

## Comment on “Boson Peak in Deeply Cooled Confined Water: A Possible Way to Explore the Existence of the Liquid-to-Liquid Transition in Water”

In their Letter, Wang *et al.* [1] report on an inelastic neutron scattering (INS) experiment where they describe the pressure evolution of a low energy ( $E \sim 6$  meV) excitation, emerging in confined protonated water only below 230 K at an exchanged momentum  $Q = 2.0 \text{ \AA}^{-1}$ . Water confinement was used to overcome the unavoidable crystallization occurring below  $\sim 250$  K in bulk water. The authors report that a similar finding was also obtained in both bulk (numerical simulations [2]) and confined water at ambient pressure. They refer to this low temperature excitation as a boson peak (BP) [3], and relate its occurrence to the Widom line, concluding that the observed pressure behavior of the BP reveals the signature of the high-density liquid (HDL) to the low-density liquid (LDL) transition proposed [4], though severely questioned [5], for bulk water. We believe these claims to be unconvincing for the following reasons.

*Comparison with corresponding findings in liquid water.*—The authors not only overlook commenting on the actual density of confined liquid water [6–9], but they also neglect to establish any physical relationship with the well known excitations of coherent or incoherent origin occurring at similar energies in bulk liquid water. Since the seminal Raman scattering room temperature studies by Bolla [10], a mode in the  $\sim 5$ – $7$  meV range has indeed been regularly observed with optical [11–14], numerical [15], and inelastic x-ray scattering [16–18] and INS techniques [19,20] over a wide thermodynamic range (0–2 kbar, 250–450 K) in both  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  (see Fig. 4 in Ref. [21]). The microscopic nature of such a mode, underdamped and still well defined at  $Q = 2.0 \text{ \AA}^{-1}$ , is the subject of controversial single-particle [12,22,23] or collective [21,24] interpretations. Irrespective of its incoherent or coherent nature, this evidence is unquestionable and cannot be ignored. This mode is not easily detectable in high temperature neutron spectra from  $\text{H}_2\text{O}$  because of the overwhelming quasielastic contribution. However, its presence always emerges in calculating the hydrogen vibrational density of states, as was done in Ref. [22] at  $T = 256$  K, and in bulk or confined  $\text{H}_2\text{O}$  from 300 K down to 242 K [25], but not mentioned in Ref. [1]. This mode, but not the BP, was also observed when investigating the vibrational dynamics in amorphous ices [26,27]. Moreover, a bulklike excitation not dependent on temperature was observed down to 205 K in an INS measurement on slightly salty liquid water [28].

*Data analysis and treatment.*—(i) INS probes at the same time the coherent and incoherent properties of matter with a weight given by their respective neutron cross section and dynamic structure factor.  $\text{H}_2\text{O}$  is considered as an incoherent scatterer by reason of the high  $\sigma_{\text{inc}}/\sigma_{\text{coh}}$  ratio. Yet, this approximation cannot be uncritically adopted as was done

in Ref. [1] and a proper estimation of the related ratio  $S_{\text{inc}}(Q, \omega)/S_{\text{coh}}(Q, \omega)$  at the thermodynamic ( $P, T$ ) and kinetic ( $Q, \omega$ ) investigated point should be addressed. (ii) An arbitrary interpolating metric is adopted to determine the locus of the BP appearance: the  $T_B$  parameter is a clumsy, large-error quantity inherent in the slowing down of the thermal diffusion. The peak associated with the low energy—and virtually temperature independent (see Fig. 4 of Ref. [1])—excitation is enhanced by the narrowing of the quasielastic signal upon lowering the temperature. (iii) The exact internal pressure existing in such tiny pores is not directly related to the He applied pressure and is therefore unknown [29]. As a consequence, the confined water phase diagram and properties cannot be unconditionally assigned to those of bulk water. (iv) In order to support the authors’ claims at a less speculative level, the correct BP shape should be determined by calculating the vibrational density of states in excess of that of the corresponding crystal-line phase.

In conclusion, the whole large body of numeric and experimental investigations on the single particle and collective properties of liquid water report the presence of a weakly dispersing excitation in the 5–7 meV range. We believe that, in order to use the BP as a marker of the HDL or LDL bulk water phases, the authors should perform a more complete data treatment and establish a relation, if any, between the supposed BP peak they observe in a confined environment and the well established bulk mode present in a wide portion of the phase diagram at the same energy.

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Received 13 March 2015; published 30 September 2015

DOI: 10.1103/PhysRevLett.115.149801

PACS numbers: 61.20.Lc, 63.50.-x

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