

Dynamics and Scission of Rodlike Cationic Surfactant Micelles in Shear Flow

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Flow-induced configuration dynamics and scission of rodlike micelles are studied for the first time using molecular dynamics simulations in the presence of an explicit solvent and salt. Predicted dependence of tumbling frequency and orientation distribution on shear rate S agrees with mesoscopic theories. However, micelle stretching increases the distance between the cationic head groups and adsorbed counterions, which reduces electrostatic screening and increases the overall energy Φ linearly with micelle length. Micelle scission occurs when Φ exceeds a threshold value, independent of S .

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Thermodynamic self-assembly in surfactant solutions is known to result in various micelle morphologies such as spheres, cylinders, vesicles, and lamellae. The rich diversity in equilibrium morphologies is of great fundamental and practical interest in applications ranging from detergency to targeted drug delivery. Amongst micellar structures, cylindrical and wormlike micelles are known to exhibit rich rheological behavior because micelle length and entanglement density depend greatly on the surfactant and salt concentration, salt hydrophobicity, temperature, and flow shear. Specifically, depending on the surfactant concentration, chemical environment, and shear rate, cylindrical micelle solutions could exhibit Newtonian, shear thinning, or shear thickening behavior, which is accompanied by shear-induced structure (SIS) formation or shear banding [1–7].

While phenomenological, kinetic-theory inspired [8] and continuum-level constitutive models [9–12] have been developed to study the dynamics and rheology of cylindrical micelles, molecular-level explorations of flow-structure coupling in micelle solutions is lacking. Understanding micelle structure and dynamics via faithful and systematic molecular simulations is an important and essential step in filling this knowledge gap [13]. Briels *et al.* [14] have developed stochastic simulations (Brownian dynamics), in which parameters such as the persistence length were obtained from atomistic molecular simulations, to study the dynamics and rheology of wormlike micelle solutions. Here, we describe for the first time insights gained from a coarse-grained molecular dynamics (CGMD) study of cylindrical cationic micelles subjected to uniform and steady shear flow. In CGMD, explicit solvent-salt-micelle interactions are incorporated by utilizing force fields which are validated against more detailed atomistic molecular dynamics (MD) simulations [15,16]. Several fundamental questions are addressed including the criterion that demarcates the Brownian (diffusive) and

flow-aligned regimes, asymptotical scaling laws that characterize micelle configurational dynamics and their comparisons with stochastic theories of flexible and semi-flexible polymers [17–20], relationship between micelle energetics and length, and a mechanism of flow-induced micelle scission.

Cates and co-workers [3] suggested that SIS formation in rodlike micelle solutions is initiated by energetic end-to-end collisions of flow-aligned micellar rods resulting in micelle fusion, which, in turn, leads to the formation of longer micelles. Significant flow alignment is induced when the rotational diffusion time scale $\tau_D \geq 1/S$ where S is the shear rate, or equivalently, when the Peclet number $Pe \equiv S\tau_D \geq 1$. Longer micelles tend to flow-align more easily because $\tau_D \sim L^3$ where L is the rod length [21]. Hence, this mechanism results in the formation of progressively longer micelles which could eventually form networks. Recent equilibrium MD simulations of cationic spherical micelles suggest that sufficient electrostatic screening is required to induce micelle fusion, which is more readily facilitated when ringed organic counter ions are present in salt to surfactant molar ratio exceeding unity (see Fig. 5 in [16]). For instance, benzoate ions physically adsorb onto the micelle-water interface and help neutralize the repulsive electrostatic interactions between the surfactant head groups, as characterized by the potential of mean force of binary interactions. This is consistent with experimental observations that hydrophobic salts such as benzoate or salicylate promote SIS formation and shear thickening in cationic micelle solutions [22].

Hence, the literature suggests that satisfaction of two criteria is required for robustly inducing SIS in ionic micelle solutions, namely, a mechanical one that requires that the inverse rate of flow deformation is comparable to that of the inherent rotational time scale of the micelle, and a chemical one that necessitates the counterion concentration to be large enough to provide sufficient electrostatic

