Comment on "Casimir Force and In Situ Surface Potential Measurements on Nanomembranes"

In Ref. [1] the frequency shift of an oscillator Δf under the influence of the residual electric force $F_{res}^{el}(z)$ and the Casimir force $F_C(z)$ was measured between Au-coated surfaces of a sphere and a membrane. Using the model for $F_{\rm res}^{\rm el}(z)$ with two fitting parameters, it was claimed that the data for frequency shift are best described by the Drude model approach to the Casimir force ($\chi^2 = 35.3$ and χ^2 probability to exceed 35%) and excludes the plasma model approach ($\chi^2 = 56.1$ and the probability of about 1%). We demonstrate that these results are incorrect, and that the data of Ref. [1] are inconsistent with both theoretical approaches.

First, we note that although a mistake in Eq. (2) in Ref. [1], indicated by us earlier [2], was corrected in an erratum [3], the factor $\eta = \sqrt{1 + A_{\rm rms}^2/z^2}$, where $A_{\rm rms}$ is the rms amplitude of membrane vibrations and z is separation, remained incorrect. In dynamic experiments the corrections due to surface roughness must be included in an expression for the external force, whereas vibrations determine the frequency shift according to corrected Eq. (2). Because of this, the factor η used on p. 4 to correct separations must be replaced with $\eta_{\text{corr}} = \sqrt{1 + 3A_{\text{rms}}^2/(2z^2)}$. Next, we recalculated Δf using the tabulated optical

data of Au [4] extrapolated to zero frequency using either the Drude [5] or the plasma [6] model with parameters $\omega_p = 7.54 \text{ eV}$ and $\gamma = 0.051 \text{ eV}$ which, according to Refs. [1,7], best fit the optical data. In Fig. 1, our computational results for $z\Delta f$ using corrected Eq. (2) and the Drude- and plasma-model extrapolations are shown by the upper and lower solid lines, respectively, for z from 0.118 to 0.230 μ m where, as stated in Ref. [1], the electrostatic force is negligible. The experimental data taken from Fig. 4(c) of Ref. [1] are indicated as crosses whose arms

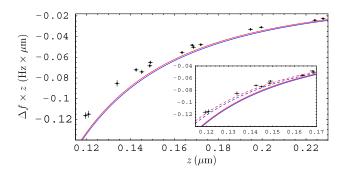


FIG. 1 (color online). The experimental data with respective errors are indicated as crosses. The upper and lower solid lines (dashed lines) show the theoretical results calculated here using the optical data extrapolated by means of the Drude and plasma models, respectively (calculated in Ref. [1] using the Drude and plasma models).

show the experimental errors (the values of z were extracted from the electrostatic force and corrected by a factor η in the Letter [1]).

As is seen in Fig. 1, all 15 data points of the total 32 exclude both the Drude- and plasma-model approaches to the Casimir force. For any fit, the contribution of these data to χ^2 will exceed 300 and 419, respectively. The 17 data points, measured in Ref. [1] at larger separations, where $F_{\rm res}^{\rm el}(z)$ may not be negligible, can only increase these values. As a result, for both approaches the χ^2 probability is much less than 10^{-6} %. Our computational results do not coincide with the theoretical results in Fig. 4(c) of Ref. [1]. The latter are reproduced when one disregards the optical data and uses instead simple Drude and plasma models over the entire frequency region and Eq. (2) with only the first term (see the upper and lower dashed lines in the inset of Fig. 1, respectively). It is known, however, that at $z \leq 0.3 \ \mu m$ core electrons contribute significantly to the dielectric permittivity, so that simple Drude and plasma models cannot be used to calculate the Casimir force [6]. Note that the effect of surface roughness only increases the magnitude of Δf , thus increasing a disagreement between experiment and theory.

To summarize, the calculations presented here demonstrate strong disagreement between the experimental data of Ref. [1] and both theoretical approaches to the Casimir force, thus suggesting the presence of an unaccounted systematic error in the data.

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- M. Bordag,¹ G. L. Klimchitskaya,^{1,2} and
- V. M. Mostepanenko^{1,2} ¹Institute for Theoretical Physics, Leipzig University D-04009 Leipzig, Germany
- ²Central Astronomical Observatory at Pulkovo of the Russian Academy of Sciences, Saint Petersburg, 196140, Russia

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