is again a unique solution:

$$(\gamma_{\pi\pi}{}^{\rho})^2 : (\gamma_{KK}{}^{\rho})^2 : (\gamma_{\pi K}{}^{K*})^2 = \tfrac{1}{3} : \tfrac{2}{3} : 1,$$

$$\gamma_{\pi\pi}{}^{\rho}\gamma_{KK}{}^{\rho} < 0;$$

but now $\lambda = -1$ whereas in cases (I) and (II), $\lambda = +1$. It can be shown that a negative λ corresponds to a repulsive force and the equation $\lambda = 1/\alpha(\mu)$ is not satisfied. This means that this case is not consistent with the N/D bootstrap approximation employed here.

It is difficult at this stage to estimate the exact physical significance of the several special cases quoted above. The authors believe that the equations derived are quite general and may well be arrived at by various arguments. It is interesting, however, that they can be derived from N/D bootstrap which is at present the only available tool with a reasonable degree of success in strong interactions. The general group-theoretical properties of these equations are also of interest and are under investigation.

We are grateful to Dr. Tsu Yao for correspondence which led to the correction of an error in the manuscript version of this paper.

Two of the authors (C.H.-M. and P.C.D.) wish to thank Professor Oppenheimer for the support and hospitality shown them during their stay at the Institute

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*A portion of this work was done when the authors were visiting the Florida State University. Work also performed, in part, under the auspices of the U. S. Atomic Energy Commission.

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‖Supported in part by the U. S. Air Force Office of Scientific Research and Development Command.

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The scale on the ordinate for the upper curve of Fig. 2 should be multiplied by four.