

**Erratum: High-Precision Mass Measurement of  $^{56}\text{Cu}$   
and the Redirection of the  $rp$ -Process Flow  
[Phys. Rev. Lett. **120**, 032701 (2018)]**

A. A. Valverde<sup>✉</sup>, M. Brodeur, G. Bollen, M. Eibach, K. Gulyuz, A. Hamaker, C. Izzo, W.-J. Ong,  
D. Puentes, M. Redshaw, R. Ringle, R. Sandler, S. Schwarz, C. S. Sumithrarachchi,  
J. Surbrook, A. C. C. Villari, and I. T. Yandow

 (Received 4 November 2019; published 3 December 2019)

DOI: [10.1103/PhysRevLett.123.239905](https://doi.org/10.1103/PhysRevLett.123.239905)

In the Letter, we presented a new  $^{55}\text{Ni}(p, \gamma)^{56}\text{Cu}$  reaction rate based on our measured  $^{56}\text{Cu}$  mass and the resonance information taken from [1]. The rate reported in the Letter is incorrect due to an incorrect scaling of the penetrability factor

TABLE I. The recommended reaction rate  $N_A \langle \sigma v \rangle$  as a function of temperature  $T$  (GK) from this work, together with  $1\text{-}\sigma$  uncertainties (higher and lower).

$T_9$	$N_A \langle \sigma v \rangle$ (cm <sup>3</sup> /s/mole)		
	Recommended	Lower	Upper
0.100	4.487e-19	4.158e-19	4.612e-19
0.200	5.301e-12	4.107e-12	8.781e-12
0.300	2.199e-08	1.742e-08	2.894e-08
0.400	1.743e-06	1.120e-06	2.621e-06
0.500	2.931e-05	1.841e-05	4.950e-05
0.600	2.141e-04	1.262e-04	4.048e-04
0.700	9.195e-04	5.080e-04	1.906e-03
0.800	2.795e-03	1.450e-03	6.198e-03
0.900	6.767e-03	3.292e-03	1.563e-02
1.000	1.403e-02	6.428e-03	3.295e-02
1.500	1.731e-01	7.422e-02	3.388e-01
2.000	8.899e-01	4.791e-01	1.450e+00
3.000	6.682e+00	4.962e+00	9.097e+00
3.500	1.336e+01	1.073e+01	1.695e+01
4.000	2.402e+01	2.046e+01	2.874e+01
4.500	4.058e+01	3.615e+01	4.627e+01
5.000	6.590e+01	6.074e+01	7.238e+01
6.000	1.604e+02	1.541e+02	1.679e+02
7.000	3.554e+02	3.486e+02	3.634e+02
8.000	7.193e+02	7.122e+02	7.273e+02
9.000	1.338e+03	1.331e+03	1.346e+03
10.000	2.313e+03	2.306e+03	2.321e+03

TABLE II. REACLIB fit coefficients for our corrected  $^{55}\text{Ni}(p, \gamma)$  reaction rate, with the fit performed over the temperature range of 0.1 to 10 GK.

$a_0$	$a_1$	$a_2$	$a_3$	$a_4$	$a_5$	$a_6$
3.046255E+00	-7.508431E+00	2.049851E+00	-2.648531E+00	1.357497E-01	-7.098837E-03	-1.061167E-01
-7.516925E+01	-9.131942E+00	5.982367E+01	2.418860E+01	-5.181901E+00	3.801069E-01	1.526078E+01
-1.968345E+00	-5.296513E+00	-3.830019E-04	5.743085E-04	-3.599905E-05	2.300254E-06	-1.500277E+00
-1.718560E+01	-2.858869E+00	8.588120E-01	-1.036879E+00	5.156693E-02	-2.666759E-03	-9.389072E-01
2.730636E+01	4.610386E-06	-3.894134E+01	1.459666E-04	-6.651521E-06	2.754834E-07	-6.667493E-01

