

Electronic Raman Scattering in Twistrionic Few-Layer Graphene

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We study electronic contribution to the Raman scattering signals of two-, three- and four-layer graphene with layers at one of the interfaces twisted by a small angle with respect to each other. We find that the Raman spectra of these systems feature two peaks produced by van Hove singularities in moiré minibands of twistrionic graphene, one related to direct hybridization of the Dirac states, and the other resulting from band folding caused by moiré superlattice. The positions of both peaks strongly depend on the twist angle, so that their detection can be used for noninvasive characterization of the twist, even in hBN-encapsulated structures.

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Twisted bilayer graphene is a van der Waals heterostructure where the relative twist of constituent atomic planes alters electronic properties of the material [1,2]. A small-angle twist in a bilayer produces a long-period moiré pattern that generates minibands for electrons with a small moiré Brillouin zone (mBZ). The minibands and gaps between them strongly depend on twist angle θ leading to Mott insulator states and superconductivity for a magic angle, $\theta \approx 1.1^\circ$ [1–11], where the lowest miniband appears to be almost flat. Narrow minibands also appear in (1 + 2) [12,13] and (1 + 1 + 1) [14] trilayers as well as in (2 + 2) tetralayers [15–21]. Theoretical studies also predict the appearance of correlated and topological states in various “twistrionic” graphene stacks [22–28], highlighting the need to expand the toolbox of fast and noninvasive methods for measuring twist angles in such structures.

Over the years, Raman spectroscopy emerged as a powerful tool for the characterization of graphene. Raman scattering with phonons provides information about defects, doping, strain, and the number of layers in the film [29]. Twist of graphene layers was found to lead to the resonant enhancement as well as variation of the width and position of various Raman-active phonon modes [30–37], however, with a limited accuracy in determining the twist angle. Here, we demonstrate that Raman spectroscopy of

the interband electronic excitations (ERS) [38–50] can be used to detect twist-angle-dependent features in the electronic spectrum of twistrionic graphene. We study the electronic minibands and electronic contributions to the Raman spectra for few-layer graphene stacks with one of the interfaces between the layers twisted by a small angle, $\theta < 2^\circ$, Fig. 1(a). We calculate ERS spectra of twisted bilayers (1 + 1), trilayers (1 + 2), and tetralayers

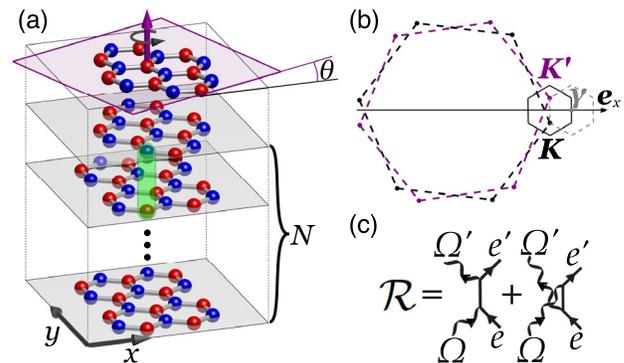


FIG. 1. (a) Sketch of (1 + N) twistrionic graphene. The red and blue balls correspond to the two different sublattice sites in each graphene layer and green ellipse represents a dimer bonding leading to direct interlayer coupling. (b) Brillouin zones of the N -layer stack (dashed black line) and the top graphene (dashed purple line) with corners \mathbf{K} (bottom layer) and \mathbf{K}' (top), respectively, as well as the effective moiré Brillouin zone, shown both centered around the valleys (solid black line) as the preferred choice in this Letter, and with valleys in its corners (dashed gray), centered at γ . (c) Feynman diagrams for the scattering amplitudes \mathcal{R} contributing to the electronic Raman features discussed in this Letter.

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$[(1+3)$ and $(2+2)]$ and show that these are formed by transitions from the n th valence to the n th conduction moiré superlattice (mSL) miniband and feature two spectral peaks. One, at a lower Raman shift, is caused by the resonant hybridization of electronic states of the two few-layer graphene crystals separated by the twisted interface forming the lowest-energy minibands [51,52]. Another higher-energy peak is due to the anticrossing of bands, backfolded by mSL. Both peaks are related to van Hove singularities in the mSL minibands and involve electronic excitations different from those responsible for the optical absorption [53,54]. We trace the peaks positions as a function of the twist angle and estimate their quantum efficiency, $I \sim 10^{-11}$, to be in the measurable range [38–45].