screening. In this Letter, we focus on studying the effects of shear flow on the configurational dynamics of a single cylindrical cetyl trimethylammonium chloride (CTAC) micelle in an aqueous solution containing sodium salicylate (NaSal) salt while preserving the underlying physical chemistry and intrinsic self-assembled structure of the micelle. In the simulations, all chemical interactions, spontaneous rearrangements within the micelle structure, and micelle shape or length fluctuations and angular diffusion are explicitly calculated. To date, atomistic non-equilibrium molecular dynamics (NEMD) simulations of micellar systems in the presence of explicit solvent and electrostatic interactions are limited to ns time scales. In the present study, we are able to conduct NEMD studies that span microseconds (\gg angular orientational relaxation time τ_r) by utilizing coarse-grained (CG) potentials of CTAC surfactant and NaSal salt in explicit solvent. Coarse graining of surfactant, salt, and water helps reduce the number of particles by approximately 4 times and increases the time step in MD by an order of magnitude since time scales on the order of atomic bond vibrations are neglected in the CG description. MARTINI-based CG force fields [23] used here have been extensively validated against atomistic simulations in terms of the radial distribution functions of the system constituents [15]. Polarizable CG water was used as a solvent because of its ability to capture electrostatic interactions in systems with charged interfaces more accurately compared to CG water models that lack a local dipole [24]. The integrity of the cylindrical structure of the micelle is maintained by the surface adsorption of Sal^- ions onto the micelle-water interface, especially when the molar ratio of NaSal to CTAC, $R \geq 1$. Hence, in the simulations we used $R = 1.5$. The methodology of assembling and equilibrating cylindrical micelle structure, implementation of shear flow and boundary conditions in the MD simulations, determination of micelle relaxation from flow-induced stretched configurations, and the calculation of the instantaneous micelle length and orientation are provided in the Supplemental Material [25]. The Weissenberg number $Wi \equiv S\tau_r$ is used to characterize the flow strength. Simulations are conducted for $0 \leq Wi \leq 125$. The vector \mathbf{l} that connects the centers of mass of the two end caps of the micelle represents the micelle axis. The angle θ that the projection of \mathbf{l} onto the flow (x) gradient (y) plane subtends with the x direction is used to quantify orientation and flow alignment. Micelle structure is not subject to any external constraints and is governed only by hydrodynamic forces as well as the interactions among the CTAC surfactants, sodium salicylate salt (NaSal), Cl^- counterions, and water molecules.

Figure 1(a) shows a schematic of the NEMD simulation box and Figs. 1(b)–1(e) show θ plotted as a function of time t for four different values of Wi . Note that $\theta = 0^\circ$ represents a fully flow-aligned state while a transition through 90° corresponds to a tumbling event. Results are shown only for

