the subject of solar neutrinos.

[2]A paper by the present authors is in preparation. It will contain a discussion of the various limits which can be set on nucleon stability and neutrino fluxes by a consideration of the unaccompanied counts. The earlier stages of this work are described in a paper by C. C. Giamati and F. Reines, Phys. Rev. 126, 2178 (1962).

[3]R. P. Feynman and M. Gell-Mann, Phys. Rev. 109, 193 (1958).

[4]R. E. Marshak and E. C. G. Sudarshan, Phys. Rev. 109, 1860 (1958); Proceedings of the Padua-Venice Conference on Mesons and Newly Discovered Particles, September, 1957 (Società Italiana di Fisica, Padua-Venice, 1958).

[5]Our results can also be used to set an upper limit on the product of the elastic scattering cross section averaged over the upper end of the ^{4}Li decay spectrum (on the unlikely assumption that it is particle stable) times the solar ^{4}Li-produced ν_e flux at the earth. The flux limit so obtained is $<2\times10^8$ ν_e/cm^2 sec. This is to be compared with the Bahcall-Davis limit of $<1\times10^8$ ν_e/cm^2 sec.

[6]C. L. Cowan, Jr., and F. Reines, Phys. Rev. 107, 528 (1957). This experiment was interpreted in terms of an upper limit on the neutrino magnetic moment. We here reinterpret these data in terms of the conserved vector current predictions.

[7]D. Strominger, J. M. Hollander, and G. T. Seaborg, Rev. Modern Phys. 30, 585 (1958).

[8]L. V. East, T. L. Jenkins, and F. Reines (unpublished).

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ABSOLUTE DECAY RATE FOR $K_2^0 \to \pi^+ + \pi^- + \pi^0$ AND THE $|\Delta I| = 1/2$ RULE*

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In this Letter we describe a measurement of the absolute decay rate $\Gamma_2(+-0) \equiv \Gamma(K_2^0 \to \pi^+ + \pi^- + \pi^0)$. Our result is based on 16 events of the type $\pi^- + p \to \Lambda + K^0$ followed by $\Lambda \to p + \pi^-$ and $K_2^0 \to \pi^+ + \pi^- + \pi^0$, and 2608 double-vee events $\pi^- + p \to \Lambda + K^0$ with $\Lambda \to p + \pi^-$ and $K_1^0 \to \pi^+ + \pi^-$. We find

$$\Gamma_2(+-0) = (2.90 \pm 0.72) \times 10^6\ \mathrm{sec}^{-1}. \tag{1}$$

This result can be compared with the prediction of the $|\Delta I| = 1/2$ rule for nonleptonic decays, that $\Gamma_2(+-0)$ and $\Gamma(+00) \equiv \Gamma(K^+ \to \pi^+ + \pi^0 + \pi^0)$ are related by

$$\Gamma_2(+-0) = 2(1.032)\Gamma(+00), \tag{2}$$

where the factor 1.032 corrects for small mass differences. Equation (2) holds for any linear combination of the three $I=1$ three-pion states.[1,2] Taking $\Gamma(+00)$ from a compilation of K^+ branching ratios and lifetimes,[3] one obtains the prediction of the $|\Delta I| = 1/2$ rule,

$$\Gamma_2(+-0) = (2.87 \pm 0.23) \times 10^6\ \mathrm{sec}^{-1}. \tag{3}$$

The excellent agreement between our experimental result (1) and the prediction (3) shows that the $|\Delta I| = 1/2$ rule is well satisfied.[4]

