Search for Leptoquarks Coupled to Third-Generation Quarks in Proton-Proton Collisions at $\sqrt{s} = 13$ TeV

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(Received 14 September 2018; published 12 December 2018)

Three of the most significant measured deviations from standard model predictions, the enhanced decay rate for $B \to D^{(*)}\tau\nu$, hints of lepton universality violation in $B \to K^{(*)}\ell\ell$ decays, and the anomalous magnetic moment of the muon, can be explained by the existence of leptoquarks (LQs) with large couplings to third-generation quarks and masses at the TeV scale. The existence of these states can be probed at the LHC in high energy proton-proton collisions. A novel search is presented for pair production of LQs coupled to a top quark and a muon using data at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 35.9 fb^{-1}, recorded by the CMS experiment. No deviation from the standard model prediction has been observed and scalar LQs decaying exclusively into $t\mu$ are excluded up to masses of 1420 GeV. The results of this search are combined with those from previous searches for LQ decays into $t\tau$ and $b\nu$, which excluded scalar LQs below masses of 900 and 1080 GeV. Vector LQs are excluded up to masses of 1190 GeV for all possible combinations of branching fractions to $t\mu$, $t\tau$ and $b\nu$. With this analysis, all relevant couplings of LQs with an electric charge of $-1/3$ to third-generation quarks are probed for the first time.

DOI: 10.1103/PhysRevLett.121.241802

The standard model of particle physics has been outstandingly successful in describing most fundamental physical phenomena. However, significant deviations from the predictions of the standard model (SM) have been observed in measurements of rare decays of B mesons. In particular, deviations have been seen in the values of the ratio $R_{D^{(*)}}$, defined as the ratio of the $B \to D^{(*)}\tau\nu$ branching fraction to the $B \to D^{(*)}\mu\nu$ branching fraction. These deviations from the SM were first reported by the *BABAR* [1,2] and Belle [3–5] Collaborations and have been confirmed by the LHCb Collaboration [6,7] with a combined significance of about four standard deviations [8]. The ratios of the branching fractions of $B \to K^{(*)}\mu\mu$ to $B \to K^{(*)}ee$, R_K and R_{K^*}, as measured by the LHCb Collaboration [9–12], show departures from lepton universality by 2.6 and 2.4 standard deviations, respectively. The measurement of the muon anomalous magnetic moment a_μ, one of the most precisely measured quantities in particle physics [13], also deviates from the SM prediction by 3.5 standard deviations [14]. These anomalies are among the most significant deviations from the SM observed so far.

The existence of leptoquarks (LQs) with masses at the TeV scale and large couplings to third-generation quarks [15–25] has been proposed as a possible explanation for one, two, or all of these deviations. Leptoquarks are hypothetical particles that can decay to SM quarks and leptons. They are triplets with respect to the strong interaction, have fractional electric charge, and can be either scalar (spin 0) or vector (spin 1) particles. Many extensions to the SM, among them grand unification [26–28], technicolor [29,30], and compositeness models [31,32], predict the existence of these particles. The effective Buchmüller-Rückl-Wyler model [33] incorporates the assumption that LQ interactions with SM fermions are renormalizable and gauge invariant, leading to restrictions on the allowed quantum numbers of LQs [34]. Depending on its quantum numbers and the coupling structure, a given LQ can decay to any one of a number of different combinations of SM fermions. The couplings of LQs to leptons and quarks of different generations introduce flavor changing neutral currents that may be observable in precision measurements [35]. While simultaneous couplings to the first and second generations are tightly constrained by experimental data, the bounds are weaker for couplings to the second and third generation, thus allowing the existence of leptoquarks with nondiagonal couplings in the generation matrix [19,24,36].

Collider searches for LQs with decays to third-generation quarks have been performed in the decay channels LQ → $t\tau$, LQ → $b\tau$, and LQ → $b\nu$ at $\sqrt{s} = 8$ TeV [37–44] and recently at $\sqrt{s} = 13$ TeV [45–49]. We present

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the first search for the pair production of LQs with decays to a top quark and a muon, LQ → $t\mu$, a decay mode that is essential to explain the anomalies in a_μ and $R_{K^{(*)}}$ [19–25]. This search is combined with previous searches that target other decay modes [48,49]. The combination provides sensitivity to all relevant couplings of LQs with an electric charge of $-1/3$ to third-generation quarks.

