

β Decays as Sensitive Probes of Lepton Flavor UniversalityAndreas Crivellin^{1,2} and Martin Hoferichter^{3,4}¹Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland²Physik-Institut, Universität Zürich, Winterthurerstrasse 190, CH-8057 Zürich, Switzerland³Albert Einstein Center for Fundamental Physics, Institute for Theoretical Physics, University of Bern, Sidlerstrasse 5, CH-3012 Bern, Switzerland⁴Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA

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Nuclear β decays as well as the decay of the neutron are well-established low-energy probes of physics beyond the standard model (SM). In particular, with the axial-vector coupling of the nucleon g_A determined from lattice QCD, the comparison between experiment and SM prediction is commonly used to derive constraints on right-handed currents. Further, in addition to the CKM element V_{us} from kaon decays, V_{ud} from β decays is a critical input for the test of CKM unitarity. Here, we point out that the available information on β decays can be reinterpreted as a stringent test of lepton flavor universality (LFU). In fact, we find that the ratio of V_{us} from kaon decays over V_{us} from β decays (assuming CKM unitarity) is extremely sensitive to LFU violation (LFUV) in W - μ - ν couplings thanks to a CKM enhancement by $(V_{ud}/V_{us})^2 \sim 20$. From this perspective, recent hints for the violation of CKM unitarity can be viewed as further evidence for LFUV, fitting into the existing picture exhibited by semileptonic B decays and the anomalous magnetic moments of muon and electron. Finally, we comment on the future sensitivity that can be reached with this LFU violating observable and discuss complementary probes of LFU that may reach a similar level of precision, such as $\Gamma(\pi \rightarrow \mu\nu)/\Gamma(\pi \rightarrow e\nu)$ at the PEN and PiENu experiments or even direct measurements of $W \rightarrow \mu\nu$ at an FCC-ee.

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Introduction.—Within the SM of particle physics the masses and mixing angles of quarks have a common origin: their Yukawa couplings to the Higgs boson. In the physical basis with diagonal mass matrices, the misalignment between the up- and down-quark Yukawa couplings is parametrized by the unitary Cabibbo-Kobayashi-Maskawa (CKM) matrix [1,2]. Therefore, the elements of the CKM matrix are fundamental quantities of the SM and their determination is of utmost theoretical and experimental importance [3].

Superaligned β decays—long-lived nuclear $0^+ \rightarrow 0^+$ transitions—are the primary source of information on the CKM matrix element V_{ud} [4–6]. However, additional information can be obtained from neutron decay, whose lifetime and decay asymmetry parameter together determine the axial-vector coupling of the nucleon g_A as well as V_{ud} , with a sensitivity close to, but not yet competitive with, the superallowed nuclear decays [7]. The present situation is illustrated in Fig. 1, based on recent results for the neutron lifetime $\tau_n = 877.7(7)^{(+0.4)}_{(-0.2)} \text{ s}$ [8] and the

asymmetry parameter $\lambda = g_A/g_V = -1.27641(56)$ [9]. The values of V_{ud} , both from superallowed β decays and the neutron lifetime, depend crucially on the applied

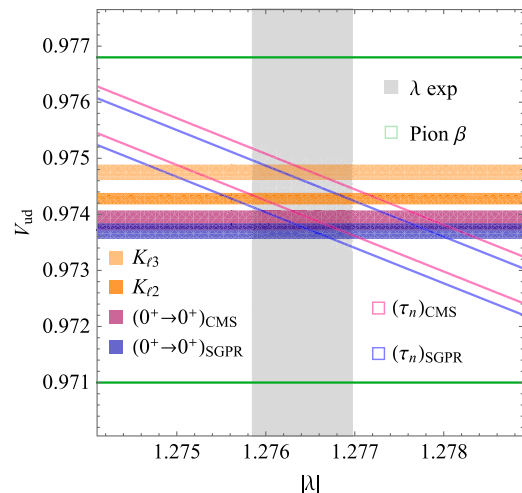


FIG. 1. Constraints on V_{ud} and $|\lambda|$ from superallowed β decays [5,6], neutron lifetime [8], and asymmetry parameter [9], for two sets of radiative corrections [11,14]. We also include the constraints from pion β decay [15,16] and as well as from V_{us} as determined from $K_{\ell 2}$ and $K_{\ell 3}$ decays when assuming CKM unitarity. On this scale, the determination of g_A from lattice QCD [20] does not yet provide a competitive constraint.