The ratio $\Gamma_2(+-0)/\Gamma(+00)$ affords a sensitive test of the $|\Delta I| = 1/2$ rule. To exhibit this sensitivity we parametrize the effect of a small $|\Delta I| = 3/2$ amplitude $A_{3/2}$ for $K \to 3\pi$ under the assumption that the dominant $|\Delta I| = 1/2$ amplitude leads to the symmetrical $I=1$ three-pion state, but with no such restriction on the $|\Delta I| = 3/2$ amplitude. Then, if $|A_{3/2}/A_{1/2}|^2$ is neglected, our experimental result expressed in the notation of reference 2 becomes

$$\sqrt{2}\,\mathrm{Re}\left(\frac{A_{3/2}}{A_{1/2}}\right) = \frac{2(1.032)\Gamma(+00) - \Gamma_2(+-0)}{4(1.032)\Gamma(+00) + \Gamma_2(+-0)} \tag{4}$$

$$= 0.00 \pm 0.09. \tag{5}$$

The remainder of this paper is concerned with experimental details.

The Alvarez 72-in. hydrogen bubble chamber was exposed to π^- beams with momenta between 1034 and 1325 MeV/c. All film was scanned for single and double vees which were analyzed using the least-squares fitting program KICK. For the present experimental events of the type $\pi^- + p^+ \to \Lambda + K^0$ were used; $\pi^- + p \to \Sigma^0 + K^0$ events were not used. All single-Λ events were rescanned along the direction of the unobserved K^0, as predicted by the fitting program. We believe that the resulting over-all efficiency for finding associated K decays is nearly 100%.

Most of the double vees fit the hypothesis $\pi^- + p \to \Lambda + K^0$ with $\Lambda \to p + \pi^-$ and $K_1^0 \to \pi^+ + \pi^-$. Those within the fiducial volume that fail because the assumption $K_1^0 \to \pi^+ + \pi^-$ fails are analyzed as follows. The fitted parameters for the decay $\Lambda \to p + \pi^-$ are propagated to the production vertex. At the production vertex we perform a three-constraint (3C) fit using the decay-fitted Λ, the incident π^-, and the measured direction of the K^0 as determined from the two-point neutral track. We accept events with production $\chi^2(3C) < 25$.[6] The K^0 momentum vector is then well known and is used as input information for the following hypotheses;

A. $K^0 \to \pi^+ + \pi^- + \pi^0$ (1C, π^0 unseen). For $\chi^2(1C) < 10$, the decay is a τ^0 candidate. About five events having an obvious $e^\pm$ or $\mu^\pm$ track are discarded. There are then 20 candidates.

B. $K_1^0 \to 2\pi^0 \to e^+ + e^- + \gamma + \pi^0$. We assume that the charged tracks are electrons and calculate their invariant mass $m(e^+e^-)$. For $m(e^+e^-) < 85$ MeV, the event is to be discarded, unless we can prove by other means that this hypothesis is wrong. We expect a total of about 32 Dalitz decays in the entire experiment, and of these we expect 99% to have $m(e^+e^-) < 85$ MeV.[7] None of the candidates is discarded.

C. $K^0 \to \pi + \mu + \nu$, $\pi + e + \nu$, or $\pi + \pi + \gamma$ (1C fit). Candidate 1 720 440 has $\chi^2(\pi^+ + e^- + \nu) = 2.6$ and $\chi^2(\pi^+ + \pi^- + \pi^0) = 2.9$; candidate 1 739 122 has $\chi^2(\pi^+ + e^- + \nu) = 0.3$, $\chi^2(\pi^+ + \pi^- + \pi^0) = 7.2$. Gap counting on the negative track unambiguously proves these events both are $\pi^+ + e^- + \nu$. Candidate 1 492 562 has $\chi^2(\mu^+ + \pi^- + \nu) = 0.0$ and $\chi^2(\pi^+ + \pi^- + \pi^0) = 3.8$; gap counting does not resolve the ambiguity. From our 1C χ^2 distribution,[8] we bet 14 to 2 against $\pi^+ + \pi^- + \pi^0$. In addition, any completely ambiguous 3-body decay would have a priori about 5:1 odds against $\pi^+ + \pi^- + \pi^0$ on the basis of known branching ratios. We should perhaps count this as 1:35 event; instead we discard the candidate. Candidate 1 458 048 is nearly unmeasurable and is completely ambiguous. It should count as 0.2 events, but we discard it. None of the remaining 16 τ^0 candidates is ambiguous.

