

Erratum: Optical Spectrum of MoS₂: Many-Body Effects and Diversity of Exciton States [Phys. Rev. Lett. 111, 216805 (2013)]

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After further convergence tests, we find that finer k grids than the $12 \times 12 \times 1$ k grid for the dielectric matrix and the $72 \times 72 \times 1$ k grid for the BSE Hamiltonian reported in our Letter are required to converge the calculation. We have increased k -point sampling for the calculation of quasiparticle energies to $24 \times 24 \times 1$ and the calculation of the optical spectrum to $300 \times 300 \times 1$. We find that the quasiparticle gap decreases to 2.7 eV and the binding energy of the lowest excitation energy A exciton decreases to 0.65 eV. The main conclusions of our Letter—that MoS₂ has a large number of exciton states with large binding energies and the analysis of the character of these states—remain unchanged, but the oscillator strengths of excited states of the A and B excitons in the optical spectrum have decreased, resulting in better agreement with experiment for higher energy states. The $300 \times 300 \times 1$ sampling in our corrected calculation was achieved by explicitly calculating matrix elements on a $24 \times 24 \times 1$ k grid for all transitions in the Brillouin zone and a $300 \times 300 \times 1$ k grid for all transitions where the scattering momentum is less than 0.1 bohr^{-1} . These explicitly calculated

TABLE I. Values from different studies: optical transition energies for peaks A , B , A' , B' , and C , the binding energy of peak A (E_b) and the convergence parameters— k -point sampling and number of valence (N_v) and conduction (N_c) bands. Low temperature experimental results were taken at 10 K [1] and 4 K [2].

| | Optical transition energies (eV) | | | | | E_b | Convergence parameters | | |
|------------------------------------|----------------------------------|------|----------|------|------|-------|---------------------------|-------|-------|
| | A | B | A' | B' | C | | k grid | N_v | N_c |
| Present work (G_0W_0) at 0 K | 2.04 | 2.17 | 2.32 | 2.45 | 2.73 | 0.63 | $300 \times 300 \times 1$ | 4 | 4 |
| Present work (G_0W_0) at 300 K | 1.96 | 2.09 | 2.24 | 2.37 | 2.69 | 0.63 | $300 \times 300 \times 1$ | 4 | 4 |
| Absorp. exp. at low T | 1.92–1.94 [1,2] | | 2.08 [1] | | | | | | |
| Absorp. exp. at 300 K | 1.88 [3] | | 2.03 [3] | | | | | | |
| PL exp. at 300 K | 1.85 [4] | | 1.98 [4] | | | | | | |
| G_0W_0 -BSE [5] | 1.78 | 1.96 | | | | 1.04 | $6 \times 6 \times 1$ | 2 | 4 |
| sc- GW -BSE [6] ^a | 2.22 | 2.22 | 2.5 | 2.5 | | 0.63 | $15 \times 15 \times 1$ | 6 | 8 |

^aSpin orbit was not included in this calculation.

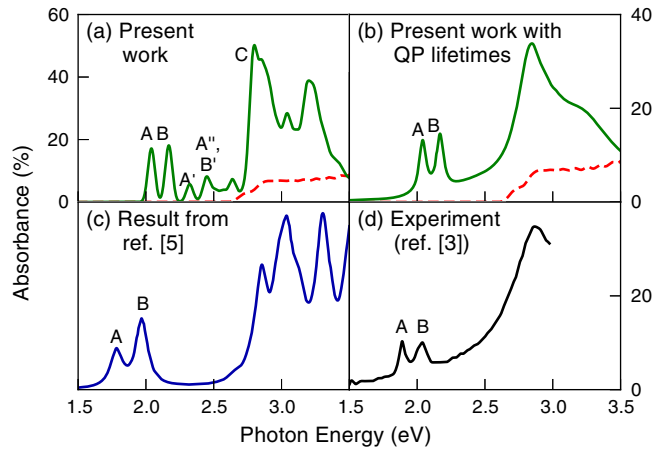


FIG. 1 (color online). (a) Absorption spectra of MoS₂ without (dashed red curve) and with (solid green curve) electron-hole interactions using a constant broadening of 25 meV. (b) Same calculated data as in (a) but using an *ab initio* broadening based on the electron-phonon interactions [7,8]. (c) Previous G_0W_0 plus BSE calculation (in arbitrary units) [5]. (d) Experimental absorbance [3].

matrix elements are interpolated to all points on a uniform $300 \times 300 \times 1$ k grid. This fine sampling is necessary to capture fast variations in screening at small wave vectors and fine features in the exciton wave functions, which are tightly localized in k space.

We have summarized the corrected quantitative results in Table I, which replaces Table II on page 3 of our Letter. In the table, we have added the effect of lattice expansion due to temperature to draw a more accurate comparison with experimental measurements taken at room temperature, as noted by Soklaski *et al.* [2]. Figure 1 replaces Fig. 2 on page 3 of our Letter.

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