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radiative corrections [10–14]. Here, we will consider two sets of corrections “SGPR” [11] and “CMS” [14] to illustrate the spread. The agreement of the bands from $0^+ \rightarrow 0^+$ transitions and neutron decays improves if the value for the neutron lifetime is moved towards its PDG value $\tau_n = 879.4(6)$ s [3], highlighting the importance of an accurate τ_n measurement for the V_{ud} determination. The figure also shows the constraint from pion β decay $\pi^\pm \rightarrow \pi^0 e^\pm \nu_e$ [15,16] as well as the preferred values for V_{ud} deduced from $K_{\ell 2}$ ($K \rightarrow \ell \nu$, $\ell = \mu, e$) and $K_{\ell 3}$ ($K \rightarrow \pi \ell \nu$) decays under the assumption of CKM unitarity. The observed discrepancy between the latter and the direct determinations of V_{ud} has been interpreted as a possible sign for the (apparent) violation of CKM unitarity and triggered recent interest in potential explanations beyond the SM (BSM) [17,18]. However, inducing a sizable violation of CKM unitarity is in general difficult due to the strong bounds from flavor-changing neutral currents, such as kaon mixing, see, e.g., Ref. [19].

In addition, the determination of g_A from neutron decay, once compared to calculations in lattice QCD [20], allows one to constrain the size of right-handed ud currents [21], albeit not yet at a level competitive with the experimental determination of g_A . Still, the comparison shows that both determinations of g_A are compatible, demonstrating that within current uncertainties there is no evidence for right-handed contributions. In particular, their effect does not suffice to remove the tension with CKM unitarity, although a small reduction in significance is possible given that $K_{\ell 2}$ and $K_{\ell 3}$ decays are affected in the opposite way, so that their determinations of V_{us} can be brought into better agreement [22]. This situation is reminiscent of previous hints for right-handed currents in semileptonic B decays [23,24] that disappeared with updated measurements and theory predictions, to the effect that currently an explanation of the discrepancies between the inclusive and exclusive determinations of V_{cb} and V_{ub} in terms of right-handed currents is disfavored [25,26].

On the other hand, experiments have accumulated intriguing hints for the violation of LFU within recent years. In particular, the measurements of the ratios $R(D^{(*)})$ [27–29] and $R(K^{(*)})$ [30,31] deviate from the SM expectation of LFU by more than 3σ [32–36] and 4σ [37–40], respectively. In addition, the anomalous magnetic moments $(g-2)_\ell$ of charged leptons also measures the violation as they vanish in the massless limit. Here, there is the longstanding discrepancy in $(g-2)_\mu$ of about 3.7σ [41–66] and a recently emerging deviation in the electron case of 2.5σ , interestingly, with the opposite sign [67–73]. Furthermore, it has been shown that LFU violating neutrino interactions with SM gauge bosons give an excellent fit of electroweak data, including LFU tests [74].

In this context, it seems natural to consider the discrepancies between the different determinations of V_{ud} (or, equivalently, V_{us} under the assumption of CKM unitarity), in particular the direct determination from β decays, not as a sign of right-handed currents or as a violation of CKM unitarity, but rather as a sign of LFUV [74]. In fact, the most precise determination of V_{ud} from $K \rightarrow \mu \nu$ involves muons, while β decays can only have electrons in the final states. Therefore, we offer the novel perspective to use β decays to search for LFUV and propose to study the observable

$$R(V_{us}) = \frac{V_{us}^{K_{\mu 2}}}{V_{us}^{\beta}} \quad (1)$$

as the corresponding measure. As we will show below, this observable proves to be extremely sensitive to LFUV in the charged current (in particular in the muon sector) and even complements the picture described above, as it deviates by 2 – 3σ from unity. Including $K \rightarrow \pi \ell \nu$ and V_{us} from τ decays into the analysis would even increase the tension towards the 4σ level [17,18,22].