At the CERN LHC, pair production of LQs is possible via gluon-gluon fusion or quark-antiquark annihilation, allowing direct searches to be performed. Single LQ production via quark-gluon scattering is subdominant for LQs coupled to heavy quarks, as it requires a heavy quark in the initial state. The pair production cross section depends on the mass of the scalar LQ and is known at next-to-leading order (NLO) precision [50]. The pair production cross section for vector LQs has been calculated at leading order (LO) [51] and is much larger than the scalar LQ cross section. The cross section for vector LQs depends on an additional parameter κ, which is a dimensionless coupling and takes a value of $\kappa = 1$ in the Yang-Mills case and $\kappa = 0$ in the minimal coupling case.

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity (η) coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [52].

This analysis uses data recorded by the CMS detector in pp collisions at a center-of-mass energy of 13 TeV in 2016. Online, potential signal events are required to pass a single-muon trigger that selects isolated muon candidates with transverse momentum $p_T > 24$ GeV [53]. Data recorded by single electron triggers are used in background-enriched control regions (CRs). The data correspond to an integrated luminosity of 35.9 fb^{-1}.

Signal events of pair-produced LQs with prompt decays to $t\mu$ are simulated with the PYTHIA 8.205 [54,55] Monte Carlo program at LO for mass values ranging from 200 to 2000 GeV. The POWHEG [56–63] v1 generator is used to simulate background events resulting from the production of single top quarks in the tW channel at NLO. The POWHEG v2 generator is used for single top production in the t channel and for simulating $t\bar{t}$ production at NLO. Single top quark production in the s channel, $t\bar{t}$ production in association with a heavy gauge boson ($t\bar{t}+\mathrm{V}$), and the production of a W boson with additional jet radiation are simulated with MADGRAPH 5_aMC@NLO (v2.2.2) [64] at NLO. Events from Drell-Yan (DY) production with additional jet radiation are simulated with MADGRAPH 5_aMC@NLO at LO and an NLO K factor is applied to the LO DY + jets production cross section. The simulation of the production of two heavy gauge bosons with additional jet radiation is performed at NLO with MADGRAPH 5_aMC@NLO and POWHEG v2. Events in which jets are produced through the strong interaction only, referred to as quantum chromodynamic multijet events, are simulated with PYTHIA at LO.

Parton showers in the simulated W boson production events and DY events with additional jet radiation are matched to the matrix element calculation with the FXFX [65] and MLM [66] algorithms, respectively. The parton shower and hadronization process is simulated with PYTHIA. The NNPDF3.0 [67] parton distribution functions (PDFs) at LO and NLO are used for processes simulated at LO and NLO, respectively. The underlying event tune CUETP8M2T4 [68] is used for the simulation of $t\bar{t}$ and single top quark production via the t channel, all other processes are generated using CUETP8M1 [69,70]. All simulated event samples include the simulation of additional inelastic pp interactions within the same or adjacent bunch crossings (pileup). The detector response is simulated with the GEANT4 package [71,72]. Simulated events are processed through the software chain used for collision data and are reweighted to match the observed distribution of the number of pileup interactions in data.

