

Nonperturbative Interband Response of a Bulk InSb Semiconductor Driven Off Resonantly by Terahertz Electromagnetic Few-Cycle Pulses

F. Junginger,¹ B. Mayer,¹ C. Schmidt,¹ O. Schubert,^{1,2} S. Mährlein,¹ A. Leitenstorfer,¹ R. Huber,^{1,2} and A. Pashkin¹

¹*Department of Physics and Center for Applied Photonics, University of Konstanz, Universitätsstr. 10, 78457 Konstanz, Germany*

²*Department of Physics, University of Regensburg, Universitätsstr. 31, 93053 Regensburg, Germany*

(Received 5 April 2012; published 4 October 2012)

Intense multiterahertz pulses are used to study the coherent nonlinear response of bulk InSb by means of field-resolved four-wave mixing spectroscopy. At amplitudes above 5 MV/cm the signals show a clear temporal substructure which is unexpected in perturbative nonlinear optics. Simulations based on a model of a two-level quantum system demonstrate that in spite of the strongly off-resonant character of the excitation the high-field few-cycle pulses drive the interband resonances into a nonperturbative regime of Rabi flopping. The rotating wave approximation breaks down in this case and the system reaches a complete population inversion.

DOI: [10.1103/PhysRevLett.109.147403](https://doi.org/10.1103/PhysRevLett.109.147403)

PACS numbers: 78.47.nj, 42.65.Re, 78.40.Fy, 78.55.Cr

Semiconductors represent a uniquely well-defined laboratory to explore novel limits of nonlinear optics. A key parameter for interaction of a coherent light field with an electronic transition is given by the Rabi frequency $\Omega_R = \mu E/\hbar$. E is the electric field strength and μ the transition dipole moment. When the dephasing rate is negligible compared to Ω_R , coherent Rabi flopping governs the dynamics of electronic systems [1]. If the detuning of the driving electromagnetic field is smaller than Ω_R the response of a system cannot be described by the perturbative approach which usually comes along with off-resonant nonlinear optics. This nonperturbative excitation regime provides access to many fascinating quantum effects in semiconductors such as Rabi splitting, self-induced transparency, and generation of high harmonics [2–8]. In particular, sufficiently intense and ultrashort laser pulses have been exploited to implement ultimate scenarios in which the duration of the light pulse, the Rabi cycle, and the oscillation period of the carrier wave all become comparable in size [9]. Under such conditions, a detailed insight into the nonlinear optical interaction calls for complete phase and amplitude resolution of all interacting light fields. However, in most of the cases, the lack of phase-stable laser pulses and fast detectors does not allow for capturing subcycle polarization dynamics of a system.

The development of ultraintense terahertz laser systems generating phase-stable transients with field amplitudes above 1 MV/cm [10–12] paves the way towards a coherent spectroscopy of extreme nonlinearities in condensed matter systems with absolute sampling of amplitude and phase. Recent experiments performed with high-field multiterahertz pulses have demonstrated a high potential of this approach [13–16]. However, off-resonant excitation of Rabi flopping has remained almost unexplored. This regime bears remarkable effects due to a distinct violation of the rotating wave approximation (RWA) [17]. The main obstacle for experimental investigations of these

phenomena was the lack of sufficiently intense and phase-locked pulses. The latest breakthrough in the generation and field-resolved detection of multiterahertz pulses with peak fields up to 100 MV/cm [18,19] opens up the possibility to explore this highly nonperturbative regime.

In this Letter, we report the nonlinear response of the bulk semiconductor indium antimonide (InSb) excited far below interband resonance using our novel high-field multiterahertz laser source [19]: time-resolved four-wave mixing (FWM) signals are recorded with amplitude and phase at different terahertz peak fields of the excitation pulses. For the highest intensities, the Rabi frequency becomes comparable to the detuning and the interband resonance is driven into a nonperturbative regime of Rabi flopping. Our simulations of the FWM response based on the Maxwell-Bloch equations provide a qualitative understanding of the phenomena observed experimentally.

The sample under study is a (100)-oriented undoped single crystal of InSb [Fig. 1(a)] mechanically polished to a thickness of 30 μm and kept at room temperature. Terahertz transients with a center frequency of $f_0 = 23$ THz, a bandwidth of 8 THz (full width at half maximum, FWHM) and variable peak fields between 2 and 5.3 MV/cm are generated by difference frequency mixing of near-infrared pulse trains with a repetition rate of 1 kHz in gallium selenide (GaSe) emitters [Figs. 1(b) and 1(c)] [19]. A specific feature of our experiment is a large detuning of 18 THz between the terahertz transients and the nearest interband resonance of InSb ($E_g/\hbar = 41.1$ THz). Thus, the entire terahertz spectrum is located well below the band gap excluding the possibility of direct linear excitation of electron-hole pairs.

