## Comment on "Strangeness -2 Hypertriton"

In a recent Letter, Garcilazo and Valcarce [1] reported on a  $\Lambda\Lambda N - \Xi NN$  coupled-channel three-body Faddeev calculation that binds  ${}_{\Lambda\Lambda}{}^3$ n and  ${}_{\Lambda\Lambda}{}^3$ H by about 0.5 MeV below the corresponding  $\Lambda\Lambda N$  thresholds. This contrasts with *ab initio*  $A \leq 6$  few-body coupled-channel calculations associating a loosely bound  ${}_{\Lambda\Lambda}{}^4$ H with the onset of  $\Lambda\Lambda$  hypernuclear binding [2]. Here, I argue that the S = -2 chiral constituent quark model (CCQM) interactions [3] that bind  ${}_{\Lambda\Lambda}{}^3$ H [1], as well as the unobserved *H* dibaryon [4], overbind the uniquely identified NAGARA emulsion event of  ${}_{\Lambda\Lambda}{}^6$ He [5] by more than 4 MeV, casting doubts on the predictive power of the CCQM for S = -2.

Listed in Table I are  $\Delta B_{\Lambda\Lambda}({}^{6}_{\Lambda\Lambda}\text{He})$  values obtained in two sets of  $\alpha \Lambda \Lambda$  three-body calculations [6,7] which use identical  $V_{\Lambda\alpha}$ ; the  $V_{\Lambda\Lambda}$  from Ref. [7] are softer than the  $V_{\Lambda\Lambda}$  from Ref. [6]. Within each set,  $\Delta B_{\Lambda\Lambda}$  increases with increasing the strength of  $V_{\Lambda\Lambda}$ , as represented by the listed values of  $-a_{\Lambda\Lambda}$ . For  $a_{\Lambda\Lambda}^{CCQM} = -3.3$  fm, corresponding to the decoupled  $V_{\Lambda\Lambda}^{CCQM}$  [9], interpolation within the first set [6] suggests that  $\Delta B_{\Lambda\Lambda}^{CCQM}({}_{\Lambda\Lambda}^{6}\text{He}) = 3.2 \pm 0.1 \text{ MeV},$ at variance with  $\Delta B_{\Lambda\Lambda}^{exp}({}_{\Lambda\Lambda}^{6}\text{He}) = 0.67 \pm 0.17 \text{ MeV}$  [8]. Interpolation within the second set [7] results in a value larger by at least 1 MeV. Since  $V_{\Lambda\Lambda}^{CCQM}$  [4] is softer than the  $V_{\Lambda\Lambda}$  of Ref. [7], which is softer than the  $V_{\Lambda\Lambda}$  of Ref. [6],  $\Delta B_{\Lambda\Lambda}^{CCQM}({}_{\Lambda\Lambda}^{6}He)$  should be even larger. Furthermore, the inclusion of the Pauli-suppressed  $\Lambda\Lambda - \Xi N$  coupling increases  $\Delta B_{\Lambda\Lambda}({}^{6}_{\Lambda\Lambda}$ He) by another 0.2–0.5 MeV [7], and probably by more in the CCQM, owing to its stronger coupling effects. Altogether, I estimate conservatively that  $\Delta B_{\Lambda\Lambda}^{CCQM}({}_{\Lambda\Lambda}^{6}He) > 4.7 \pm 0.5$  MeV, overbinding  ${}_{\Lambda\Lambda}^{6}He$ by more than  $4.0 \pm 0.5$  MeV and thereby destroying the consistency among the bulk of  $\Lambda\Lambda$  hypernuclear data [10].

The CCQM  $\Lambda\Lambda - \Xi N$  coupled-channel interactions used in Ref. [1] are not unambiguously constrained by the scarce, imprecise free-space scattering data [11]. Figure 5 in Ref. [12] shows a variety of S = -2 interactions satisfying such constraints. In particular, there are no  $\Lambda\Lambda$  scattering data to constrain  $a_{\Lambda\Lambda}$ . Recent analysis of the  $\Lambda\Lambda$  invariant mass from the in-medium reaction  $^{12}C(K^-, K^+\Lambda\Lambda X)$  [13] results in  $a_{\Lambda\Lambda} = -1.2 \pm 0.6$  fm [14], consistently with  $a_{\Lambda\Lambda} \sim -0.5$  fm from  ${}_{\Lambda\Lambda}^{6}$  He [6,15], in disagreement with  $a_{\Lambda\Lambda}^{CCQM} = -3.3$  fm. Furthermore, the very strong CCQM  $\Lambda\Lambda - \Xi N$  coupling interaction, which leads to a bound H below the  $\Lambda\Lambda$  threshold [4] and is also responsible for binding  ${}_{\Lambda\Lambda}{}^{3}$ H, is at odds with the latest HAL QCD lattice-simulation analysis which locates the H dibaryon near the  $\Xi N$  threshold [16]. For all these reasons, foremost for heftily overbinding  $^{6}_{\Lambda\Lambda}$  He, the predictive power of the CCQM for S = -2,

TABLE I.  $\Delta B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^{6}\text{He}) = B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^{6}\text{He}) - 2B_{\Lambda}({}_{\Lambda}^{5}\text{He})$  (in MeV) from  $\alpha\Lambda\Lambda$  calculations [6,7] with no  $\Lambda\Lambda - \Xi N$  coupling, and scattering lengths  $a_{\Lambda\Lambda}$  and effective ranges  $r_{\Lambda\Lambda}$  (in fm) of the input  $\Lambda\Lambda$  interaction  $V_{\Lambda\Lambda}$ .  $\Delta B_{\Lambda\Lambda}^{exp}({}_{\Lambda\Lambda}^{6}\text{He}) = 0.67 \pm 0.17$  MeV [8].

	Ref. [6]	Ref. [7]	Ref. [7]				
$-a_{\Lambda\Lambda}$	0.31	0.77	2.81	5.37	10.6	1.90	21.0
$r_{\Lambda\Lambda}$	3.12	2.92	2.95	2.40	2.23	3.33	2.54
$\Delta B_{\Lambda\Lambda}$	0.79	1.51	2.91	3.91	4.51	4.12	8.29

including the prediction of a  ${}^{3}_{\Lambda\Lambda}H$  bound state [1], is questionable.

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