The CMS experiment uses a particle-flow (PF) event reconstruction algorithm [73], which makes use of an optimized combination of information from the various elements of the CMS detector. The reconstructed vertex with the largest value of summed physics object p_T^2 is taken to be the primary pp interaction vertex. The physics objects here are the objects returned by a jet finding algorithm [74,75] applied to all charged tracks associated with the vertex, plus the associated missing transverse momentum, taken as the negative vector p_T sum of those jets. More details are given in Ref. [76]. All detected particles are reconstructed either as electrons, muons, photons, charged hadrons, or neutral hadrons. In this analysis, electrons and muons are required to have $p_T \geq 30$ GeV, $|\eta| \leq 2.4$, and to be isolated. The isolation [77,78] is defined as the summed p_T of all neutral particles and charged hadrons in a cone with radius $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$, ϕ being the azimuthal angle in radians, of 0.4 (for muons) or 0.3 (for electrons) around the lepton. The sum is corrected for the contribution of neutral pileup inside the cone. Jets are clustered from charged and neutral PF candidates using the anti-k_T jet-clustering algorithm [74,75] with a distance parameter of 0.4. Charged PF candidates originating from vertices other than the primary vertex are not clustered. A jet energy correction (JEC) is applied [79] to account for remaining contributions arising from a different vertex than the primary one as well as for nonuniformity of the jet response in η and nonlinearity in p_T. Finally, a correction is applied to account for the residual differences in the jet response between data and simulated events. The jet energy resolution (JER) in simulated events is smeared to match the wider resolution in data. All jets are required to have $p_T \geq 30$ GeV and $|\eta| \leq 2.4$. The combined secondary vertex v2 [80] algorithm is used to identify jets originating from bottom quarks (b-tagged jets). The loose working point is chosen, which has an efficiency of about 90% and a mistag rate of approximately 10%. The missing transverse momentum p_T^{miss} is calculated as the magnitude of the negative vectorial p_T sum of all PF candidates in an event. Both the jet energy scale and resolution corrections are propagated to the calculation of p_T^{miss}.

Offline, events are required to contain at least two muons and at least two jets, of which at least one must be b tagged. By requiring the invariant mass of each pair of muons in an event to exceed the Z boson mass by at least 20 GeV, events arising from the production of a Z boson with additional jet radiation are suppressed. As the decay of heavy LQs is expected to produce highly energetic leptons and jets, minimum values of S_T^{lep} and S_T of 200 and 350 GeV are required, respectively. Here, S_T^{lep} is the scalar p_T sum of all selected muons and electrons and S_T is defined as the scalar sum of S_T^{lep}, p_T^{miss}, and the p_T of all selected jets. The phase space region resulting from these selection criteria is referred to as the signal region (SR) in the following.

Two categories of events are defined, based on the number of muons or electrons. If at least three such charged leptons are present, of which at least two are muons, and at least one pair of muons with opposite electric charge is found, the event falls into category A. Category B contains all remaining events in the SR. In category A, the LQ mass $M_{\text{LQ}}^{\text{rec}}$ is reconstructed under the assumption that one of the top quarks decays into the leptonic final state with a muon or an electron (leptonic top) and the other one decays into the hadronic final state (hadronic top). The distribution of $M_{\text{LQ}}^{\text{rec}}$ is used for the final statistical analysis in this category, while the distribution of S_T is used for this purpose in category B.

For each event, the leptonic top quark candidates are constructed from permutations of one or more of the seven p_T-leading jets, one of the three p_T-leading muons or the p_T-leading electron, and p_T^{miss}. The hadronic top quark candidates are constructed using all permutations of jets not assigned to the leptonic top quark. The LQ candidates are assembled from top quark candidates and the two p_T-leading muons that have not been associated to the leptonic top quark. The muon charge is used when assigning it to one of the top quark candidates. In events with more than two muons, all possible permutations of muons are considered. A χ^2 variable that takes into account the invariant mass of each top quark candidate and the relative mass difference between the two LQ candidates is then used to select the best pair of LQ candidates for each event. Events with four leptons, which could originate from dileptonic $t\bar{t}$ decays, are included in category A and contribute to the signal efficiency. In order to provide a more accurate SM background prediction in category A, which contains a minimum of three charged leptons, the misidentification rate of electrons and muons is measured using jets in a DY + jets enriched CR in data. The CR is defined by selecting two muons with an invariant mass close to the Z boson mass, and the misidentification rate is measured on events where a jet is misidentified as a third lepton. The resulting data-to-simulation correction factors are applied to simulation for each misidentified charged lepton in a given event in the SR. The effect of charge misidentification on the analysis was found to be negligibly small.

