Signatures of Weyl Fermion Annihilation in a Correlated Kagome Magnet

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The manipulation of topological states in quantum matter is an essential pursuit of fundamental physics and next-generation quantum technology. Here we report the magnetic manipulation of Weyl fermions in the kagome spin-orbit semimetal $Co_3Sn_2S_2$, observed by high-resolution photoemission spectroscopy. We demonstrate the exchange collapse of spin-orbit-gapped ferromagnetic Weyl loops into paramagnetic Dirac loops under suppression of the magnetic order. We further observe that topological Fermi arcs disappear in the paramagnetic phase, suggesting the annihilation of exchange-split Weyl points. Our findings indicate that magnetic exchange collapse naturally drives Weyl fermion annihilation, opening new opportunities for engineering topology under correlated order parameters.

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Quantum magnets exhibiting electronic topology are attracting considerable interest for the magnetic manipulation of Weyl and Dirac quasiparticles, as well as their topological surface states [1–7]. To date, spectroscopic signatures of electronic topological ground states have been observed in several magnetic semimetals, comprising magnetic Weyl loops [8]; Weyl points [9–13]; and massive Dirac fermions [14]. In parallel, the magnetic manipulation of Weyl and Dirac fermions has been extensively explored in transport [15–20]. However, direct spectroscopic observation of magnetic control of topology remains challenging. Demonstrating coherent evolution of topological quasiparticles under varying magnetic order, such as the annihilation of Weyl points, offers the possibility to directly verify fundamental notions of topological band theory [2,3,21,22]. Furthermore, novel transport and optical effects are enabled by tuning the relative energies of Weyl loops and points [23–25], controlling their positions relative to the Fermi level [8,26–29] and switching their topological surface states on and off [30,31].

We have investigated magnetic modulation of topological semimetallic states in a range of materials by spectroscopy, including Fe_3Sn_2, Co_2MnGa, PrAlGe, Fe_3GeTe_2, $TbMn_6Sn_6$, and $Co_3Sn_2S_2$ [8,13,32–34]. Some of these materials exhibit high magnetic transition temperatures > 600 K, so that thermal broadening may fundamentally overwhelm magnetic evolution of the Weyl or Dirac state [8,14]. Other systems, such as PrAlGe, exhibit low transition temperatures of ~ 10 K, associated with only small magnetic perturbations to the electronic structure [13]. Even in materials such as Fe_3GeTe_2, with intermediate $T_C = 230$ K, the thermal evolution appears to be dominated by a suppression of quasiparticle lifetime, without significant coherent evolution of the dispersion [35–37]. Using newly available high-quality single crystals combined with state-of-the-art variable-temperature photoemission spectroscopy, we have found that a large and previously overlooked energy shift of a topological spin-orbit gapped Weyl loop occurs in $Co_3Sn_2S_2$ across

$T_C = 176$ K [38–43]. This shift takes place together with a magnetic exchange gap collapse that suggests a ferromagnetic Weyl to paramagnetic Dirac loop transition on raising temperature. This transition is further accompanied by the removal of candidate topological Fermi arc surface states and the annihilation of Weyl points.

Materials with inversion symmetry, mirror symmetry, and ferromagnetism provide a unique platform for a magnetic-topological phase transition. The ferromagnetism produces singly degenerate spin-split bands. In the limit of weak spin-orbit coupling (SOC), mirror symmetry can then give rise to Weyl loops on mirror planes of the bulk Brillouin zone [8,44–46]. A Weyl loop is a closed curve along which the bands are twofold degenerate everywhere; it is characterized by a π Berry phase topological invariant and a linear energy-momentum dispersion everywhere along the loop. If the magnetic order is removed and no spin splitting remains, opposite-spin partner Weyl loops naturally collapse into a Dirac loop, a closed curve along which the bands are fourfold degenerate everywhere [46–48]. Weyl loops under SOC typically gap out, concentrating a loop of Berry curvature in momentum space, leading to a giant anomalous Hall response [8,10], large anomalous Nernst effect [29,49], large optical Hall conductivity [40], and other exotic response. Under SOC, Weyl loops may also leave behind some discrete number of Weyl points. By contrast, Dirac points are generically unstable under inversion and time-reversal symmetry [50], so that Dirac loops under SOC gap out fully. As a result, in this scenario upon magnetic exchange collapse the Weyl points generically annihilate.