We use a two-dimensional scheme of nonlinear spectroscopy [Fig. 1(d)]. In order to perform a terahertz multiwave mixing experiment, a pair of mutually synchronized pulses is obtained by splitting the terahertz beam after the

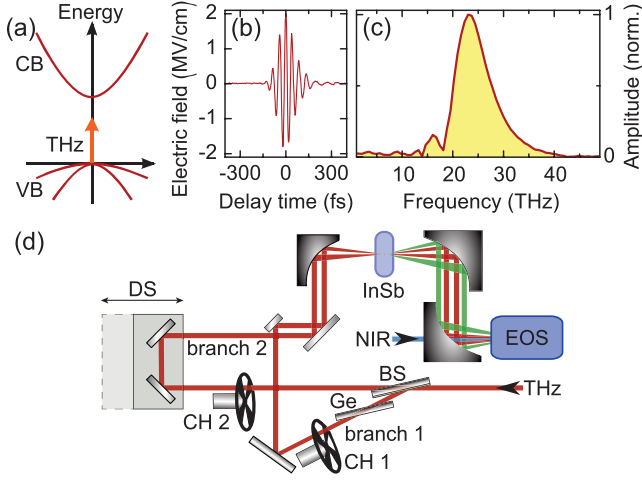


FIG. 1 (color). (a) Schematic band structure of InSb with conduction (CB) and valence bands (VB). The vertical arrow illustrates a nonresonant excitation by a terahertz pulse. (b) Terahertz transient generated by a $370\text{-}\mu\text{m}$ -thick GaSe crystal. (c) Amplitude spectrum of the terahertz transient in panel (b). (d) Scheme of the FWM setup: two phase-locked terahertz pulse trains are generated by a Ge beam splitter (BS). Both branches are modulated by mechanical choppers (CH) and delayed with respect to each other via a translation stage (DS). The terahertz transients are noncollinearly focused on the sample and the transmitted signals are detected by electro-optic sampling (EOS).

GaSe emitter crystal. In this way nonlinear mixing effects of two pulses within the emitter itself are precluded. A germanium (Ge) beam splitter set at Brewster's angle and coated with a 6-nm-thick gold layer provides a splitting ratio of 1:1. The transmitted electric field in branch 2 (E_2) is delayed by a retroreflector stage. The transient in branch 1 (E_1) propagates through a second Ge wafer to match the dispersion of E_2 . The terahertz beams in both branches are individually chopped with frequencies of 500 and 250 Hz, respectively, and the chopper phases are locked to the 1 kHz pulse train. This configuration enables a fast and efficient acquisition of the transmitted signals from only branch 1, only branch 2, or both branches at the same time (E_{12}). The terahertz beams are tightly focused onto the same spot on the sample (FWHM: $85\ \mu\text{m}$). The emerging linear and nonlinear fields are collected with large numerical aperture and directed onto a $140\text{-}\mu\text{m}$ -thick GaSe electro-optic sensor gated by near-infrared pulses with a duration of 8 fs.

Figure 2(a) shows the total transmitted field E_{12} as a function of the electro-optic sampling delay time t and the relative temporal offset τ between the terahertz pulses. The signal from branch 1 with a fixed temporal position appears as a set of vertical lines of constant phase centered around $t = 0$ ps, whereas the diagonal lines correspond to the delayed signal from branch 2. The external peak fields are $2\ \text{MV/cm}$ per pulse. The field strengths within InSb