β decays and LFUV.—We are interested in testing LFU of the charged current, i.e., of W - ℓ - ν couplings, which we parametrize in terms of small corrections ε_{ij} according to

$$\mathcal{L} \supset -i \frac{g_2}{\sqrt{2}} \bar{\ell}_i \gamma^\mu P_L \nu_j W_\mu (\delta_{ij} + \varepsilon_{ij}), \quad (2)$$

with the SM recovered for $\varepsilon_{ij} \rightarrow 0$. Here we neglected tiny neutrino masses and set the leptonic mixing matrix to unity. Furthermore, we will disregard flavor-violating ε_{ij} parameters in the following since they are tightly bounded by radiative lepton decays $\ell \rightarrow \ell' \gamma$ and lead to contributions that do not interfere with the SM in observables testing LFU. Note that in Eq. (2) we simply parametrize the BSM effect by ε_{ij} , but do not consider the $SU(2)_L$ gauge invariance in SMEFT as discussed in Ref. [75].

The first crucial observation is that the corresponding modification of the charged current affects the determination of the Fermi constant G_F from the muon lifetime [76]

$$\frac{1}{\tau_\mu} = \frac{(G_F^L)^2 m_\mu^5}{192\pi^3} (1 + \Delta q)(1 + \varepsilon_{ee} + \varepsilon_{\mu\mu})^2, \quad (3)$$

where G_F^L is the Fermi constant appearing in the Lagrangian (excluding BSM contamination) and Δq subsumes the phase space, QED, and hadronic radiative corrections. Accordingly, we conclude that the Fermi constant measured in muon decay (extracted under the SM assumption) is related to the one at the Lagrangian level (containing the fundamental parameters M_W and g_2) as

$$G_F = G_F^L (1 + \varepsilon_{ee} + \varepsilon_{\mu\mu}). \quad (4)$$

TABLE I. Ratios sensitive to LFUV in the μ - e sector, indicating the dependence on the LFU violating parameters ε_{ij} . For $R(V_{us})$ we give the values corresponding to the radiative corrections from Refs. [11,14]. The last column gives the constraints on $(\varepsilon_{\mu\mu} - \varepsilon_{ee}) \times 10^3$ and $\varepsilon_{\mu\mu} \times 10^3$, respectively.

| Observable | Measurement | Constraint |
|--|----------------------|------------|
| $K \rightarrow \pi\mu\bar{\nu}/K \rightarrow \pi e\bar{\nu} \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 1.0010(25) [77] | 1.0(2.5) |
| $K \rightarrow \mu\nu/K \rightarrow e\nu \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 0.9978(18) [3,78,79] | -2.2(1.8) |
| $\pi \rightarrow \mu\nu/\pi \rightarrow e\nu \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 1.0010(9) [3,80–82] | 1.0(9) |
| $\tau \rightarrow \mu\nu\bar{\nu}/\tau \rightarrow e\nu\bar{\nu} \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 1.0018(14) [3,32] | 1.8(1.4) |
| $W \rightarrow \mu\bar{\nu}/W \rightarrow e\bar{\nu} \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 0.9960(100) [83,84] | -4(10) |
| $B \rightarrow D^{(*)}\mu\nu/B \rightarrow D^{(*)}e\nu \simeq 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee}$ | 0.9890(120) [85] | -11(12) |
| $R(V_{us}) \simeq 1 - (V_{ud}/V_{us})^2 \varepsilon_{\mu\mu}$ | 0.9891(33) [11] | 0.58(17) |
| | 0.9927(39) [14] | 0.39(21) |

This correction has to be taken into account whenever considering a weak decay, unless normalized to another branching ratio subject to the same correction.