$Co_3Sn_2S_2$ crystallizes in space group $R\bar{3}2/m$ (No. 166), with dihedral point group D_{3d}, which includes inversion symmetry and three mirror planes [Fig. 1(a), S3]. The system is ferromagnetic, with Curie temperature $T_C =$ 176 K [51,52]. Keeping in mind the mirror symmetry and ferromagnetic order, we explore our $Co_3Sn_2S_2$ samples by ARPES at 20 K. Measuring with incident photon energy

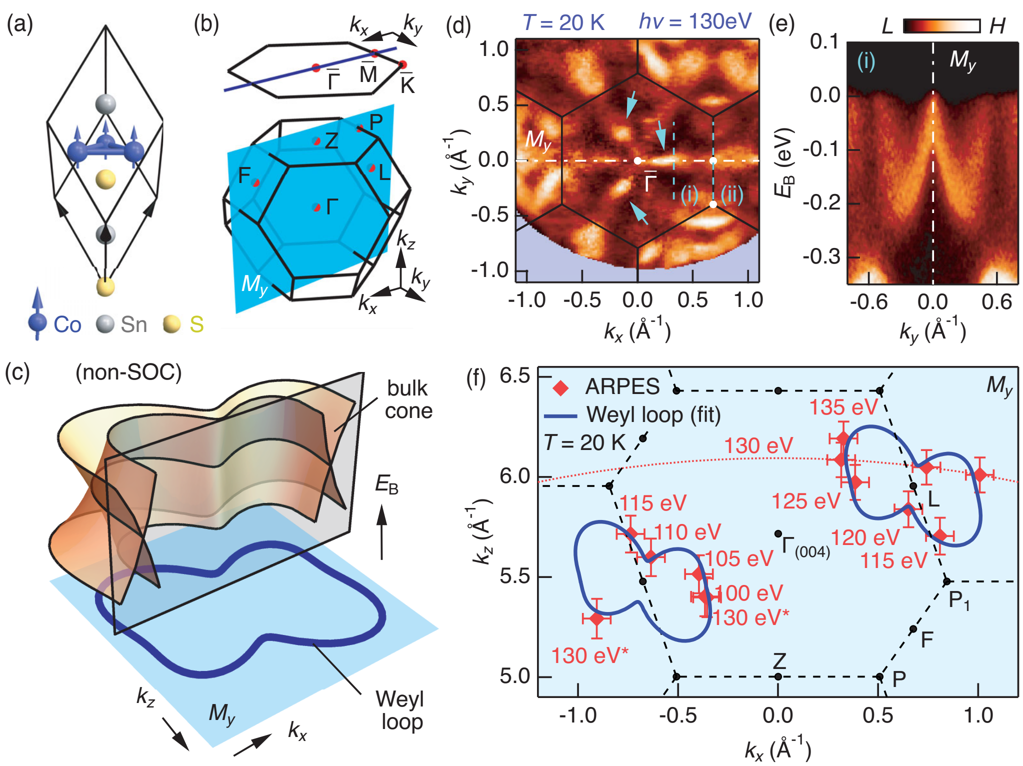

FIG. 1. Topological magnetic Weyl loop. (a) Primitive unit cell of ferromagnetic $Co_3Sn_2S_2$, with mirror symmetry. (b) Bulk and (001) surface Brillouin zones with bulk mirror plane (M_y, cyan) and several high-symmetry points (red). (c) In the absence of spin-orbit coupling (SOC), the combination of mirror symmetry and ferromagnetism generically gives rise to Weyl loops, which live in the mirror planes of the bulk Brillouin zone. A Weyl loop exhibits a ring of band crossings along a closed curve in momentum space (blue loop) with a linear cone dispersion on any energy-momentum slice through the loop. Under SOC, the Weyl loop typically gaps, possibly leaving behind Weyl points. (d) ARPES Fermi surface at $T = 20$ K and photon energy $h\nu = 130$ eV, exhibiting multiple dot features (cyan arrows) on the mirror planes ($\bar{\Gamma} - \bar{M}$). (e) Cone dispersion at the Fermi level on an energy-momentum spectrum through the dot feature (Cut (i)). (f) Collecting cone dispersions for a range of $h\nu$ suggests a loop of band crossings (red diamonds) living in M_y and encircling the bulk L point (Figs. S4, S5). Different $h\nu$ sample different out-of-plane k_z momenta; representative example shown for 130 eV (dotted red curve). The crossing points can be fit by a low-order polar coordinate Fourier decomposition around the L point (blue curve [53]), mapping out the trajectory of the Weyl loop.

$h\nu = 130$ eV, we observe pointlike electronic structures on M_y [cyan arrows, Fig. 1(d); the mirror planes correspond to $\bar{\Gamma} - \bar{M}$]. On cuts along k_y through the pointlike features, we observe cone dispersions straddling the mirror plane M_y [Fig. 1(e)]. On an energy-momentum cut along k_x, within the mirror plane, we again observe conelike dispersions (Fig. S4). The observation of cone dispersions along both k_x and k_y, coming together at point Fermi surfaces, suggests a set of band crossings living in the momentum-space mirror plane. To systematically understand the evolution of the band crossings along the out-of-plane momentum-space direction k_z we acquire analogous datasets at a range of photon energies, from $h\nu = 100$ to 135 eV (Fig. S5). We find that the cones persist in $h\nu$, with crossing points consistently on the M_y plane, but at varying (k_x, k_z) coordinates [red diamonds, Fig. 1(f)]. Taken together, these crossing points appear to form an extended nodal electronic state encircling the L point of the bulk Brillouin zone, suggesting the observation of a bulk loop node in $Co_3Sn_2S_2$. Since the system is ferromagnetic with generically singly degenerate bands, we interpret this loop node as a Weyl loop [Fig. 1(c)]. To extract the complete trajectory of the loop, we fit the ARPES locations of the cone dispersions to a low-order polar Fourier decomposition around the L point of the bulk Brillouin zone [blue loop, Fig. 1(f); see Supplemental Material [53] for fitting parameter values]. In this way we extract the full momentum-space trajectory of the Weyl loop from photoemission data alone.

Next we explore the evolution of the Weyl loop with temperature, focusing on Cut (i). We systematically cycle the temperature of our samples from 20 K to 290 K and back to 20 K, moving across $T_C = 176$ K. On raising the temperature, we observe a dramatic evolution of the Weyl cone on a large energy scale of ~0.1 eV [Figs. 2(a) and 2(b); S8], with the cone appearing to recede above E_F. We next assemble the momentum distribution curves (MDCs) of Cut (i) at E_F for all temperatures [Fig. 2(c)]. Upon cycling the temperature, we observe a prominent and reversible evolution of the Weyl cone across T_C, consistent with a magnetic phase transition. For further insight, we examine additional spectra on Cut (ii), obtained during the course of the same measurement, and we consider a set of deep bands ~0.3 eV below E_F, which are predominantly formed from the same exchange-split Co $3d$ a_{1g} and e_g manifolds as the Weyl loop [Figs. 2(d); S11]. At 20 K, these