To model twistrionic graphene, we use a hybrid $k \cdot p$ theory-tight-binding model, where we describe electrons' states in each flake using the $k \cdot p$ expansion around $\pm\mathbf{K}$ and $\pm\mathbf{K}'$ Brillouin zone corners of the bottom and anticlockwise rotated (by angle θ) top crystal, respectively [see Fig. 1(b)], and the interlayer hybridization using tunneling Hamiltonian [52,55],

$$\hat{H} = \begin{array}{c|ccc|ccc} \dots & \vdots & \vdots & \vdots & \vdots & \vdots & \dots \\ \dots & \hat{H}_+ & \hat{T}_t & 0 & 0 & 0 & \dots \\ \dots & \hat{T}_t^\dagger & \hat{H}_+ & \hat{T} & 0 & 0 & \dots \\ \hline \dots & 0 & \hat{T}^\dagger & \hat{H}_- & \hat{T}_b & 0 & \dots \\ \dots & 0 & 0 & \hat{T}_b^\dagger & \hat{H}_- & \hat{T}_{b'} & \dots \\ \dots & 0 & 0 & 0 & \hat{T}_{b'}^\dagger & \hat{H}_- & \dots \\ \dots & \vdots & \vdots & \vdots & \vdots & \vdots & \dots \end{array}, \quad (1)$$

$$\hat{H}_\pm = v\sigma_\xi \cdot \left(\mathbf{p} \mp \xi\hbar\frac{\theta}{2}K\mathbf{e}_y \right).$$

Here, $\sigma_\xi = (\xi\sigma_x, \sigma_y)$ and $\sigma_x, \sigma_y, \sigma_z$ are sublattice Pauli matrices for electrons within each monolayer, and the vertical and horizontal lines mark the twisted interface. Momentum $\mathbf{p} = (p_x, p_y)$ is measured from the center of the valley $\xi\mathbf{K}$ in the bottom flake (with valley index $\xi = \pm 1$), $K = |\mathbf{K}|$, \mathbf{e}_y is a unit vector along wave vector y axis, and v is the monolayer Dirac velocity. To describe the $(MY+NX)$ twistrionic system, where M and N are the thicknesses of the stack above and below the twisted interface and X and Y describe the stacking of the corresponding crystal [$X=B$ for Bernal and $X=R$ for rhombohedral trilayer, $X=AB$ and $Y=AB$ or $Y=BA$ for $(2+2)$ structures], one needs to truncate \hat{H} at M blocks in the top left and N in the bottom right parts. Coupling across the twisted interface is described by [52,55]

$$\hat{T} = t(\hat{T}_0 + \hat{T}_1 e^{i\xi\mathbf{G}_1 \cdot \mathbf{r}} + \hat{T}_2 e^{i\xi\mathbf{G}_2 \cdot \mathbf{r}}),$$

$$\hat{T}_n = 1 + \cos\left(\frac{2\pi n}{3}\right)\sigma_x + \xi \sin\left(\frac{2\pi n}{3}\right)\sigma_y, \quad (2)$$

where $\mathbf{G}_n \approx [(-1)^n \frac{1}{2}, (\sqrt{3}/2)]\sqrt{3}\theta K$ and $t \approx 110$ meV [55]. Coupling \hat{T} is responsible for the mSL effects and defines the hexagonal mBZ, chosen as in Fig. 1(b). Interlayer couplings across the nontwisted interfaces are set using

$$\hat{T}\left(\pm\frac{\theta}{2}\right) = \begin{bmatrix} -v_4\hat{\pi}^\dagger(\pm\frac{\theta}{2}) & v_3\hat{\pi}(\pm\frac{\theta}{2}) \\ \gamma_1 & -v_4\hat{\pi}^\dagger(\pm\frac{\theta}{2}) \end{bmatrix},$$

$$\hat{\pi}\left(\pm\frac{\theta}{2}\right) = \begin{bmatrix} \xi p_x + ip_y \mp i\xi\hbar\frac{\theta}{2}K \end{bmatrix},$$

where $\gamma_1 = 0.39$ eV, $v_3 \approx 0.11v$, and $v_4 \approx 0.015v$ [56]. For a $(1+2)$ trilayer, $\hat{T}_b = \hat{T}[-(\theta/2)]$. For a $(1+3R)$ stack, $\hat{T}_b = \hat{T}_{b'} = \hat{T}(-\theta/2)$. For a $(1+3B)$ tetralayer, $\hat{T}_b = \hat{T}(-\theta/2)$ and $\hat{T}_{b'} = \hat{T}^\dagger(-\theta/2)$. For a $(2AB+2AB)$ structure, $\hat{T}_t = \hat{T}(\theta/2)$, $\hat{T}_b = \hat{T}(-\theta/2)$, while for $(2AB+2BA)$, $\hat{T}_b = \hat{T}(-\theta/2)$, $\hat{T}_t^\dagger = \hat{T}(\theta/2)$.