This redefinition of the Fermi constant affects the determination of V_{ud} from all β -decay observables in the same way

$$V_{ud}^\beta = V_{ud}^\mathcal{L}(1 - \varepsilon_{\mu\mu}), \quad (5)$$

again denoting by $V_{ij}^\mathcal{L}$ CKM matrix elements without any BSM contamination, which therefore, by definition, fulfill CKM unitarity exactly. In particular, the indirect corrections introduced via G_F imply that β -decay observables are actually sensitive to LFUV in the muon, not the electron sector. To construct the LFU violating observable in Eq. (1) we further define

$$V_{us}^\beta \equiv \sqrt{1 - (V_{ud}^\beta)^2 - |V_{ub}|^2} \simeq V_{us}^\mathcal{L} \left[1 + \left(\frac{V_{ud}^\mathcal{L}}{V_{us}^\mathcal{L}} \right)^2 \varepsilon_{\mu\mu} \right]. \quad (6)$$

It is this crucial enhancement by $(V_{ud}/V_{us})^2 \sim 20$ that generates the amplified sensitivity to LFUV of our proposed observable $R(V_{us})$.

Before turning to the numerical analysis, we compare the sensitivity of $R(V_{us})$ to that of other probes of LFUV. Apart from $K_{\ell 3}$ decays, as given in Eq. (16) below, this includes $K_{\ell 2}$ and $\pi_{\ell 2}$, $\tau \rightarrow \ell\nu\bar{\nu}$, and $W \rightarrow \ell\nu$, see Table I for their dependence on the ε_{ii} as well as current experimental constraints. Concerning B decays, only $B \rightarrow D^{(*)}e\nu/B \rightarrow D^{(*)}\mu\nu$ provides a relevant constraint [85].

A crucial advantage of $R(V_{us})$ is that all these ratios are sensitive to the difference $\varepsilon_{\mu\mu} - \varepsilon_{ee}$, not LFUV in either sector separately, and thus can only test LFU in case the ε_{ii} differ. In addition, none of the other ratios can probe LFU at a level below $\mathcal{O}(10^{-3})$ yet, demonstrating the superior sensitivity of $R(V_{us})$ thanks to the CKM enhancement. We illustrate the comparison and complementarity of $R(V_{us})$ with respect to the other observables testing LFU in the W - ℓ - ν couplings in Fig. 2, anticipating the results from the following numerical analysis.

We start from the master formula for superallowed β decays [6]

$$|V_{ud}|^2 = \frac{2984.432(3) \text{ s}}{\mathcal{F}t(1 + \Delta_R^V)}, \quad (7)$$

with $\mathcal{F}t$ value $\mathcal{F}t = 3072.07(63) \text{ s}$ [6] and two sets of radiative corrections

$$\Delta_R^V|_{\text{SGPR}} = 0.02467(22) \quad [11], \quad (8)$$

$$\Delta_R^V|_{\text{CMS}} = 0.02426(32) \quad [14], \quad (9)$$

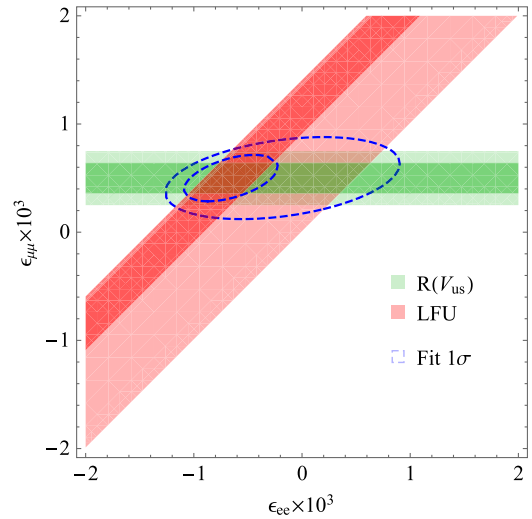


FIG. 2. Fit to modified W - ℓ - ν couplings. Here the light bands (large ellipse) correspond to the current status (1σ), where the band for $R(V_{us})$ has been increased to include both the CMS and the SGPR results. The dark bands (small ellipse) correspond to a future scenario in which the errors on $\tau \rightarrow \mu\nu\bar{\nu}/\tau \rightarrow e\nu\bar{\nu}$ and $\pi \rightarrow \mu\nu/\pi \rightarrow e\nu$ have decreased by a factor 3 due to forthcoming Belle II and PEN/PiENU results, respectively. The projected $R(V_{us})$ band assumes that the difference between the radiative corrections of CMS and SGPR has been understood and that a competitive V_{ud} determination from neutron decay has become available.