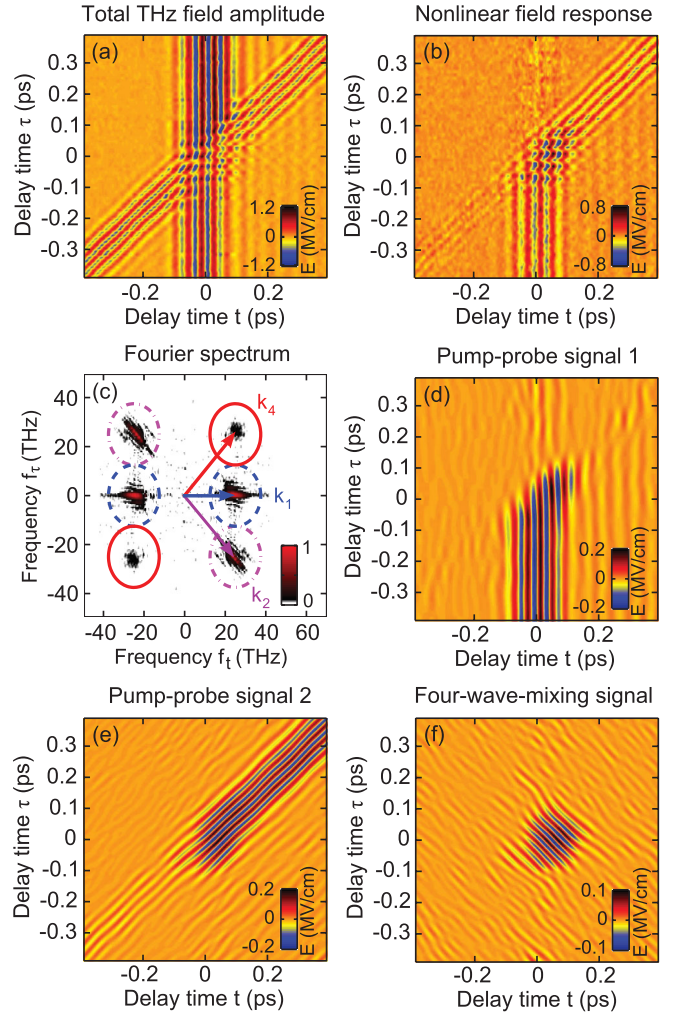


FIG. 2 (color). (a) Electric field of two terahertz transients traversing InSb plotted as a function of delay time τ and sampling time t . (b) Nonlinear signal E_{NL} . (c) Two-dimensional FT of panel (b) revealing pump-probe signals at wave vector positions \mathbf{k}_1 and \mathbf{k}_2 and the FWM signature at \mathbf{k}_4 . (d) Selective inverse FT of the pump-probe signal emerging in the direction of \mathbf{k}_1 . (e) Inverse FT of the pump-probe signature at \mathbf{k}_2 only. (f) Inverse FT of the FWM signal in the direction of \mathbf{k}_4 only.

are attenuated by a Fresnel factor $\tilde{t} = 2/(n+1) = 0.4$ defined by a refractive index of $n = 4$ in the frequency range of interest [20]. The nonlinear signal E_{NL} shown in Fig. 2(b) is retrieved by subtracting the contributions of individual transients E_1 and E_2 from the total response E_{12} : $E_{NL} = E_{12} - E_1 - E_2$. The Fourier transform (FT) of E_{NL} [Fig. 2(c)] has a direct correspondence with the wave vector space and, thus, allows us to disentangle different contributions to the total nonlinear field response [14]. The inverse Fourier transform of selected regions in frequency space depicted in Fig. 2(c) provides the temporal fingerprints of multiwave mixing signals of various orders. The pump-probe signals for each terahertz pulse, located around wave vectors $\pm\mathbf{k}_1$ ($f_t = \pm f_0$, $f_\tau = 0$) and $\pm\mathbf{k}_2$

($f_t = \pm f_0, f_\tau = \mp f_0$) [Fig. 2(c)], are shown in Figs. 2(d) and 2(e), respectively. These signals correspond to the transmission change of the sample excited by the first and probed by the second terahertz transient. They do not depend on the relative phase.

Most remarkably, a FWM signal at the wave vector $\pm \mathbf{k}_4 = \pm(2\mathbf{k}_1 - \mathbf{k}_2)$ ($f_t = \pm f_0, f_\tau = \pm f_0$) becomes clearly discernible [Fig. 2(c)]. This signal features lines of constant phase that point along a downward diagonal [Fig. 2(f)] indicating its dependence on the respective phases of both terahertz transients. Thus, it bears information about coherence induced in the sample. Therefore, we will concentrate in the following on the FWM signal which is free from the influence of incoherent excitation processes dominating the pump-probe response such as two-photon absorption [21] or impact ionization [22].