In the electronic Raman scattering, photon of energy Ω carrying vector potential $\mathbf{A} = \sqrt{(\hbar/2\epsilon_0\Omega)}[l e^{i(\mathbf{q}\cdot\mathbf{r}-\Omega t)/\hbar} \hat{b}_{\mathbf{q},l} + \text{H.c.}]$ ($\hat{b}_{\mathbf{q},l}$ annihilates a photon with momentum \mathbf{q} and polarization l), arrives at the sample (here, we assume normal incidence of light). This photon scatters to another one with energy $\tilde{\Omega} = \Omega - \omega$, momentum $\tilde{\mathbf{q}}$, and polarization \tilde{l} , leaving behind an electron-hole pair with energy ω . In contrast to classical plasmas, where the amplitude of such process is controlled by contact interaction, in graphene, the dominating contribution comes from a two-step process, described by the Feynman diagrams shown in Fig. 1(c). It corresponds to absorption (emission) of a photon with energy Ω ($\tilde{\Omega}$) transferring an electron with momentum \mathbf{p} from an occupied state in the valence band into a virtual intermediate state (energy is not conserved at this stage), followed by emission (absorption) of a second photon with energy $\tilde{\Omega}$ (Ω) [38,39]. The amplitude of this process is [45]

$$\mathcal{R} = i \frac{(e\hbar v)^2}{\epsilon_0 \Omega^2} (\mathbf{l} \times \tilde{\mathbf{l}}^*)_z \hat{I}_{M+N} \otimes \sigma_z, \quad (3)$$

where \hat{I}_{M+N} is a $(M+N) \times (M+N)$ unit matrix. The main contribution to the Raman signal comes from $n^- \rightarrow n^+$ miniband transitions, where n^s denotes the n th miniband on the valence ($s = -1$) or conduction ($s = 1$) side. The structure of the two-photon coupling vertex in Eq. (3) [38] suggests that, for circularly polarized photons, both the incoming and outgoing light have the same polarization. In turn, for linearly

polarized light, the incoming and outgoing light carries perpendicular polarization (so-called linear cross polarization). To compare, the phonon G -line is observed in both parallel and perpendicular linear polarization geometries whereas for circular polarization the incoming and outgoing photons have opposite polarization [42,43,47–49,57]. The overall line shape $g(\omega)$ of inelastic photon scattering with Raman shift ω is given by

$$g(\omega) = \frac{1}{c} \int \frac{d\tilde{q} w(\omega)}{(2\pi\hbar)^3} \delta(\tilde{\Omega} - c|\tilde{q}|) \approx \frac{\Omega^2}{(2\pi\hbar)^3 c^4} w(\omega),$$

$$w(\omega) = \frac{2}{\pi\hbar^3} \sum_{n^s, m^{s'}} \int d\mathbf{p} |\langle \mathbf{p}, m^{s'} | \mathcal{R} | \mathbf{p}, n^s \rangle|^2 \times f_{p, n^s} (1 - f_{p, m^{s'}}) \delta(\epsilon_{p, m^{s'}} - \epsilon_{p, n^s} - \omega), \quad (4)$$

where f_{p, n^s} is the occupation factor of a state $|\mathbf{p}, n^s\rangle$ with momentum \mathbf{p} and energy ϵ_{p, n^s} in the band n^s .

In Fig. 2 we present the ERS intensity map for a $(1+1)$ twisted bilayer graphene with $0.8^\circ < \theta < 2^\circ$, neglecting its atomic reconstruction [58]. Two bright features stand out in this plot, corresponding to transitions between the minibands $1^- \rightarrow 1^+$ and $2^- \rightarrow 2^+$. For two selected cuts at $\theta = 1.8^\circ$ and $\theta = 1.1^\circ$, the peaks