leading to

$$\begin{aligned} V_{ud}^\beta|_{\text{SGPR}} &= 0.97370(14), & V_{us}^\beta|_{\text{SGPR}} &= 0.22782(62), \\ V_{ud}^\beta|_{\text{CMS}} &= 0.97389(18), & V_{us}^\beta|_{\text{CMS}} &= 0.22699(78), \end{aligned} \quad (10)$$

where we used $|V_{ub}| = 0.003683$ from Refs. [86,87], although the precise value of $|V_{ub}|$ is immaterial here. In addition to the universal electroweak corrections Δ_R^V , it has been pointed out in Refs. [12,13] that also $\mathcal{F}t$ may be subject to additional nuclear corrections. The final recommendation $\mathcal{F}t = 3072(2)$ s in Ref. [13] leaves the central value largely unchanged, but implies a significant increase in uncertainty. Since the role of nuclear corrections in $0^+ \rightarrow 0^+$ transitions is far from settled, we continue to employ $\mathcal{F}t$ from Ref. [6], keeping in mind that the nuclear uncertainties are potentially underestimated.

For the neutron lifetime we use the master formula [7,14]

$$|V_{ud}|^2 \tau_n (1 + 3g_A^2)(1 + \Delta_{\text{RC}}) = 5100.1(7) \text{ s}, \quad (11)$$

with radiative corrections

$$\Delta_{\text{RC}}^{\text{SGPR}} = 0.03992(22), \quad \Delta_{\text{RC}}^{\text{CMS}} = 0.03947(32). \quad (12)$$

The experimental value for g_A from the asymmetry parameter is confronted with the lattice-QCD calculation $g_A = 1.271(13)$ [20] (see Ref. [88] for a critical assessment of the error estimate). Finally, pion β decay gives [15,16]

$$V_{ud}^\beta = 0.9739(29). \quad (13)$$

With the enhanced sensitivity to LFUV originating solely from V_{us}^β , it is less crucial which determination enters the numerator in Eq. (1). V_{us} can be determined directly from semileptonic kaon decays $K_{\ell 3}$. Using the compilation from Ref. [77] (updating Ref. [89]) as well as the form factor normalization $f_+(0) = 0.9698(17)$ [77,90,91], we have

$$\begin{aligned} V_{us}^{K_{\mu 3}} &= 0.22345(54)(39) = 0.22345(67), \\ V_{us}^{K_{e 3}} &= 0.22320(46)(39) = 0.22320(61), \end{aligned} \quad (14)$$

where the first error refers to experiment and the second to the form factor. LFUV affects these values according to

$$\begin{aligned} V_{us}^{K_{\mu 3}} &= V_{us}^{\mathcal{L}}(1 - \varepsilon_{ee}), \\ V_{us}^{K_{e 3}} &= V_{us}^{\mathcal{L}}(1 - \varepsilon_{\mu\mu}), \end{aligned} \quad (15)$$

leading to the constraint

$$R(K_{\ell 3}) = \frac{V_{us}^{K_{\mu 3}}}{V_{us}^{K_{e 3}}} = 1 + \varepsilon_{\mu\mu} - \varepsilon_{ee} = 1.0010(25), \quad (16)$$

where several uncertainties cancel in the ratio [77].