The contributions from the dominant SM backgrounds in category B, the production of $t\bar{t}$ and DY + jets events, are estimated simultaneously in a data-driven procedure. A CR similar to the SR is defined by requiring a minimum of two electrons without additional muons. The invariant mass of any pair of electrons must be at least 20 GeV above the Z boson mass and all other SR requirements have to be fulfilled for the selected events. We correct for small differences in the distribution of S_T between the SR and the CR with an extrapolation function $\alpha(S_T)$, which is derived from simulated $t\bar{t}$ and DY + jets events by fitting both S_T distributions with an empirical functional form to obtain smoothed distributions that are then used to compute the ratio. The number of data events in the CR, after all simulated minor backgrounds have been subtracted, is multiplied by $\alpha(S_T)$ to extrapolate into the SR. Using the ratio of the fitted functions results in a significantly smaller impact of systematic uncertainties on the estimated backgrounds.

Various uncertainties affecting the rate and the shape of the signal and background contributions are taken into account. In general, uncertainties in this analysis are treated similarly to those in Ref. [48]. For the background in category A, the uncertainties in the renormalization and factorization scales as well as the uncertainties in the lepton misidentification rates are dominant. In category B the major backgrounds are derived from data. Lepton efficiencies and the background extrapolation procedure are the most important sources of uncertainties for these backgrounds. Uncertainties in the renormalization and factorization scales and in those associated with the choice of PDFs [67,81] used to simulate the events dominate for the minor backgrounds. The signal in both categories is most affected by the uncertainties in lepton and b-tagging [80] efficiencies. Other uncertainties considered are related to SM cross sections [82–90], the integrated luminosity [91], JEC and JER [79], and the pileup reweighting [92].

The THETA software package [93] is used to perform a maximum-likelihood template fit to the binned $M_{\text{LQ}}^{\text{rec}}$ and S_T distributions for the background and to extract the cross section of a potential signal. The statistical uncertainties in the SM backgrounds and the signal, as well as all systematic uncertainties, are taken into account as nuisance parameters in the fit. The uncertainty in the luminosity is assigned a log-normal prior distribution, for all other systematic uncertainties a Gaussian prior is used. The statistical uncertainty in the predicted background and the signal is taken into account by defining one additional nuisance parameter with a Gaussian distribution for each bin. A flat prior distribution is assumed for the signal cross section. The data are found to be compatible with the SM prediction in both categories. The distributions of $M_{\mathrm{LQ}}^{\mathrm{rec}}$ and S_T after the background-only fit are shown in Fig. 1. A Bayesian method [93–95] is used to set upper limits at 95% confidence level (C.L.) on the cross section for pair production of LQs decaying into a top quark and a muon. Pseudoexperiments are performed to determine the median along with the regions expected to contain 68% and 95% of the distribution of limits under the background-only hypothesis.