come from transitions indicated by arrows in the corresponding miniband plots. The first peak, $1^- \rightarrow 1^+$, is due to transitions from (or to) flat regions of electronic dispersion resulting from hybridization of the Dirac cones of the two layers (the Dirac points can be identified in the dispersion for $\theta = 1.8^\circ$ as touching points of the minibands shown in red), which are responsible for a van Hove singularity (vHs) in the density of states [51,52] for larger twist angles. Note that due to the chiral nature of graphene electrons, the vHs is not positioned on the line connecting the cones but is shifted off it in the opposite directions in the conduction and valence bands (as indicated by black arrows in Fig. 2 for $\theta = 1.8^\circ$). The inset below the $\theta = 1.8^\circ$ miniband structure highlights in yellow parts of the mBZ that contribute to the $1^- \rightarrow 1^+$ peak. As the twist angle is decreased, this ERS peak moves to lower energies. The second peak, $2^- \rightarrow 2^+$, is due to the transitions between the flat regions of the second valence and conduction minibands, indicated by orange arrows in the $\theta = 1.1^\circ$ miniband plot. Its intensity comes from the mBZ section painted in yellow in the leftmost inset below the ERS map. The other two insets show the real-space distribution of the saddle point states across the moiré supercell in the top and bottom monolayers.

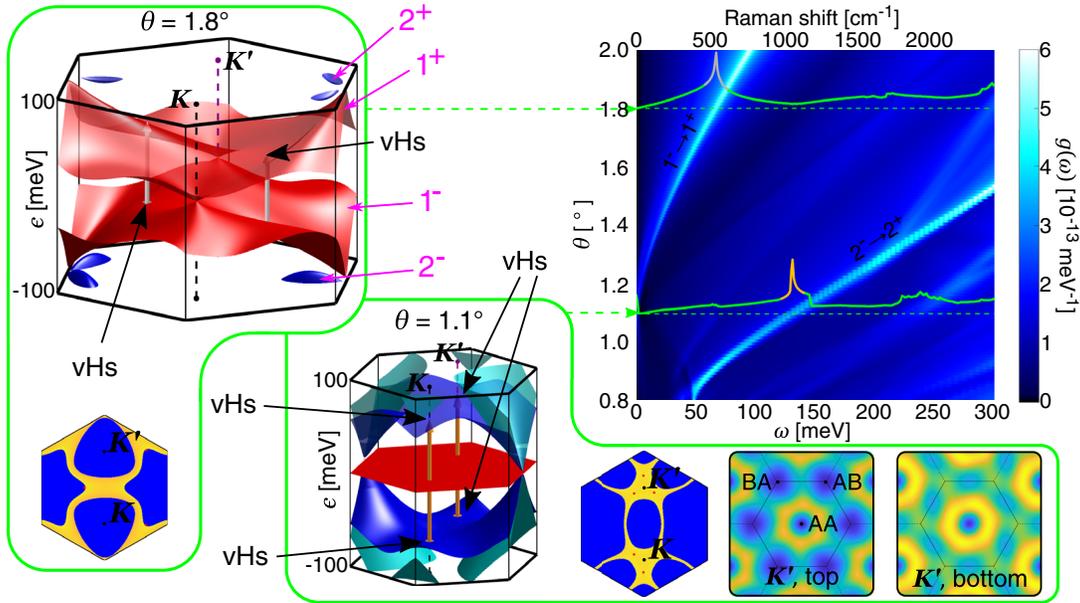


FIG. 2. Left: Electronic miniband structures of twisted bilayer graphene for $\theta = 1.8^\circ$ and $\theta = 1.1^\circ$ across the corresponding mBZ (solid black hexagons). The dashed lines indicate the positions of the valleys \mathbf{K} (bottom layer) and \mathbf{K}' (top) within the mBZ. The pink arrows identify the minibands and the black arrows point to saddle points which give rise to van Hove singularities (vHs). Right: ERS intensity map for $0.8^\circ < \theta < 2^\circ$. The green solid lines show cuts for the angles $\theta = 1.8^\circ$ and $\theta = 1.1^\circ$, indicated with the green dashed lines. The gray and orange peaks mark spectral features corresponding to excitations indicated with vertical arrows of the same color in the miniband dispersions which involve dispersion saddle points. The hexagonal insets next to band structures show in yellow mBZ regions which contribute to the $1^- \rightarrow 1^+$ ERS peak for $\theta = 1.8^\circ$ and $2^- \rightarrow 2^+$ peak for $\theta = 1.1^\circ$; red dots indicate positions of dispersion saddle points on the valence side. The additional two insets for $\theta = 1.1^\circ$ depict real space distribution of the saddle points wave function in the top or bottom layers; black solid lines mark boundaries of the moiré supercell and the letters indicate local stacking.