For the purely leptonic kaon decays $K_{\ell 2}$, one typically considers the ratio $K \rightarrow \mu\nu$ over $\pi \rightarrow \mu\nu$ to cancel the dependence on absolute decay constants. This allows one to directly determine $V_{us}^{\mathcal{L}}/V_{ud}^{\mathcal{L}}$ once the ratio of decay constants f_{K^\pm}/f_{π^\pm} as well as the treatment of isospin-breaking corrections are specified [92,93]. Here, we follow the strategy in Ref. [77] to use the recent results from lattice QCD [93], at the same time adjusting the FLAG average [94] back to the isospin limit $f_{K^\pm}/f_{\pi^\pm} = 1.1967(18)$ [95–97], to obtain

$$V_{ud}^{K_{\mu 2}} = 0.97427(10), \quad V_{us}^{K_{\mu 2}} = 0.22534(42). \quad (17)$$

The tension with the determinations from $K_{\ell 3}$ (14) cannot be explained with LFUV. For definiteness, we will use the $K_{\mu 2}$ value as reference point in Eq. (1), given that in contrast to $K_{\ell 3}$ it is, by definition, not sensitive to LFUV. Note that this is a conservative choice in the sense that the value of $R(V_{us})$ lies closer to unity if $K_{\mu 2}$ and not $K_{\ell 3}$ is used. One could extend the analysis further to determinations from τ decays, see Ref. [32], but here the errors are larger and again there are tensions between the inclusive and exclusive determinations. Therefore, for simplicity we restrict the analysis to observables sensitive only to the μ - e sector.

Numerically, V_{us}^β from Eq. (10) and $V_{us}^{K_{\mu 2}}$ from Eq. (17) provide the constraint

$$\varepsilon_{\mu\mu}|_{\text{SGPR}} = 0.00058(17), \quad \varepsilon_{\mu\mu}|_{\text{CMS}} = 0.00039(21), \quad (18)$$

and thus a sensitivity to LFUV below $\mathcal{O}(10^{-3})$. If instead the V_{us} values from $K_{\ell 3}$ decays were used, the central values would increase to $\varepsilon_{\mu\mu} \sim 0.001$ with similar errors as in Eq. (18), thus implying a much higher significance.

Future prospects.—Future improvements of the analysis presented here are foreseen at several frontiers: (i) the numerical value and accuracy of $R(V_{us})$ could be consolidated with improved radiative corrections; (ii) improved experimental input for neutron lifetime and asymmetry parameter could make the resulting V_{ud} determination competitive with superallowed β decays; (iii) new data could shed light on the tension between $K_{\ell 2}$ and $K_{\ell 3}$ decays.

First, for the superallowed $0^+ \rightarrow 0^+$ transitions the main uncertainty at this point originates from radiative corrections, both universal radiative corrections that also affect neutron decay as well as additional nuclear effects. The latter should be amenable to refined calculations with modern nuclear structure theory (see, e.g., Ref. [98] for

a recent *ab initio* calculation of nuclear β decays and Ref. [99] for a discussion of the nuclear theory requirements). Meanwhile, improving the universal radiative corrections rests on a better understanding of the nucleon matrix elements $\langle p | T \{ j_{\text{em}}^\mu j_{W,A}^\nu \} | n \rangle$ involving the electromagnetic current j_{em}^μ and the axial part of the charged weak current $j_{W,A}^\mu$, for which either new input from experiment or lattice QCD [100] is required.

Second, the measurement of λ is currently dominated by the measurement of Ref. [9] (with small changes if earlier results from Refs. [101,102] are included in the average), leading to a relative precision of 4×10^{-4} . There are several ongoing and planned developments that promise to extend the sensitivity towards or even beyond the level of 10^{-4} [103–107]. To establish the V_{ud} determination from neutron decay at a level competitive with superallowed β decays commensurate improvements in the lifetime are necessary. In addition to the long-standing discrepancy between bottle and beam measurements (see Refs. [7,108] for reviews), also the difference between recent bottle measurements [8,109,110] currently leads to a non-negligible scale factor in the PDG average [3]. Fortunately, there are plans to probe τ_n at a level down to hundreds of ms [111–114].