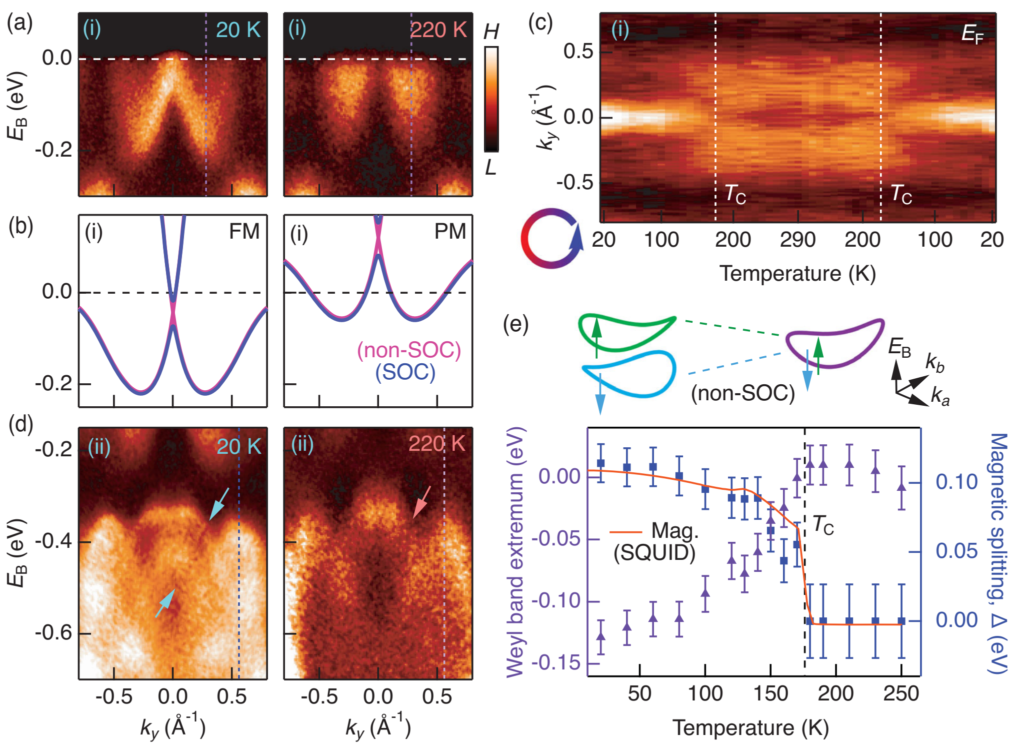

FIG. 2. Magnetic Weyl to Dirac loop collapse. (a) Cut (i) at 20 K and 220 K, with (b) corresponding *ab initio* calculation. Left: calculation in the ferromagnetic state through the Weyl loop, without SOC (magenta) and with SOC (blue). Right: calculation in the nonmagnetic state through the Dirac loop. (c) Momentum distribution curves (MDCs) of Cut (i) at the Fermi level for the full temperature cycle, 20 K → 290 K → 20 K. (d) Cut (ii), defined in Fig. 1(d), exhibiting clear splittings in deeper energy bands at 20 K (left), which collapse at 220 K (right). (e) Energy extremum of the Weyl loop band, extracted from the temperature dependence on Cut (i), obtained by Lorentzian fitting of energy distribution curves [EDCs, dotted lines in (a)]. Also, the magnetic exchange splitting as a function of temperature, obtained from Cut (ii) by Lorentzian fitting to EDCs [dotted lines in (d)] and compared with the magnetization $M(T)$. Cartoon: exchange gap collapse of two opposite-spin Weyl loop partners (blue and green) into a single Dirac loop (purple).

deep valence bands exhibit clear splitting, consistent with the material's ferromagnetic order. Upon raising the temperature, the splitting appears to vanish and these deep bands collapse together, suggesting a paramagnetic state with spin-degenerate bands. By examining the evolution of the deep bands, we circumvent the limitations of the photoemission E_F cutoff and observe direct signatures of a prominent magnetic exchange gap collapse across T_C in $Co_3Sn_2S_2$.