D. $K_1^0 \to \pi^+ + \pi^-$; π^+ (or π^-) suffers a small (unobserved) scatter. We delete the π^+ (or π^-) and fit (1C) to the hypothesis $K_1^0 \to \pi^+$ (unmeasured) + π^- (measured) (and also with the signs reversed). For $\chi^2(1C) < 10$, the event is a possible Coulomb scatter. However, we do not reject the event as a τ^0 candidate unless the "scattered" pion satisfies $(p\beta)_{\text{fitted}} |\theta_{\text{fitted}} - \theta_{\text{measured}}| \leq 2000$ (MeV/c) deg. This condition is chosen after considering the form and magnitude of the Rutherford scattering cross section, so that out of 3000 normal double vees, only a calculated

Table I. Details of the decays. χ_p^2 is for the production (3C), χ_d^2 is for the decay (1C); p_{K^0} (lab) is obtained from the production fit; t_{K^0} is the K_2^0 proper time from production to decay. T_+, T_-, and T_0 are the decay pion kinetic energies in the K_2^0 rest frame.

Event	χ_p^2	χ_d^2	p_{K^0} (lab) (MeV/c)	t_{K^0} (10^{-10} sec)	T_+ (MeV)	T_0 (MeV)	T_- (MeV)
554 595	1.3	0.6	354 ± 3	10.68	29.5	17.4	36.7
559 553	0.3	2.2	622 ± 7	6.94	36.0	36.4	11.2
760 339	3.9	0.0	527 ± 9	4.80	22.4	16.8	44.4
853 275	12.3	2.2	499 ± 8	5.88	25.0	37.7	20.9
1 357 217	5.4	1.7	637 ± 8	3.11	31.4	4.2	48.0
1 422 603	15.3	1.4	169 ± 3	1.95	41.6	11.0	31.0
1 424 205	5.7	2.3	590 ± 6	6.84	52.8	21.5	9.3
1 852 169	2.4	3.7	276 ± 3	7.80	13.3	34.6	35.7
1 852 302	3.5	4.7	605 ± 9	4.97	35.7	39.2	8.7
1 859 423	3.9	0.3	515 ± 6	0.48	51.4	19.8	12.4
1 519 387	3.1	5.2	640 ± 6	2.95	45.1	24.5	14.0
1 560 262	1.1	0.0	678 ± 6	2.35	43.4	9.5	30.7
647 539	0.2	0.3	856 ± 7	11.13	44.9	12.2	26.5
1 327 269	3.0	2.4	791 ± 9	0.65	28.2	20.7	34.7
1 839 341	8.5	0.3	81 ± 3	35.5	5.3	39.4	38.9
1 870 430	4.5	1.0	576 ± 9	14.1	49.0	30.8	3.8

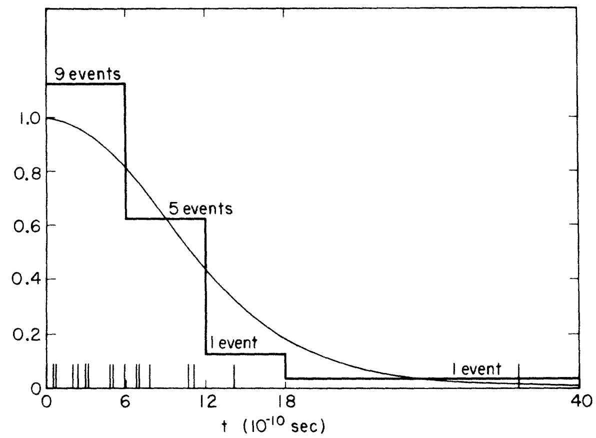

FIG. 1. Proper time distribution of $K_2^0 \to \pi^+ + \pi^- + \pi^0$ events. The vertical lines give the individual times of the 16 events. The smooth curve is their expected time distribution; its shape is entirely determined by the fiducial volume and the K^0 momentum distribution, because attenuation of the K_2^0 by decay is almost negligible.[10] The histogram of the 16 events has the same normalization as the smooth curve.