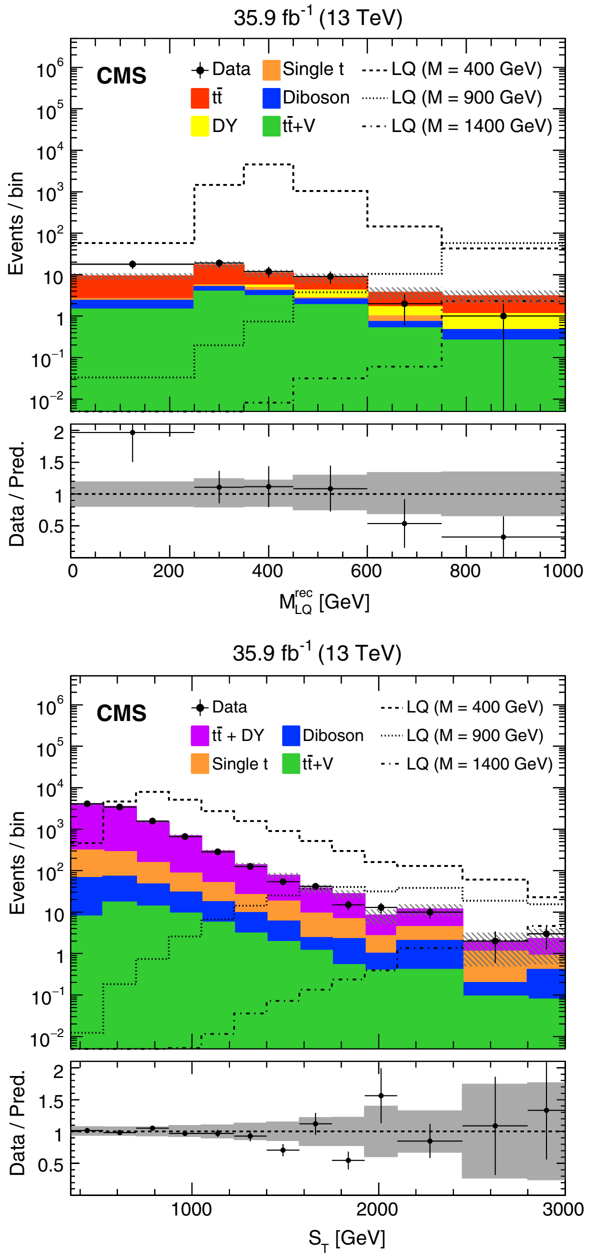

FIG. 1. Distributions for $M_{\mathrm{LQ}}^{\mathrm{rec}}$ (category A, left) and S_T (category B, right) after applying the full selection and estimating the $t\bar{t}$ and DY + jets background contributions from data in category B. All backgrounds are normalized according to the post-fit nuisance parameters based on the corresponding SM cross sections. In the upper panels, the hatched areas correspond to the total uncertainty. In the lower panels, the gray bands indicate the total uncertainty.

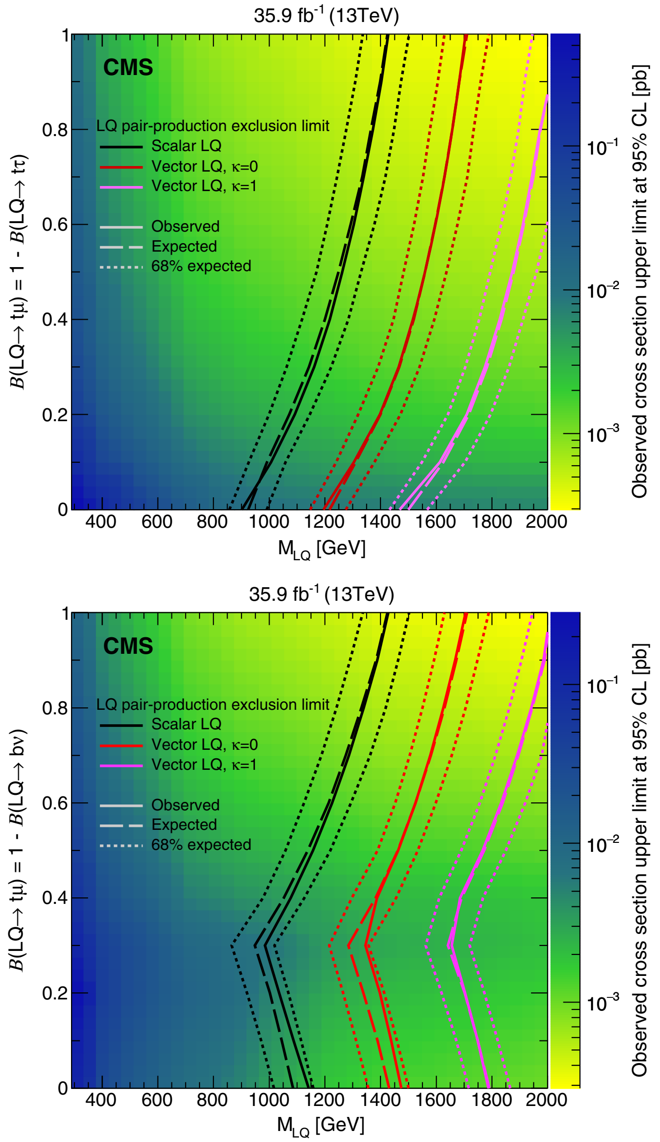

FIG. 2. Observed upper limits on the production cross section for pair production of LQs decaying into a top quark and a muon or a τ lepton (upper) and LQs decaying into a top quark and a muon or into a bottom quark and a neutrino (lower) at 95% C.L. in the M_{LQ}–$\mathcal{B}(\mathrm{LQ} \to t\mu)$ plane. The lines show the lower mass exclusion limits for scalar (black) and vector (colored) LQs. They are derived by using the prediction for the scalar and vector LQ signal calculated at NLO [50] and LO [51], respectively.