To relate the Weyl loop temperature evolution to the magnetic exchange gap collapse, we consider more carefully the interplay between topology and ferromagnetism. In *ab initio* calculation, in the absence of SOC and in the ferromagnetic state, the Weyl loop arises as a crossing of two spin-majority bands, with a spin-minority partner Weyl loop above the Fermi level [schematic blue and green loops, Figs. 2(e); S6]. In a nonmagnetic *ab initio* calculation, the exchange gap vanishes and these two Weyl loops coincide, forming a spinless loop crossing—a Dirac loop (purple loop). Comparing the *ab initio* calculations with ARPES, we find that the magnetic Weyl and nonmagnetic Dirac nodes exhibit overall agreement with the ferromagnetic and paramagnetic spectra, respectively [magenta traces, Fig. 2(b)]. Note that including SOC in our *ab initio* results does not alter this interpretation, although the expected gap appears in both loop nodes (blue traces). The observation that the loop recedes above E_F on increasing temperature is also consistent with maintenance of charge balance in the spin-degenerate electronic structure, further indicating a paramagnetic Dirac loop. We further reduced the magnetic moment in our samples via nickel (Ni) doping, and again observed a persistent loop node electronic structure despite suppression of the ferromagnetism (Fig. S1). Taken together, our systematic ARPES spectra and *ab initio* calculations suggest that we have observed the collapse of two opposite-spin ferromagnetic Weyl loops into a paramagnetic Dirac loop.

To quantitatively characterize the Weyl loop collapse with temperature, we perform a Lorentzian fit of energy distribution curves (EDCs) through the extremum of the Weyl loop band [Cut (i), dotted line, Fig. 2(a)]. The extracted Weyl band extremum exhibits a clear evolution upward in E_B as the temperature increases, 20 K → 250 K, consistent with exchange gap collapse [Fig. 2(e)]. We further extract the exchange gap $\Delta(T)$ on EDCs through the deep bands [Cut (ii), dotted line, Fig. 2(d)] and compared the resulting $\Delta(T)$ with the magnetization $M(T)$ as measured by a SQUID. For $T < T_C$ we find that the exchange splitting tracks $M(T)$. For $T > T_C$ we no longer observe an exchange splitting within our spectral linewidth, consistent with the absence of magnetization. Remarkably, the observed exchange gap and Weyl band shift are both ~0.12 eV, suggesting a complete collapse of the opposite-spin partner Weyl loops across T_C and the formation of a spin-degenerate paramagnetic Dirac loop.