0.3 $K_1^0 \to \pi^+ + \pi^-$ decays followed by a single Coulomb scatter will fail to be cut off. This cutoff also eliminates $K_1^0 \to \pi^+ + \pi^-$ decays followed by a small-angle decay $\pi^\pm \to \mu^\pm + \nu$ in flight. Small-angle nuclear scatters are also eliminated, but in any case they are estimated to be negligible. None of the 16 τ^0 candidates are rejected by this cutoff.

We are left with 16 τ^0 events. There are no correction factors for lost or cutoff events. The characteristics of the events are exhibited in Table I. Their time distribution is shown in Fig. 1. There is no evidence for an enhanced τ^0 decay rate within the first K_1^0 mean life. This agrees with the expectation that the rate for $K_1^0 \to \pi^+ + \pi^- + \pi^0$ is negligible compared to that for $K_2^0 \to \pi^+ + \pi^- + \pi^0$.[9] We therefore use all of the τ^0 decays, including those within a few K_1^0 mean lives.

The rate $\Gamma_2(+-0)$ is equal to $1.014\ N(+-0)/T_2$, where T_2 is the total of the K_2^0 flight times in the fiducial volume, and $N(+-0) = 16$ is the number of observed τ^0 decays.[10] We obtain T_2 by using the 2608 acceptable normal double vees. The acceptance criteria for $\Lambda \to p + \pi^-$ are independent of the K decay mode. For a given acceptable Λ, the probability of obtaining an acceptable $K_1^0 \to \pi^+ + \pi^-$ decay and hence an acceptable double vee is given by $P \equiv (1/2)\ B[\exp(-t_0/\tau_1) - \exp(-t_1/\tau_1)]$. Here the factor 1/2 arises from $|K^0\rangle|^2 = (1/2)\ |K_1^0\rangle|^2 + (1/2)\ |K_2^0\rangle|^2$, B is $\Gamma(K_1^0 \to \pi^+ + \pi^-)/\Gamma(K_1^0 \to \text{all } 2\pi) = 0.725$,[11] τ_1 is 0.90×10^{-10} sec,[12] and t_0 and t_1 are the minimum and maximum acceptance times for $K_1^0 \to \pi^+ + \pi^-$.[13] Corresponding to an accepted normal double vee, the average number of acceptable Λ's is $1/P$, and their expected contribution to T_2 is $(1/2)(1/P)\ t_1$.[14] We then obtain T_2 by summing over all acceptable normal double vees, $T_2 \equiv \sum(1/2)(1/P)t_1 = B^{-1}\sum[\exp(-t_0/\tau_1) - \exp(-t_1/\tau_1)]^{-1}t_1$.[15] We find $BT_2 = 4.06 \times 10^{-6}$ sec. Our final result is $\Gamma_2(+-0) = (1.014)16/[4.06 \times 10^{-6}/0.725] = (2.90 \pm 0.72) \times 10^6\ \text{sec}^{-1}$.

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[1]R. F. Sawyer and K. C. Wali, Nuovo Cimento 17, 938 (1960); also, S. B. Treiman (private communication)

[2]An elementary spurion-type derivation of the result of Sawyer and Wali[1] is given by Frank S. Crawford, Jr., in Proceedings of the 1962 Varenna Summer School (Academic Press, New York, 1964), Course 26.

[3]Gideon Alexander, Silverio P. Almeida, and Frank S. Crawford, Jr., Phys. Rev. Letters 9, 68 (1962).