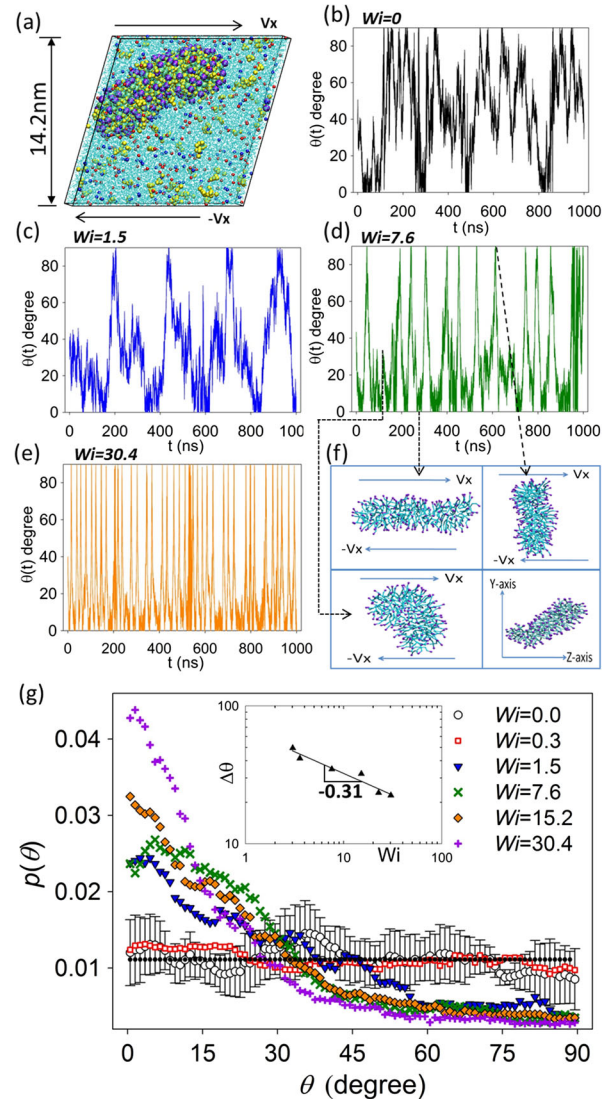


FIG. 1 (color online). (a) Schematic of NEMD simulations (high resolution image in SI). CTA^+ surfactants (purple head and green tails), Sal^- ions (yellow), Na^+ (red) and Cl^- (blue) ions and water (cyan). (b)–(e) Plots of θ versus t at for various Wi . (f) (i) Flow-aligned, (ii) gradient-aligned, (iii) misaligned (in the xy plane), and (iv) misaligned (in the yz plane) states (not to scale). (g) Probability distribution function $p(\theta)$ at various Wi . The dotted line represents a uniform distribution and the vertical bars show the standard deviation of $p(\theta)$ at $Wi = 0$ in the simulations within each interval. Inset shows a plot of $\Delta\theta$ as a function of Wi showing $\Delta\theta \propto Wi^{-0.31}$.

the first $1 \mu\text{s}$ for clarity although the trends remain the same throughout the duration of the simulations ($3 \mu\text{s}$). It is evident that, in the absence of flow, i.e., for $Wi = 0$, the intrinsic rotational diffusion of the micelle dominates the orientational distribution. This is clearly seen in Fig. 1(g) where the probability distribution function (PDF) $p(\theta)$ is plotted. The PDF is uniform for $Wi = 0$ suggesting the absence of any preferred orientation. As Wi is increased, the effect of flow on micelle orientation becomes

progressively more prominent. For example, for $Wi = 0.3$ the departure of $p(\theta)$ from the uniform equilibrium distribution is relatively small, suggesting that the rotational diffusion dominates over flow alignment. However, at $Wi = 1.5$, higher probabilities are clearly seen associated with flow-aligned states. A crossover from a diffusive to a flow-aligned regime occurs at $Wi \sim 1$. For $Wi > 1$, quasi-periodic tumbling events dominate the micelle dynamics. Snapshots (i) and (ii) of the micelle displayed in Fig. 1(f) correspond to $\theta = 0^\circ$ (flow-aligned state) and $\theta = 90^\circ$ (gradient-aligned state), respectively, at $Wi = 7.6$. The tumbling events occur at regular intervals while the tumbling frequency increases with Wi . For $Wi > 1$, $p(\theta)$ follows a Gaussian distribution and the width at half maximum of the orientational distribution $(\Delta\theta) \propto Wi^{-0.31}$, which is consistent with previous predictions of mesoscopic theories for semiflexible polymers [20,26,27]. Shape fluctuations in the z dimension also influence the tumbling cycles, as seen in snapshot (iii) in Fig. 1(f), in which the micelle is slightly bent and misaligned with respect to the flow direction (more snapshots are shown in [25]). In this case, the micelle axis is out of the flow-gradient plane, as seen in snapshot (iv). Such events occur more frequently for relatively low shear rates and tend to randomize the tumbling cycles. The micelle eventually recovers from such misaligned states to flow-aligned ones before undergoing a tumbling event or transitioning to another misaligned state.

The autocorrelation function (ACF) of θ defines the temporal correlation of the angular orientation. The ACF at $Wi = 0$ (at equilibrium) can be used to define an orientational relaxation time τ_r of θ by fitting it to a decaying exponential of the form

$$\text{ACF} = \exp(-t/\tau_r), \quad (1)$$

where the ACF has been normalized using its value at $t = 0$. The relaxation time τ_r was calculated to be 28.7 ns and was used in our calculations of Wi . The power spectral density can be obtained from the fast Fourier transform of the ACF. We obtain the tumbling frequency $f^* \equiv (1/\tau_c)$ from the peak in the power spectral density [28]. The normalized tumbling frequency $f \equiv \{1/[(\tau_c/\tau_r)]\}$, plotted as a function of Wi , is shown in Fig. 2. For $Wi < 1$, no systematic change in f is observed. However, for $Wi > 1$, f increases, indicating more frequent and periodic tumbling. Power law fitting of f versus Wi provides an exponent of 0.68 which is also observed in Brownian dynamics simulations of polymers in dilute solutions [29–32]. The rotational diffusivity D was obtained from equilibrium simulations as one sixth of the slope of the linear fit of the mean squared angular displacement versus time [25]. This yielded $D = 6 \times 10^6 \text{ rad}^2 \text{ s}^{-1}$, corresponding to $\tau_D = 6.58 \mu\text{s}$ which is in good agreement with theoretical predictions for rodlike polymers [21]. However, since

