\Lambda p)| = 1.29 \pm 0.13$  in excellent agreement with  $f_V = -(\frac{2}{3})^{1/2} = -1.22$ . For  $K^0 \rightarrow \pi^- + e^+ + \nu$  we write a matrix element  $(1/\sqrt{2})(G \sin\theta) f(p+q) \bar{\mu} \gamma_\mu \times (1 + \gamma_5) u_\nu$ . The prediction is  $f = 1$ . From data on  $K_2^0$  [D. Luers, I. S. Mitra, W. J. Willis, and S. S. Yamamoto, Phys. Rev. 133, B1276 (1964); Proceedings of the Sienna International Conference on Elementary

Particles (Società Italiana di Fisica, Bologna, Italy, 1963), Vol. 1, p. 23; for the branching ratio we take a weighted average of  $0.56 \pm 0.03$ ], we obtain  $|f| = 0.96 \pm 0.20$ . From  $K^+$  data [B. Roe et al., Phys. Rev. Letters 7, 346 (1961)] using  $\Delta T = \frac{1}{2}$  rule, we obtain instead  $|f| = 1.18 \pm 0.06$ , in apparent disagreement with the predicted value and with the  $K_2^0$  data.

<sup>7</sup>The amplitude for  $\gamma \rightarrow \bar{\Sigma}^0 + \Sigma^0$  can be expressed by charge independence in terms of the other  $\bar{\Sigma}\Sigma$  amplitudes; the (first-class) amplitude for  $\gamma \rightarrow \bar{\Sigma}^0 + \Lambda$  is equal to the amplitude for  $\gamma \rightarrow \bar{\Lambda} + \Sigma^0$  and is expressed as a linear combination of the other amplitudes for neutral baryons [Okubo's relation: S. Okubo, Phys. Letters 4, 14 (1963)].

<sup>8</sup>L. Wolfenstein, to be published.

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#### E R R A T U M

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NUCLEAR SPIN ORDERING IN ADSORBED He<sup>3</sup>.  
M. H. Lambert [Phys. Rev. Letters 12, 67 (1964)].

Further experiments have shown that the observed specific heat anomaly is not due to spin ordering. A complete report is in preparation.