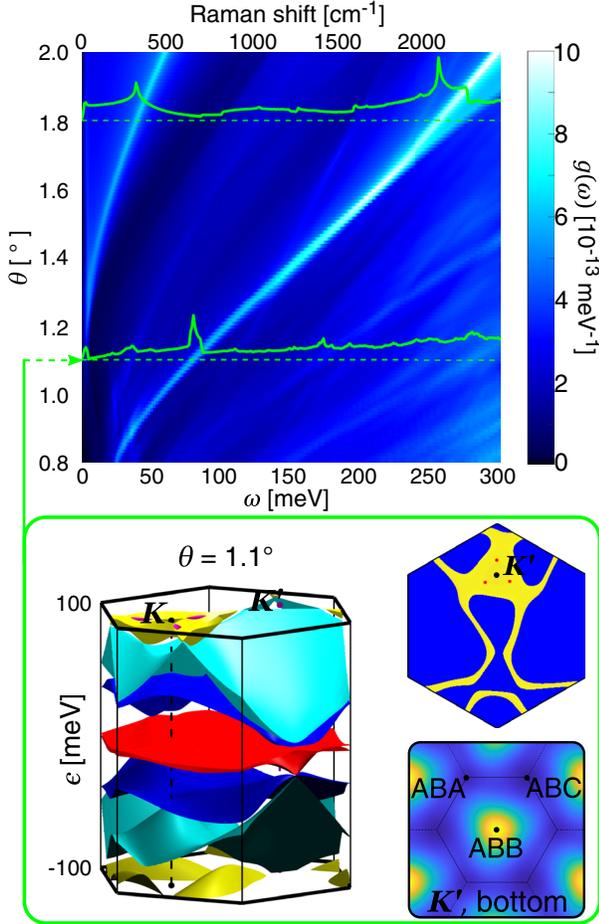


FIG. 3. Top: ERS intensity map for (1 + 2) graphene. The green dashed lines indicate twist angles $\theta = 1.8^\circ$ and $\theta = 1.1^\circ$ and the green solid lines show the ERS spectra for these angles. Bottom: Miniband structure for the (1 + 2) stack for $\theta = 1.1^\circ$. The top inset shows in yellow the mBZ regions which contribute to the $2^- \rightarrow 2^+$ ERS peak; red dots indicate positions of dispersion saddle points on the valence side. The saddle points wave function is predominantly located in the bottom layer and its real space distribution is presented in the bottom inset; black solid lines mark boundaries of the moiré supercell and the letters indicate local stacking (from the bottom layer to the top).

In Fig. 3, we show the ERS intensity map for (1 + 2) twistronic graphene and an exemplary miniband structure for $\theta = 1.1^\circ$. Similarly to (1 + 1) graphene, the dominant contributions come from the $1^- \rightarrow 1^+$ and $2^- \rightarrow 2^+$ electronic transitions. The two peaks also have the same origins: the first one is due to direct hybridization of the monolayer and bilayer states while the second is formed by states backfolded by moiré superlattice. However, for a given twist angle, the peaks appear at lower Raman shifts than in the (1 + 1) structure. This is because the unperturbed dispersion

in bilayer is flatter than in a monolayer—as a consequence, the anticrossings vHs form at lower energies than in a (1 + 1) stack.

In Figs. 4(a) and 4(b), we plot ERS maps of two monolayer-on-trilayer structures, (1 + 3B) and (1 + 3R). Bernal-stacked trilayer graphene hosts both a bilayer- and monolayerlike low-energy bands [59]. Both of these hybridize with the states of the top monolayer to form the first miniband and contribute to the $1^- \rightarrow 1^+$ ERS peak. The next peak, marked as (1 + 2') in Fig. 4(a), is due to the anticrossing of backfolded bilayerlike and top monolayer bands. Another peak, (1 + 1'), is due to an anticrossing of the backfolded top monolayer band and monolayerlike band of the trilayer. In contrast, rhombohedral trilayers only host one low-energy band, with a flat dispersion in the vicinity of the valley center [60]. This low-energy band is localized on the top and bottom surfaces of the crystal (representing surface states generic for rhombohedral graphitic films [61–63]) and is significantly affected by mSL, leading to a pair of clear spectral features, Fig. 4(b), as in (1 + 1) and (1 + 2) twistronic graphenes (in the Supplemental Material, we describe a simple model that can be used to determine approximate positions of these ERS peaks [64]).