To study the field dependence of the FWM signal, external peak fields of 2, 3.5, and 5.3 MV/cm per pulse are selected. The corresponding pulse intensities are I_0 , $3I_0$, and $7I_0$ where $I_0 = 10.6 \text{ GW/cm}^2$. For the lowest peak electric field we observe an oval-shaped envelope of the FWM signal, identical to the cross-correlation function of both pulses [Fig. 3(a)]. This result is as expected in the limit of perturbative nonlinear optics far from resonance. In contrast, increasing the field strength up to 3.5 MV/cm leads to a deviation from the symmetric profile resulting in an *S*-shaped signal [Fig. 3(b)]. Surprisingly, the maximum field of 5.3 MV/cm leads to a splitting of the FWM signal [Fig. 3(c)]. A minimum appears in the temporal region where the strongest total excitation field is present. This signature is an unequivocal indication of an extremely nonlinear interaction.

Our experimental results may be understood by using a simplified model of a two-level system representing the interband resonance in InSb. The simulation is based on the Maxwell-Bloch equations [23] which are solved numerically without applying the RWA and the slowly varying envelope approximation (SVEA). The solution is obtained by an iterative predictor-corrector finite element method. For the simulations we assume a dephasing time of $T_2 = 1 \text{ ps}$ [24] and a depopulation time of $T_1 = 10 \text{ ps}$. Owing to the extremely short terahertz transients the choice of the relaxation time barely affects the results of the simulations as long as T_1 and T_2 are longer than the duration of the terahertz pulse. The transition dipole moment $\mu_{12} = 2.4 e\text{\AA}$ and the density of the two-level systems $N = 2.9 \times 10^{20} \text{ cm}^{-3}$ were adjusted in order to provide the best agreement with the shapes and intensities of the FWM signals measured experimentally.

Figures 3(d)–3(f) show the time domain FWM signals simulated for the same peak fields as those used for the experimental results depicted in the panels on the left. For the simulation we assume Fourier-limited Gaussian terahertz transients with a FWHM duration set to the pulse width measured experimentally. A perturbative response

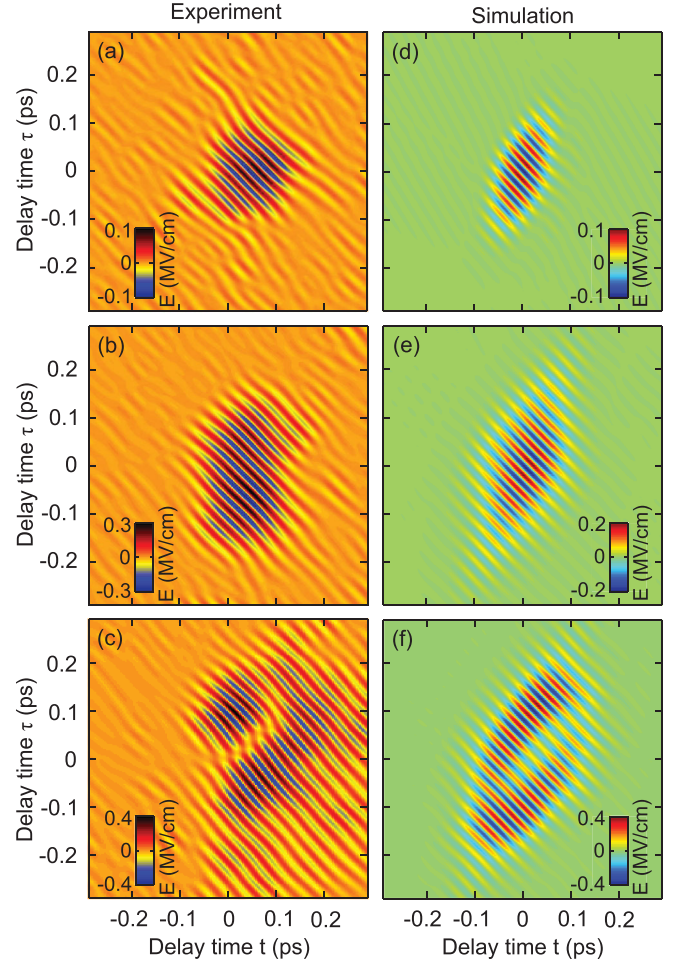


FIG. 3 (color). (a) Oval FWM signal by applying an external field of 2 MV/cm per pulse. (b) *S*-shaped FWM signature driven by an external field of 3.5 MV/cm. (c) Split FWM signal at an external exciting field strength of 5.3 MV/cm. (d)–(f) Calculated FWM signatures reproducing the main features of the corresponding measured signals on the left side of each image.