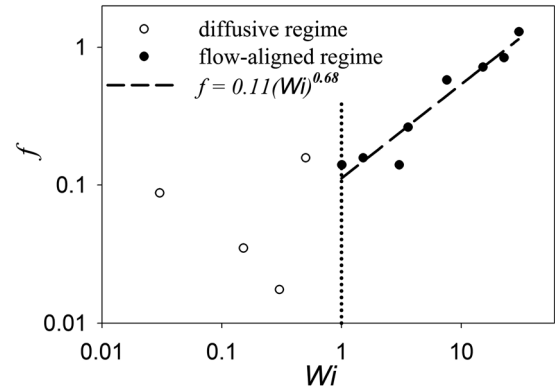


FIG. 2. Normalized tumbling frequency f versus Wi . Open circles indicate stochastic diffusion regime while filled circles lie in the flow-aligned regime. Power law fitting in the flow-aligned regime yields $f \sim Wi^{0.68}$.

temporal correlations in θ are lost during micelle tumbling, τ_r is much smaller and comparable to the tumbling period. Transition from the diffusive to flow-aligned regime occurs when $\tau_r \sim 1/S$, or equivalently, when $Wi \approx 1$ ($Pe \approx 230$).

We computed the change in the total pair potential energy $\Phi(l)$ of the surfactants in the micelle and adsorbed Sal^- ions as a function of length l of the micelle, defined as the magnitude of the vector \mathbf{l} . Note that l , therefore, is smaller than the end-to-end distance of the micelle. The use of \mathbf{l} to calculate l , instead of using the end-to-end micelle vector, has an advantage of reducing the statistical uncertainty in l arising from shape fluctuations. We carried out 10 different equilibrium MD simulations of previously sheared systems where flow-aligned stretched micelles were allowed to relax to their equilibrium length structures. Over the length relaxation spectrum of the micelle from about $l = 7.2$ to 4.7 nm, $\Phi^*(l) \equiv \Phi(l) - \langle \Phi(l, Wi = 0) \rangle$ was averaged into equally spaced bins of l with a width $dl = 0.0122$ nm [Fig. 3(a)]. The relaxation of length happens over a period of roughly 8 ns, and therefore, only the first 8 ns of data were considered. Energy $\Phi^*(l)$ needed to stretch a micelle from its equilibrium length is observed to vary linearly with l [Fig. 3(a)]. The slope of the linear fit provides a constant stretching force of $\approx 650 \text{ kJ mol}^{-1} \text{ nm}^{-1}$ which corresponds to ~ 1 nN per molecule. In comparison, the force required to stretch a covalent bond is 2–3 orders of magnitude greater [33], while that for hydrogen bond stretching in water is 1–2 orders of magnitude greater [34]. This is reasonable considering that the micelle is self-assembled by weak nonbonded forces. Our estimate is in very good agreement with experimental measurements of similar forces in biological systems [35].

The linear relationship between $\Phi^*(l)$ versus l may be interpreted as an asymptotic behavior valid for a relatively short ($\sim \text{nm}$), stiff rodlike molecular assembly. From a further coarse grained, mesoscopic modeling point of view, this would suggest that a micelle with length on the order of

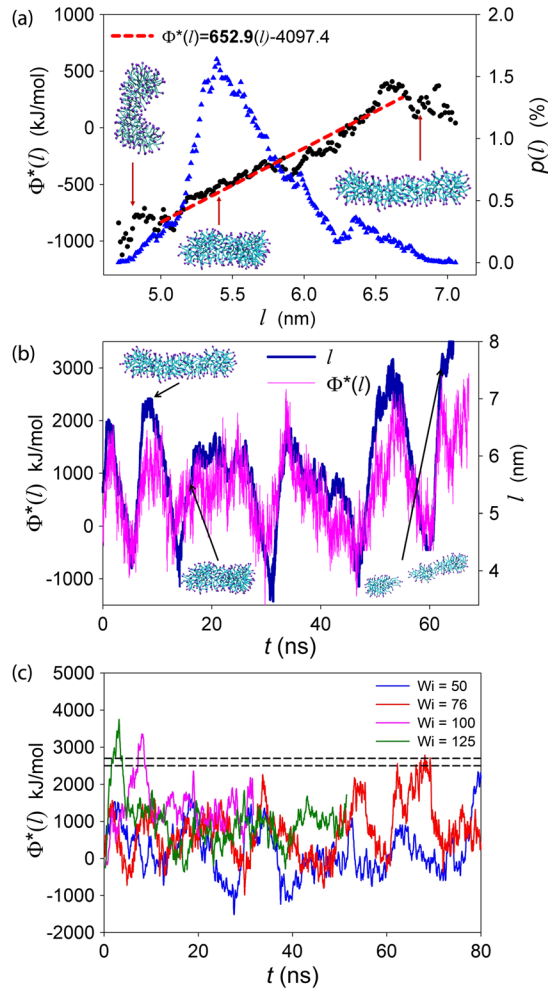


FIG. 3 (color online). (a) Pair potential energy $\Phi^*(l)$ (black dots) and probability distribution of length $p(l)$ (blue triangles) from the length relaxation simulations are plotted as a function of l . (b) $\Phi^*(l)$ and l versus t showing micelle scission for $Wi = 76$. Snapshots at different locations show (1) stretched micelle, (2) micelle with most probable l , and (3) micelle after scission. (c) $\Phi^*(l)$ versus t for four different simulations. Dotted lines indicate the range for micelle scission for $Wi > 70$.