In Figs. 4(c) and 4(d), we present ERS spectra for two structurally inequivalent double-bilayer stacks, (2AB + 2BA) and (2AB + 2AB). The details of their lattice structure and exemplary miniband dispersions are shown in the Supplemental Material [64]. The miniband spectra of AB/AB and AB/BA tetralayers are nearly identical, leading to the $1^- \rightarrow 1^+$ and $2^- \rightarrow 2^+$ ERS peaks at the same energies. However, due to the difference in the wave functions, the intensity of the $1^- \rightarrow 1^+$ is much higher for (2AB + 2AB) than for (2AB + 2BA) tetralayers (see Supplemental Material [64]).

Overall, electronic Raman scattering spectra of twistronic graphene contain characteristic features related to the van Hove singularities of moiré superlattice minibands, which depend on the twist angle between the layers. Currently, accurate determination of the twist angle, especially in the small-angle regime, requires time-consuming microscopic investigations of the moiré pattern, or magnetotransport measurements at cryogenic temperatures. We suggest that, based on our results, calibration of the positions of the Raman features in the structures with known θ would allow to identify the orientations of the component crystals in other samples, enabling a noninvasive method for measuring the twist angle, even in structures encapsulated with other materials, where the graphene or graphene moiré pattern is

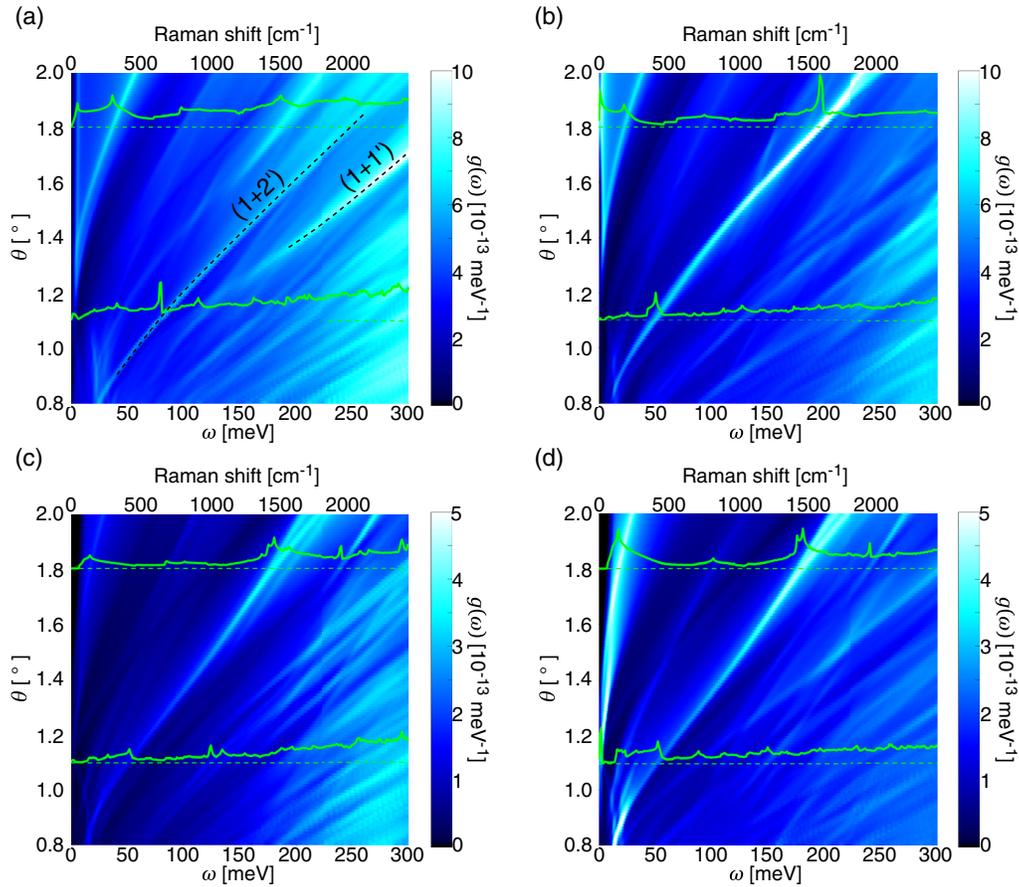


FIG. 4. Comparison of ERS intensity maps for (a) $(1 + 3B)$, (b) $(1 + 3R)$, (c) $(2AB + 2BA)$, and (d) $(2AB + 2AB)$ tetralayers. Green solid lines show the Raman spectra for twist angles $\theta = 1.8^\circ$ and $\theta = 1.1^\circ$.

not directly accessible for tunnelling spectroscopy studies [1,2,5,6,9–21].

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