at moderate field strengths leads to an oval envelope of the FWM signal [Fig. 3(d)]. At the intermediate terahertz intensity this envelope evolves into an *S* shape [Fig. 3(e)]. Finally, external peak fields above 5 MV/cm induce a minimum in the center of the signal, where the total field is strongest. This results in two sidelobes similar to those observed in the experiment [see Figs. 3(c) and 3(f)]. As one can clearly see, the simulations with a single energy electronic resonance allow us to reproduce the essential features observed in the experiment. This result is highly surprising since it is well known that interband excitations in bulk semiconductors like InSb include a continuum of electronic resonances with a broad distribution of frequencies above the absorption edge. Nevertheless, simulations using a set of two-level systems with a density distribution fitting the joint density of states in InSb essentially lead to the same results as those shown in Figs. 3(c) and 3(f). The reason for that is

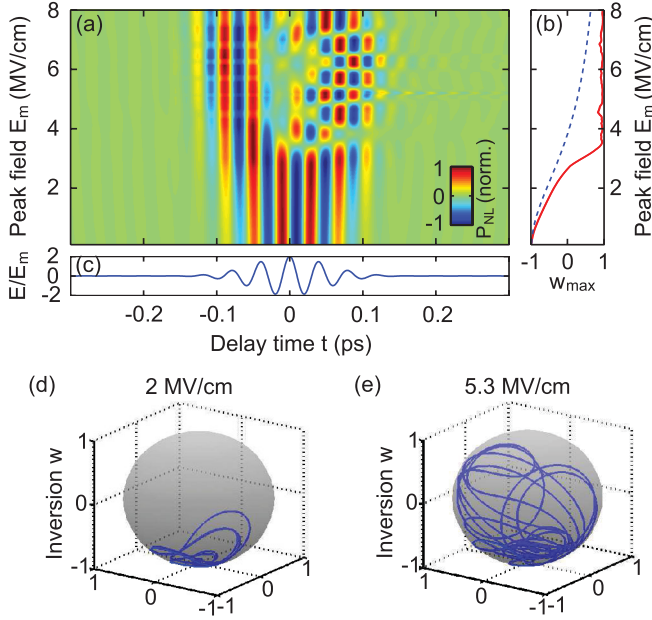


FIG. 4 (color). Simulated response of the two-level system for $\tau = 0$ ps: (a) Normalized polarization of the two-level system at the fundamental frequency of the driving field as a function of the delay time t and the peak electric field E_m . (b) Maximum inversion w_{\max} as a function of E_m (red solid line) and of the amplitude of a continuous wave excitation (blue dashed line) (c) Driving field of the terahertz transients. (d)–(e) Corresponding pathways of the Bloch vector for a moderate ($E_m = 2$ MV/cm) and a high ($E_m = 5.3$ MV/cm) peak electric field.

the steep dependence of the nonperturbative response on the detuning frequency. This finding is supported by the density of the two-level systems N estimated from our simulations which constitutes only about 1% of the total number of available states in the conduction band of InSb. These are the states near the band edges which provide the interband transitions with the smallest detuning.

A qualitative physical understanding of the splitting observed in the FWM signal can be obtained by considering the polarization response of a two-level system driven by two terahertz transients with zero delay time ($\tau = 0$ ps). Figure 4(a) depicts the simulated time-resolved polarization at the frequency of the driving pulses for different peak electric fields E_m . In the case of moderate fields the maximum population inversion w_{\max} shown in Fig. 4(b) remains negative and the deflection of the Bloch vector from the ground state is minor as illustrated in Fig. 4(d). Therefore, the response of the two-level system is perturbative and the shape of the polarization signal [Fig. 4(a)] follows the profile of the driving fields shown in Fig. 4(c). The FWM signal, thus, can be described in terms of an instantaneous third-order nonlinearity. However, this picture breaks down as soon as field amplitudes exceed 3 MV/cm and a clear splitting of the temporal signature of the FWM signal starts to develop. In this

case the maximum population inversion shown by the solid red line in Fig. 4(b) becomes positive ($w_{\max} > 0$), indicating the onset of a strongly nonperturbative regime. As illustrated in Fig. 4(e), a terahertz peak field of 5.3 MV/cm promotes the system almost to the limit of complete population inversion. This surprising result radically differs from the case of an off-resonant excitation by a continuous wave where, assuming the RWA, the complete population inversion can be achieved only in the limit of infinitely high electric fields [see the dashed blue line in Fig. 4(b)]. This fact indicates a clear violation of the RWA and SVEA for our conditions of the off-resonant excitation where the generalized Rabi frequency is comparable to the central frequency and the bandwidth of the excitation pulses. Finally, in the center of the FWM signal, where the driving electric field reaches its maximum, the polarization oscillates mainly at high harmonic frequencies [25]. The response at the fundamental frequency becomes weak leading to the observed minimum [Fig. 4(a)]. This regime of a nonperturbative excitation sets in when the maximum Rabi frequency $\Omega_R/2\pi = 2\tilde{\mu}_{12}E_m/h$ becomes comparable to the large detuning of 18 THz at $E_m = 3$ MV/cm.