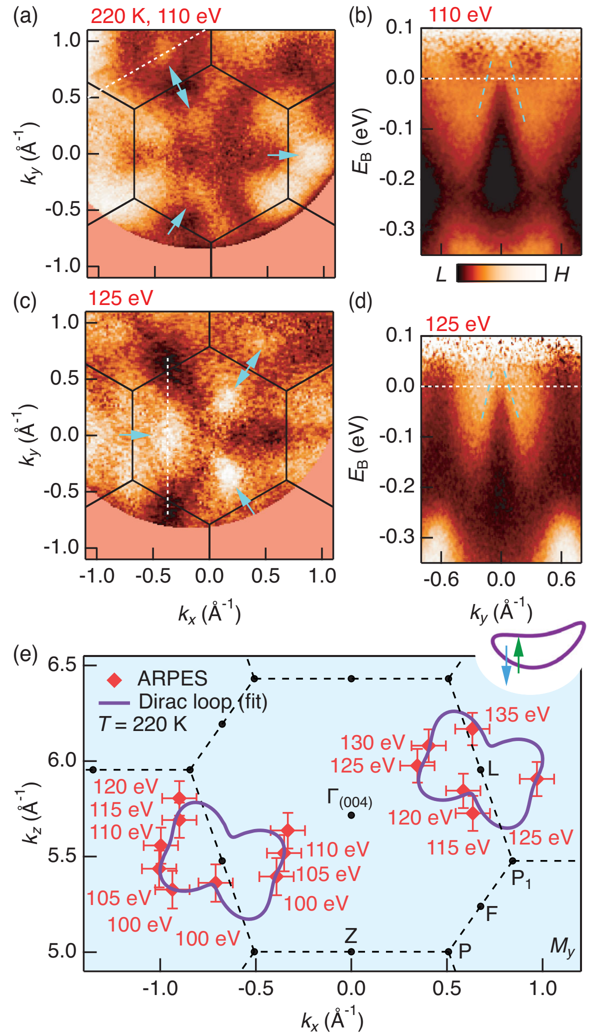

FIG. 3. Paramagnetic Dirac loop. (a) ARPES isoenergy contour slightly above E_F, at $h\nu = 110$ eV, acquired at 220 K, exhibiting pointlike features (cyan arrows) along $\bar{\Gamma}$-$\bar{M}$ (corresponding to M_y and the symmetry-related mirror planes). (b) Energy-momentum cut through the pointlike feature, exhibiting conelike spectral weight [cut location: white line in (a)]. (c),(d) Analogous to (a), (b), at 125 eV. (e) Locations of cones observed for all $h\nu$ (red diamonds, Fig. S7). Cones on symmetry-related mirror planes are plotted all together in a single momentum-space mirror plane M_y. Data points fit to a low-order polar coordinate Fourier decomposition around L (purple curves), mapping out the trajectory of the Dirac loop.

To further explore the paramagnetic Dirac loop we park our apparatus at 220 K, well into the paramagnetic phase. At a range of $h\nu$ we observe characteristic pointlike isoenergy contours on M_y and related mirror planes [Figs. 3(a) and 3(c); S7]. Energy-momentum spectra through these pointlike contours further exhibit conelike spectral weight straddling M_y, indicative of Dirac loop cone dispersions above E_F [Figs. 3(b) and 3(d)]. The presence of multiple cone features straddling M_y at a range of $h\nu$ again suggests an extended nodal electronic structure