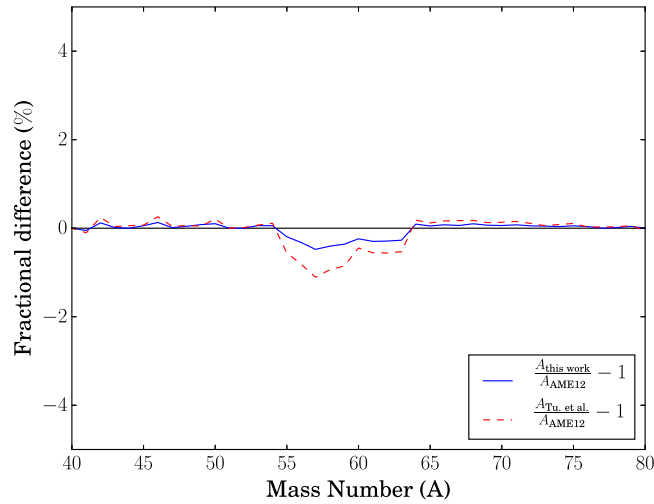


FIG. 1. Fractional difference of abundance by mass number using the mass of  $^{56}\text{Cu}$  measured in this work (solid blue line), and [3] (dashed red line), compared to that using the mass suggested in AME2012 [5].

in the proton widths. In this Erratum, we report an updated rate along with recommended REACLIB parameters, and we reevaluate the change in the mass distribution of  $rp$ -process ashes.

The corrected rate for the 579.8 keV  $Q$  value reported in the Letter is given in Table I, with the REACLIB parameters given in Table II. The change in the mass distribution of the  $rp$ -process ashes (similar to Fig. 5 in the Letter) is shown in Fig. 1. The recalculated ash composition, using a single zone model, for a typical x-ray burst trajectory [2] shows a less dramatic effect on the ashes using both the mass reported in this work as well as the mass reported in [3] (which was also incorrectly scaled in the Letter). We would like to stress, however, that these calculations assume a large  $\beta$ -delayed proton emission branch of 78% for  $^{57}\text{Zn}$  from [4], which would return most of the bypassed flow back into the  $^{56}\text{Ni}$  waiting point. With the mass measurement of  $^{56}\text{Cu}$  reported in this Letter significantly constraining the  $Q$  value of the first reaction in the bypass sequence  $^{55}\text{Ni}(p, \gamma)^{56}\text{Cu}$ , the  $\beta$ -delayed proton emission branch of  $^{57}\text{Zn}$  is now the most significant uncertainty to determine if a bypass of the  $^{56}\text{Ni}$   $rp$ -process waiting point is strong.

We would like to thank the authors of [6] for pointing out this error.

- [1] W.-J. Ong *et al.*, Low-lying level structure of  $^{56}\text{Cu}$  and its implications for the  $rp$  process, *Phys. Rev. C* **95**, 055806 (2017).
- [2] R. Cyburt, A. Amthor, A. Heger, E. Johnson, L. Keek, Z. Meisel, H. Schatz, and K. Smith, Dependence of X-ray burst models on nuclear reaction rates, *Astrophys. J.* **830**, 55 (2016).
- [3] X. Tu, Y. A. Litvinov, K. Blaum, B. Mei, B. Sun, Y. Sun, M. Wang, H. Xu, and Y. Zhang, Indirect mass determination for the neutron-deficient nuclides  $^{44}\text{V}$ ,  $^{48}\text{Mn}$ ,  $^{52}\text{Co}$  and  $^{56}\text{Cu}$ , *Nucl. Phys.* **A945**, 89 (2016).
- [4] B. Blank, C. Borcea, G. Canchel, C.-E. Demonchy, F. de Oliveira Santos, C. Dossat, J. Giovanazzo, S. Grévy, L. Hay, P. Hellmuth *et al.*, Production cross-sections of proton-rich  $^{70}\text{Ge}$  fragments and the decay of  $^{57}\text{Zn}$  and  $^{61}\text{Ge}$ , *Eur. Phys. J. A* **31**, 267 (2007).
- [5] M. Wang, G. Audi, A. Wapstra, F. Kondev, M. MacCormick, X. Xu, and B. Pfeiffer, The ame2012 atomic mass evaluation, *Chin. Phys. C* **36**, 1603 (2012).
- [6] S.-B. Ma, L.-Y. Zhang, and J. Hu, Stellar reaction rate of  $^{55}\text{Ni}(p, \gamma)^{56}\text{Cu}$  in type I X-ray bursts, *Nucl. Sci. Tech.* **30**, 141 (2019).