Pair-produced scalar LQs decaying exclusively into a top quark and a muon, $\mathcal{B}(\mathrm{LQ} \to t\mu) = 1$, are excluded at 95% C.L. for LQ masses up to 1420 GeV, exceeding the best previous limit, obtained from a reinterpretation [36] of a search for supersymmetry [96], by more than 600 GeV. These results are combined with results from the LQ → $t\tau$ [48] and LQ → $b\nu$ [49] decay channels to set exclusions limits in the plane of M_{LQ} and $\mathcal{B}(\mathrm{LQ} \to t\mu)$. Figure 2 presents upper limits on the product of the production cross section and the branching fraction squared for $\mathcal{B}(\mathrm{LQ} \to t\mu) = 1 - \mathcal{B}(\mathrm{LQ} \to t\tau)$ (upper) and $\mathcal{B}(\mathrm{LQ} \to t\mu) = 1 - \mathcal{B}(\mathrm{LQ} \to b\nu)$ (lower). The values for $\mathcal{B}(\mathrm{LQ} \to t\mu) = 0$ correspond to the results of the search for pair-produced LQs in the LQ → $t\tau$ decay channel (upper) and the search for pair-produced LQs in the LQ → $b\nu$ channel (lower). These analyses excluded pair-produced scalar LQs in the targeted decay channels up to $M_{\mathrm{LQ}} = 900$ and 1080 GeV, respectively. In the upper (lower) part of Fig. 2 the sensitivity is driven by the present analysis for values of $\mathcal{B}(\mathrm{LQ} \to t\mu) > 0.1(0.3)$ and by the LQ → $t\tau(b\nu)$ search for smaller values. Scalar LQs decaying into a top quark and either a muon or a τ lepton are excluded below masses of 900 GeV for all values of $\mathcal{B}(\mathrm{LQ} \to t\mu)$, whereas LQs decaying either into a top quark and a muon or into a bottom quark and a neutrino are excluded up to $M_{\mathrm{LQ}} = 980$ GeV. The simulated samples of scalar LQ pair production are also used to derive mass exclusion limits for pair-produced vector LQs, as the acceptance for both types of LQs is similar. The lower limit of excluded vector LQ masses is shown in Fig. 2 for the two coupling cases $\kappa = 1$ and $\kappa = 0$. Vector LQs are excluded up to masses of 1190 GeV for all values of $\mathcal{B}(\mathrm{LQ} \to t\mu)$ and κ considered.

In summary, this analysis represents the first search for leptoquarks decaying to top quarks and muons, reaching LQ masses of $\mathcal{O}(1\ \mathrm{TeV})$ and placing direct constraints on the corresponding LQ coupling, thus probing the region of interest of models including LQs. With this result, all relevant couplings of LQs with an electric charge of $-1/3$ to third-generation quarks are examined for the first time.

We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centers and personnel of the Worldwide LHC Computing Grid for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC and the CMS detector provided by the following funding agencies: BMWFW and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, and FAPESP (Brazil); MES (Bulgaria); CERN; CAS, MoST, and NSFC (China); COLCIENCIAS (Colombia); MSES and CSF (Croatia); RPF (Cyprus); SENESCYT (Ecuador); MoER, ERC IUT, and ERDF (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRT (Greece); OTKA and NIH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); MSIP and NRF (Republic of Korea); LAS (Lithuania); MOE and UM (Malaysia); BUAP, CINVESTAV, CONACYT, LNS, SEP, and UASLP-FAI (Mexico); MBIE (New Zealand); PAEC (Pakistan); MSHE and NSC (Poland); FCT (Portugal); JINR (Dubna); MON, RosAtom, RAS, RFBR and RAEP (Russia); MESTD (Serbia); SEIDI, CPAN, PCTI and FEDER (Spain); Swiss Funding Agencies (Switzerland); MST (Taipei); ThEPCenter, IPST, STAR, and NSTDA (Thailand); TUBITAK and TAEK (Turkey); NASU and SFFR (Ukraine); STFC (United Kingdom); DOE and NSF (USA).

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