[4]Four of the present 16 decays $K_2^0 \to \pi^+ + \pi^- + \pi^0$ are included in the experiment of Alexander et al.[3] They obtained the direct result $\Gamma_2(+-0) = (2.66 \pm 1.34) \times 10^6\ \text{sec}^{-1}$, based on these four decays and a correspondingly smaller sample of associated production events $\pi^- + p \to \Lambda + K^0$. This agrees with the present result based on the complete sample of 16. Alexander et al. also obtained $\Gamma_2(+-0)$ indirectly, by combining their own direct absolute decay rate for $K_2^0 \to$ leptons, $\Gamma_2(L) = (9.31 \pm 2.49) \times 10^6\ \text{sec}^{-1}$, with the branching ratio $R \equiv \Gamma_2(+-0)/\Gamma_2(\text{all charged}) = 0.134 \pm 0.018$ of D. Luers, I. S. Mittra, W. J. Willis, and S. S. Yamamoto, Phys. Rev. Letters 7, 255 (1961). In this way they indirectly obtained $\Gamma_2(+-0) = (1.44 \pm 0.43) \times 10^6\ \text{sec}^{-1}$, in poor agreement with their own direct result, with the direct result of the present complete sample, and with the value predicted by the $|\Delta I| = 1/2$ rule. We can now turn around the above comparison. From our result given in Eq. (1) and, for example, the final result $R = 0.157 \pm 0.03$ of Luers et al. [Phys. Rev. 133, B1276 (1964)], we obtain the indirect result $\Gamma_2(L) = (15.6 \pm 5.3) \times 10^6\ \text{sec}^{-1}$. This agrees well with the prediction (see, for instance, reference 3) $\Gamma_2(L) = (16.5 \pm 1.2) \times 10^6\ \text{sec}^{-1}$ that follows from the $I = 1/2$ current rule for leptonic decays of kaons.

It is in rather poor agreement with the direct result $(9.31 \pm 2.49) \times 10^6$ sec^{-1} of Alexander et al.; however, the statistical errors are large. Moreover, the decay rate for $K_2^0 \rightarrow \pi^+ + \pi^- + \gamma$ is unknown, and has been taken to be zero in all published analyses. [It gives no background for the present result, Eq. (1).] We have not finished our analysis of the complete sample of K_2^0 (and K_1^0) leptonic decays corresponding to our present sample of $\pi^+\pi^-\pi^0$ decays; we postpone a more detailed discussion of $\Gamma_2(L)$ until we have obtained its direct determination from the complete sample. For a discussion of the results of other relevant experiments, see for example Luers et al., op. cit. (1964).

[5]The decay rate $\Gamma_2(+-0)$ can also be obtained indirectly by combining the K_2^0 lifetime τ_2 with the branching ratios $\nu = \Gamma_2(000)/\Gamma_2(\text{ch})$ and $\lambda = \Gamma_2(+-0)/\Gamma_2(\text{ch})$, under the assumption that there are no additional unobserved neutral modes, like $K_2^0 \rightarrow \gamma + \gamma$. R. H. Dalitz, Proceedings of the Conference on the Fundamental Aspects of Weak Interactions, Brookhaven National Laboratory, Upton, N.Y., September 1963 (unpublished), has combined the available data on τ_2, ν, and λ from eight different experiments, to obtain $\Gamma_2(+-0) = (1.93 \pm 0.35) \times 10^6$ sec^{-1}. This result differs by 2.3 standard deviations from the prediction of Eq. (3). In the notation of Eq. (4), it corresponds to $\sqrt{2}\,\text{Re}(A_{3/2}/A_{1/2}) = 0.12 \pm 0.06$.

[6]The 3C production χ^2 distribution for the final 16 τ^0 events is as follows: $\chi^2 = 0$ to 3.67, 8 events (we expect 11.20); 3.67 to 7.82, 5 events (we expect 4.00); 7.82 to 16.27, 3 events (we expect 0.78); >16.2, zero events (we expect 0.02). Thus the expected and observed χ^2 distributions are in excellent agreement. Since Λ production and decay occurs only about once in 30 pictures, there is only one chance in 900 of finding a 3-body K^0 decay with a possibly ambiguous origin. (We do not use single-vee 3-body K^0 decays.)

