

Comment on “Tuning the Magnetic Dimensionality by Charge Ordering in the Molecular TMTTF Salts”

Yoshimi *et al.* [1] have attempted to explain the pressure (P)-dependent behavior of Fabre salts which exhibit charge order (CO), antiferromagnetic (AFM), and spin-Peierls (SP) phases. Experiments find two AFM phases [2,3], AFM₁ at large P and AFM₂ at small P . Yoshimi *et al.* suggest that there also exist two distinct zero-temperature SP phases, SP₁ and SP₂. Here we point out that the occurrence of two distinct SP phases contradicts experiments [2,3], and is found in Ref. [1] because of unrealistic model parameters.

The experiments of Refs. [2,3] emphasize *cooperative interaction* between the ferroelectric charge order (FCO) and AFM₂ phases. In the experimental phase diagram [2,3] T_{CO} and the Néel temperature in the AFM₂ phase both decrease with P . Thus charge occupancies in the FCO and AFM₂ phases are likely the same. In contrast, P increases [2,3] the SP transition temperature, indicating that FCO and SP₂ phases *compete*. No CO was detected for $P > 0.5$ GPa in (TMTTF)₂SbF₆ [2,3], in the P region where the SP₂ phase occurs at lower temperature. It is then unlikely that SP₂ and FCO coexist at zero temperature.

The hopping parameters used by the authors in their model calculations are realistic. Their choice of Coulomb interactions is, however, unrealistic. The onsite Coulomb interaction assumed, $U/t_{a2} = 4$, is too small—in the purely electronic one-dimensional model no CO occurs for this U [4,5]. The assumed intersite Coulomb interactions $V_b = 0$ and $V_q = V_a$, are also unrealistic. Given the lattice geometry (see Fig. 1) it is highly unlikely that $V_b \ll V_q$, and with large interchain separation $V_q = V_a$ is equally unrealistic. $4 \lesssim U \lesssim 8$ and $V_b = V_q \ll V_a$ is more appropriate.

We repeated the calculations with more realistic $V_a = V$, $V_b = V_q = 0$, and $4 \leq U \leq 8$. For these parameters, the

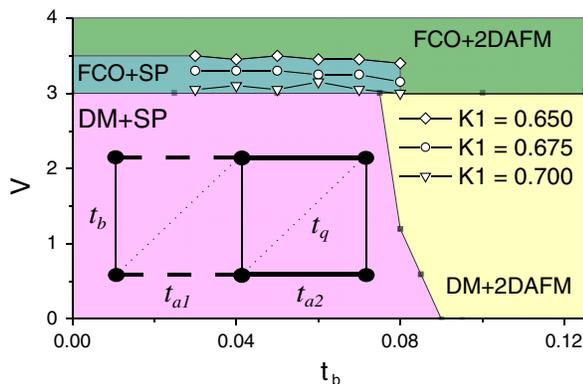


FIG. 1 (color online). 8×2 phase diagram for $U = 6$, $V_a = V$, and $K_2 = 1$. The inset shows the lattice structure assumed by Ref. [1]. As K_1 increases, the size of the FCO + SP phase shrinks. Other points do not significantly change with K_1 .

intradimer charge structure factor ($C_-(\mathbf{q})$ in Ref. [1]) peaks at several \mathbf{q} values, indicating comparable energies for both FCO and the checkerboard pattern CO, in agreement with experiments [6]. Peaks in $S_{\pm}(\mathbf{q})$ remain at the same \mathbf{q} values as in Fig. 2 of [1]. We conclude that the V_{ij} assumed in Ref. [1] is *not* required to explain coexisting FCO and AFM order in the AFM₂ state.

We also repeated (see Fig. 1) the 8×2 calculations with these parameters. We have three main observations. (i) For $V_a = V$, $V_b = V_q = 0$, we find a phase diagram similar to that in Ref. [1], but with the FCO entering at larger V as expected [4,5]. The choice $V_q = V$, $V_b = 0$ is also *not* required to realize the FCO phase; FCO can be stabilized by antiferromagnetic superexchange along the t_b bonds. (ii) As U increases the FCO + SP phase narrows. (iii) For both these and the parameters assumed in Ref. [1], the width of the FCO + SP phase is directly proportional to the strength of the intersite electron phonon coupling (larger K_1 gives weaker coupling). Unconditional transitions in the thermodynamic limit occur in the limit of 0^+ phonon coupling. Importantly, point (iii) was not discussed in Ref. [1], and together with (ii) suggests that in the thermodynamic limit the FCO + 2DAFM and DM + SP phases may share a common border.

To understand the phase diagram one must consider thermodynamics. For large Coulomb interactions the free energy is dominated by spin excitations. We have previously shown that the same DM + SP ground state can have two kinds of soliton spin excitations: (i) with local CO, or (ii) with uniform charge but local bond distortion [7]. In this picture, to the left of the line bisecting the SP phase [2], soliton excitations with local CO dominate at finite T ; to the right occur excitations with uniform site charges. A unique SP *ground state* is expected at all pressures between AFM₁ and AFM₂.

We acknowledge support from the Department of Energy Grant No. DE-FG02-06ER46315.

A. B. Ward,¹ R. T. Clay¹ and S. Mazumdar²
¹Department of Physics & Astronomy and HPC²
 Center for Computational Sciences
 Mississippi State University
 Mississippi State
 Mississippi 39762-5167, USA
²Department of Physics
 University of Arizona
 Tucson, Arizona 85721, USA

Received 16 December 2013; published 11 July 2014

DOI: 10.1103/PhysRevLett.113.029701

PACS numbers: 71.10.Fd, 71.20.Rv, 71.30.+h, 75.30.Kz

[1] K. Yoshimi, H. Seo, S. Ishibashi, and S. E. Brown, *Phys. Rev. Lett.* **108**, 096402 (2012).

-
- [2] W. Yu, F. Zhang, F. Zamborszky, B. Alavi, A. Baur, C. A. Merlic, and S. E. Brown, *Phys. Rev. B* **70**, 121101 (2004).
- [3] F. Iwase, K. Sugiura, K. Furukawa, and T. Nakamura, *Phys. Rev. B* **84**, 115140 (2011).
- [4] R. T. Clay, S. Mazumdar, and D. K. Campbell, *Phys. Rev. B* **67**, 115121 (2003).
- [5] H. Seo, J. Merino, H. Yoshioka, and M. Ogata, *J. Phys. Soc. Jpn.* **75**, 051009 (2006).
- [6] T. Nakamura, *J. Phys. Soc. Jpn.* **72**, 213 (2003).
- [7] R. T. Clay, R. P. Hardikar, and S. Mazumdar, *Phys. Rev. B* **76**, 205118 (2007).