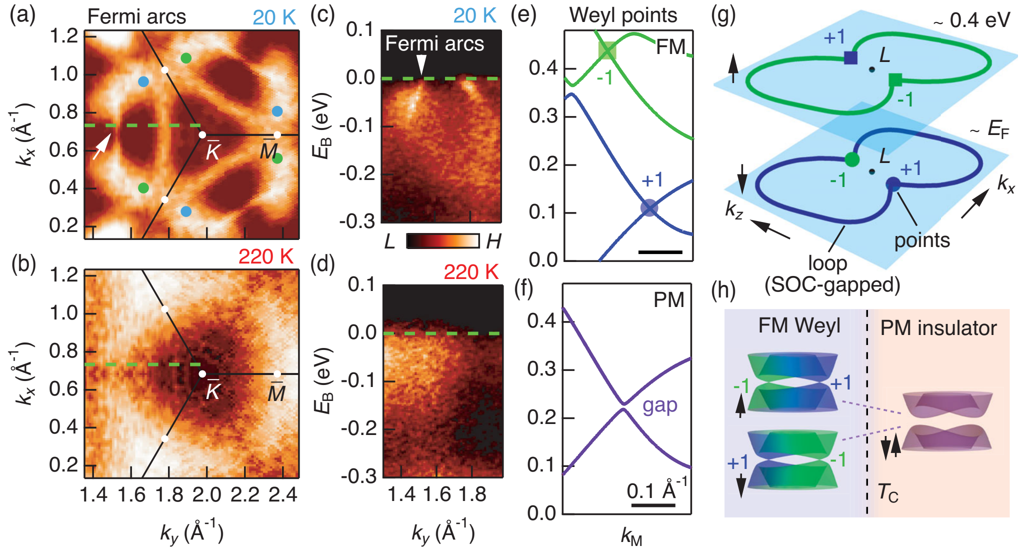

FIG. 4. Evidence for Fermi arc and Weyl point annihilation. (a) Fermi surface acquired at $T = 20$ K, $h\nu = 130$ eV, exhibiting candidate Fermi arc surface states (white arrow) near the Weyl point locations, as predicted by *ab initio* calculations (blue, green circle; Fig. S9). Data symmetrized for clarity. (b) Analogous Fermi surface at 220 K, with no signature of Fermi arcs. (c) Energy-momentum cut through the candidate Fermi arcs at 20 K [green dotted line, (a)]. (d) Analogous cut at 220 K. (e) Calculation in the ferromagnetic state, in the presence of SOC, slicing through a pair of exchange-split Weyl points of opposite chirality. (f) Analogous calculation, nonmagnetic state. The two partner Weyl points annihilate, opening a gap of 12 meV. (g) Schematic of the ferromagnetic Weyl loop (non-SOC) and Weyl point (SOC) configuration. (h) Schematic phase diagram: spin-orbit gapped ferromagnetic Weyl loops collapse to a paramagnetic Dirac loop across T_C. Concurrently, the exchange-split Weyl points annihilate.

confined to the mirror plane. Since we are in the paramagnetic phase with generically spin-degenerate bands, we interpret these candidate band crossings as fourfold degenerate, forming a Dirac loop. By analogy with our analysis in the ferromagnetic phase, we again systematically collect the locations of all cone features observed in the paramagnetic phase for $h\nu$ from 100 to 135 eV and experimentally extract the full momentum-space trajectory of the Dirac loop [red diamonds, Fig. 3(e)]. *Ab initio* calculations of the Weyl and Dirac loops in the ferromagnetic and nonmagnetic states also agree with the experimentally observed trajectories (Fig. S2). A loop node electronic structure persisting into the paramagnetic phase of $Co_3Sn_2S_2$ again suggests the observation of a ferromagnetic Weyl to paramagnetic Dirac loop collapse.

We next consider the fine structure of the Weyl loop collapse associated with SOC. In $Co_3Sn_2S_2$, *ab initio* calculations along with ARPES and STM investigations suggest that each Weyl loop under SOC produces two Weyl points above the Fermi level, with signatures of topological Fermi arc surface states extending below the Fermi level. [9–12,38–41]. Our observation of a Weyl to Dirac loop transition naturally motivates investigation of Fermi arc and Weyl point annihilation across T_C. At 20 K, we observe sharp arc-shaped states near the expected Weyl points, consistent with topological Fermi arcs in *ab initio* calculation [Figs. 4(a) and 4(c); S9]. At 220 K these Fermi arcs vanish, leaving behind bulk pockets broadly consistent with the low temperature spectra [Figs. 4(b) and 4(d)]. The disappearance of the Fermi arcs above T_C provides evidence for the annihilation of Weyl points in the paramagnetic phase. To further characterize this annihilation, we consider the *ab initio* band structure under ferromagnetic order on a momentum-space path connecting a pair of exchange-split Weyl points [Fig. 4(e)]. Upon exchange collapse, these partner Weyl points come together and annihilate, opening a gap [Fig. 4(f)].

Our systematic variable-temperature ARPES experiments suggest that pairs of ferromagnetic Weyl loops collapse into paramagnetic Dirac loops across T_C in $Co_3Sn_2S_2$ [Figs. 4(g) and 4(h)]. Taken together with *ab initio* calculations, our results additionally provide evidence for the annihilation of Fermi arcs and Weyl points concomitant with this transition. Our findings suggest a general mechanism for Weyl fermion annihilation, where the annihilation is driven by magnetic exchange gap collapse and takes place predominantly along the energy axis, rather than in momentum [3,21,62]. This novel mechanism should occur naturally in many quantum magnets and motivates exploration of the rich topological evolution associated with the onset of magnetic order, including novel magnetic topological insulators. Such interplay between magnetism and topology may also pave the way to magnetic design of correlated topological states with exotic transport and optical response.

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