[7]N. P. Samios, Phys. Rev. 121, 275 (1961).

[8]The 1C decay χ^2 distribution for the 16 final τ^0 events is as follows: for $\chi^2 = 0$ to 1.07, 7 events (we expect 11.2); 1.07 to 3.84, 7 events (expect 4.00); 3.84 to 6.63, 2 events (expect 0.64); >6.63, zero events (expect 0.02).

[9]The decay $K_1^0 \rightarrow \pi^+ + \pi^- + \pi^0$ is forbidden for totally symmetric 3π states, and the contribution from nonsymmetric states is expected to be small because of angular-momentum barrier effects. See, for instance, S. Treiman and S. Weinberg, Phys. Rev. 116, 239 (1959).

[10]The mean decay distance for K_2 is large compared to the bubble chamber. A small correction factor of 1.104 arises from the total attenuation by decay of the K_2's. The attenuation by interaction in the hydrogen is even less important and is neglected.

[11]This is our weighted average of the compilation by M. Chretien, V. K. Fischer, H. R. Crouch, Jr., R. E. Lanou, Jr., J. T. Massimo, A. M. Shapiro, J. P. Averell, A. E. Brenner, D. R. Firth, L. G. Hyman, M. E. Law, R. H. Milburn, E. E. Ronat, K. Strauch, J. C. Street, J. J. Szymanski, L. Guerriero, I. A. Pless, L. Rosenson, and G. A. Salandin, Phys Rev. 131, 2208 (1963), Table IV.

[12]This is our weighted average of the compilation by Frank S. Crawford, Jr., in Proceedings of the 1962 International Conference on High-Energy Physics at CERN (Cern Scientific Information Service, Geneva, Switzerland, 1962), p. 839.

[13]In the film analyzed at Wisconsin $t_0 = 0$ was used. At Berkeley t_0 corresponded to a cutoff at 0.8 cm. The time t_1 is the potential proper time corresponding to the decay fiducial volume. The production fiducial volume is slightly smaller than the decay fiducial volume, so that large values of $1/P$ are excluded.

[14]We impose no t_0 cutoff for τ^0 decays.

[15]In reference 3, the procedure was to use all of the acceptable Λ decays, irrespective of whether there is an acceptable K_1^0 decay, and sum over the calculated potential K^0 times. In that case one need not use the value of B. However, B is extremely well known, so that the two methods are equivalent. This was verified by comparing the methods in the film analyzed at Berkeley (75% of the total).

TESTS FOR SPIN AND PARITY OF THE B MESON

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Experimental evidence has been recently reported of a π-ω resonance called B meson.[1-3] Different assignments for its spin s and parity P have been proposed.[4-7] Methods for determining s and P have been suggested[8,9] and preliminary results seem to favor $s^P = 1^-$.[10] We present here alternative tests for determining the quantum numbers and decay parameters of B.

These tests are similar to those previously discussed for fermions.[11,12] We consider the cascade decay of B: $B \rightarrow \pi + \omega$ followed by $\omega \rightarrow \pi^+ + \pi^- + \pi^0$. We denote by $\rho_{\mu\mu'}^{(B)}$ (with $-s \leq \mu$, $\mu' \leq s$) the elements of the density matrix of B, in the B rest frame, normalized to unit trace. The density matrix of ω is then given, in the ω rest frame, by

$$I(\vec{v})\rho^{(\omega)} = \sum_{\mu\mu'} \rho_{\mu\mu'}^{(B)} \sum_{ll'} T_l T_{l'}^* \sum_{mm'} \sum_{\nu\nu'} (lm, 1\nu|s\mu)(l'm', 1\nu'|s\mu') Y_l^m(\vec{v}) Y_{l'}^{m'*}(\vec{v}) |\nu\rangle\langle\nu'|, \qquad (1)$$