Third, preliminary data on $K_{\ell 3}$ decays exist from the OKA [115] and KLOE-2 [116] experiments, with further input potentially from LHCb [117], NA62, and TREK [77]. Further insights on V_{us} could be obtained from semileptonic hyperon decays [118–122] given renewed experimental interest at BESIII [123,124], but would also require progress in lattice-QCD calculations of the hyperon form factors [125]. All of these developments (i)–(iii) should help establish or refute the current 2–3 σ hint (18) for LFUV in β decays.

In addition, there are several experimental developments dedicated to improving the LFU tests in Table I. The J-PARC E36 experiment aims at improving $K \rightarrow \mu\nu/K \rightarrow e\nu$ [126], while the ratio of $K_{\ell 3}$ decays could profit from the developments mentioned above. A similar sensitivity as in $R(V_{us})$ may be possible for $\tau \rightarrow \mu\nu\bar{\nu}/\tau \rightarrow e\nu\bar{\nu}$ at Belle II [127], where approximately 1 order of magnitude more τ leptons will be produced than at Belle or BABAR. At this level, one would directly probe the same parameter space as in Eq. (18), barring of course a significant cancellation between $\epsilon_{\mu\mu}$ and ϵ_{ee} . The most promising observable, however, is currently $\pi \rightarrow \mu\nu/\pi \rightarrow e\nu$, for which the PEN [128] and PiENU [129] experiments anticipate in the near future an improvement by more than a factor of 3, which would bring the limit on $\epsilon_{\mu\mu} - \epsilon_{ee}$ well below $\mathcal{O}(10^{-3})$ as well. Taking into account all these potential improvements, we also included an optimistic but realistic projection of future constraints in Fig. 2.

Moving beyond pure modification of W - ℓ - ν couplings, one can see from an analysis of gauge-invariant dimension-6

operators that a simultaneous modification of Z - ℓ - ℓ and/or Z - ν - ν is unavoidable: there are only two operators modifying these couplings [130,131], so that the effects in at most one of these three couplings can be canceled. The LEP bounds on Z - ℓ - ℓ couplings are already now at the per mille level [84] and also the bounds on the invisible Z width (corresponding to Z - ν - ν in the SM) are excellent. These bounds could be significantly improved by future e^+e^- colliders such as an ILC [132], CLIC [133], or an FCC-ee [134,135]. Furthermore, W pair production will allow for a direct determination of $W \rightarrow \mu\nu/W \rightarrow e\nu$. In particular, an FCC-ee could produce up to 10^8 W bosons (compared to the LEP number of 4×10^4), leading to an increase in precision that would render a direct discovery of LFUV in W - ℓ - ν conceivable.

Conclusions.— β decays are high-precision low-energy tests of the SM. While so far these decays were used to constrain CKM unitarity or right-handed currents (from superallowed β decays and neutron decay), we showed in this Letter that they are also an exquisite probe of LFU: due to the conventional definition of the Fermi constant they actually probe LFU in the muon sector, with a sensitivity CKM enhanced by $(V_{ud}/V_{us})^2 \sim 20$. Therefore, we proposed and examined the ratio (1) to test LFU. This measure $R(V_{us})$ is complementary to other probes of LFU, most notably $\pi \rightarrow \mu\nu/\pi \rightarrow e\nu$ and $\tau \rightarrow \mu\nu\bar{\nu}/\tau \rightarrow e\nu\bar{\nu}$, which are sensitive to the difference of muon and electron couplings.

Current data show a deviation of $R(V_{us})$ from unity at the level of 2–3 σ , depending on assumptions for the radiative corrections. In light of the accumulated hints for LFUV in $R(K^{(*)})$, $R(D^{(*)})$, and $(g-2)_{\mu,e}$, it seems natural to consider β decays as a probe of LFU and refine complementary tests of LFU with this connection in mind. We discussed several avenues how the present constraints can be improved in the future, including experimental developments in kaon and neutron decays, which, however, should be accompanied by adequate efforts on the theory side aiming at improving our understanding of radiative corrections. Similar improvements are anticipated in related tests of LFU, with results expected soon from the PEN and PiENU experiments, while future e^+e^- colliders, in particular the FCC-ee, would even have the potential to directly observe LFUV in W decays.

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