
 E R R A T A

RESOLUTION OF THE Σ^- -MASS ANOMALY.

Walter H. Barkas, John N. Dyer, and Harry H. Heckman [Phys. Rev. Letters 11, 26 (1963)].

There is a misprint in line 10, p. 27. For "and we observe 1.5μ ," please read "and we observe 15μ . It is possibly significant that it is the anomalous range that displays the anomalous straggling."

REGGE POLES IN RENORMALIZABLE FIELD THEORIES. Paolo Budini [Phys. Rev. Letters 10, 384 (1963)].

Form (5) should be replaced by

$$\alpha\left(\frac{t}{m_1^2}, \frac{m_1^2}{m_2^2}, e_0^2\right) = \lim_{\lambda \rightarrow 0} \alpha\left[\frac{t}{\lambda}, \frac{m_1^2}{\lambda}, \frac{m_2^2}{\lambda}, e_0^2 d_0(\lambda, m_1, m_2, e_0^2)\right]. \quad (5)$$

As a consequence one can only say that given a solution of (4') in the form $\alpha[y, \mu_1, \mu_2, e_0^2 d_0(1/\mu_1, 1/\mu_2)]$, the physical one is obtained taking the ratio of the variables y, μ_1, μ_2 and putting $d_0 = 1$.

Since (4') asserts that $\alpha(y, \mu, e^2)$ is an invariant of the renormalization group, it obeys the differential equation

$$y \frac{\partial \alpha}{\partial y} + \mu^2 \frac{\partial \alpha}{\partial \mu^2} = e^2 \frac{\partial \alpha}{\partial e^2} \frac{\partial d(x, \mu^2, e^2)}{\partial x} \Big|_{x=1} \quad (6)$$

(we have put for simplicity $\mu_1 = \mu_2 = \mu$). One sees that, in fact, $d \equiv 1$ gives a particular solution, of the form $\alpha = f(y/\mu, e_0^2)$ with f arbitrary, but not the most general; solutions with $d \neq 1$ are possible.

NEW UPPER BOUND FOR THE HIGH-ENERGY SCATTERING AMPLITUDE. T. Kinoshita, J. J. Loeffel, and A. Martin [Phys. Rev. Letters 10, 460 (1963)].

Professor O. W. Greenberg has been kind enough to point out to us a faulty formulation in reference 3 of our paper. The correct version of this footnote is: "The less restrictive assumption that $f(s, \cos\theta)$ be analytic in a Lehman ellipse (semimajor axis $\approx 1 + \alpha/k^4$) gives a weaker bound than (1). See O. W. Greenberg and F. E. Low, Phys. Rev. 124, 2047 (1961)."