a few nanometers may be considered as a stiff spring whose length may only fluctuate within a small range around an average value, e.g., as in the stiff FENE-Fraenkel spring proposed by Hsieh *et al.* [36] to model practically inextensible rods in the freely jointed bead-rod chain model of polymer chains. In order to verify this, we calculated the probability distribution of l from a total of 160 000 configurations. A probability distribution of l during the relaxation process is also plotted in Fig. 3(a). The distribution shows a pronounced peak at 5.4 nm which is the most probable length. The micelle appears to explore lengths smaller than $l = 5.4$ nm which is possible because of a mechanism in which the recoil from a stretched state to equilibrium length would cause either micelle bending or compression. In the three different micelle states shown in

the plot, a stretched micelle ($l \approx 7.0$ nm) can be seen to relax to a shorter state with $l \approx 5.4$ nm and then further recoil into a state bent along its axis ($l < 5.0$ nm).

Unlike polymer chains, in which the monomers are covalently bonded to each other, micelles are self-assembled by weak nonbonded interactions. This presents a limitation in the range of Wi that can be explored without causing micelle scission. We performed multiple simulations between $50 \leq Wi \leq 125$ to study micelle scission. It was observed that micelle scission does not occur for $Wi < 70$. For larger Wi , flow shear is sufficiently strong to cause micelle breakage as shown in Fig. 3(b), in which the micelle is seen to undergo a few tumbling events before scission. A plot of $\Phi^*(l)$ as a function of time is shown in Fig. 3(c) for a case where no breakage was observed ($Wi = 50$) and for three cases for which the micelle broke apart into two shorter ones ($Wi = 76$ 100 125). This plot, as well as data from other simulations, suggests the existence of an energy threshold for scission. From a series of simulations, the threshold value was estimated to be ~ 2500 – 2700 kJ mol $^{-1}$. The observed extension before breakage, compared to the equilibrium configuration, is 2–3 nm. This, based on the stretching force estimated above, corresponds to an increase in $\Phi^*(l)$ of 1300–1950 kJ mol $^{-1}$, which is comparable to the estimated threshold. As the micelle stretches, the distance between the CTA $^+$ head group and adsorbed Sal $^-$ ions increases, thereby reducing electrostatic screening and increasing the overall micelle energy. Above the threshold, it is energetically favorable to have two smaller micelles. This is evidenced further by the fact that the magnitudes of $\Phi^*(l)$ and $\Phi_{EL}^*(l)$, which is the electrostatic component of the pair potential energy between CTA $^+$ and Sal $^-$, are approximately equal. The other components of $\Phi^*(l)$ remain practically unchanged.

In summary, we have provided a quantitative description of the shear-induced orientation dynamics, stretching, and scission of rodlike surfactant micelles. An appropriate parameter that demarcates the diffusive and flow aligned regimes is the ratio of the micelle angular relaxation time to the inverse shear rate. Analyses of tumbling frequencies, orientation distributions, and energy-extension relationships suggest that from a mesoscopic point of view, short rodlike micelles may be represented by a stiff bead-spring unit subject to hydrodynamic drag, Brownian fluctuations, and a constant entropic spring force. Micelle scission happens through a mechanism in which shear-induced stretching causes the surfactant head groups and adsorbed counterions to be farther apart, resulting in reduced electrostatic screening and an increase in the overall micelle energy. Thus, consistent with experiments [e.g., Fig. 1(a) in [4]], for $Wi \sim O(1)$ flow-alignment of micelles would promote SIS formation leading to shear thickening, while for $Wi \gg 1$, flow-induced micelle scission would result in shear thinning. The CGMD methodology used here can be

extended to solutions that contain multiple micelles and/or nanoparticles for studying emerging morphologies, flow-structure interactions, and rheological properties to benefit applications ranging from enhanced oil recovery and hydrofracking [37] to the design of active nanofluids for energy harvesting and sensing [38].

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