In conclusion, we have studied the off-resonant FWM response of bulk InSb which provides a direct access to the coherent dynamics of the interband polarization response at terahertz frequencies. The observed splitting of the FWM signals for electric fields above 5 MV/cm manifest the onset of a nonperturbative response of Rabi flopping. This extremely nonlinear behavior underpins the high potential of high-field multiterahertz technology for high harmonics generation [26] and coherent control of quantum states in semiconductors [15,27]. The high fields and excellent signal-to-noise performance of terahertz multiwave mixing open a way for investigations of the coherent response in a large variety of important resonances in complex systems such as intermolecular librations in hydrogen-bonded liquids or energy gaps in superconducting condensates.

-
- [1] C. Cohen-Tannoudji *et al.*, *Quantum Mechanics* (Wiley, New York, 1977).
 - [2] M. Wegener, D.S. Chemla, S. Schmitt-Rink, and W. Schäfer, *Phys. Rev. A* **42**, 5675 (1990).
 - [3] S.T. Cundiff, A. Knorr, J. Feldmann, S.W. Koch, E.O. Göbel, and H. Nickel, *Phys. Rev. Lett.* **73**, 1178 (1994).
 - [4] H. Giessen, A. Knorr, S. Haas, S.W. Koch, S. Linden, J. Kuhl, M. Hetterich, M. Grün, and C. Klingshirn, *Phys. Rev. Lett.* **81**, 4260 (1998).
 - [5] A. Schülzgen, R. Binder, M.E. Donovan, M. Lindberg, K. Wundke, H.M. Gibbs, G. Khitrova, and N. Peyghambarian, *Phys. Rev. Lett.* **82**, 2346 (1999).
 - [6] O.D. Mücke, T. Tritschler, M. Wegener, U. Morgner, and F.X. Kärtner, *Phys. Rev. Lett.* **87**, 057401 (2001).

- [7] S. G. Carter, V. Birkedal, C. S. Wang, L. A. Coldren, A. V. Maslov, D. S. Citrin, and M. S. Sherwin, *Science* **310**, 651 (2005).
- [8] H. Choi, V.-M. Gkortsas, L. Diehl, D. Bour, S. Corzine, J. Zhu, G. Höfler, F. Capasso, F. X. Kärtner, and T. B. Norris, *Nature Photon.* **4**, 706 (2010).
- [9] T. Tritschler, O. D. Mücke, M. Wegener, U. Morgner, and F. X. Kärtner, *Phys. Rev. Lett.* **90**, 217404 (2003).
- [10] K. Reimann, R. P. Smith, A. M. Weiner, T. Elsaesser, and M. Woerner, *Opt. Lett.* **28**, 471 (2003).
- [11] J. Hebling, K.-L. Yeh, M. C. Hoffmann, B. Bartal, and K. A. Nelson, *J. Opt. Soc. Am. B* **25**, B6 (2008).
- [12] H. Hirori, A. Doi, F. Blanchard, and K. Tanaka, *Appl. Phys. Lett.* **98**, 091106 (2011).
- [13] C. W. Luo, K. Reimann, M. Woerner, T. Elsaesser, R. Hey, and K. H. Ploog, *Phys. Rev. Lett.* **92**, 047402 (2004).
- [14] W. Kuehn, K. Reimann, M. Woerner, and T. Elsaesser, *J. Chem. Phys.* **130**, 164503 (2009).
- [15] S. Leinß, T. Kampfrath, K. v. Volkman, M. Wolf, J. T. Steiner, M. Kira, S. W. Koch, A. Leitenstorfer, and R. Huber, *Phys. Rev. Lett.* **101**, 246401 (2008).
- [16] W. Kuehn, K. Reimann, M. Woerner, T. Elsaesser, R. Hey, and U. Schade, *Phys. Rev. Lett.* **107**, 067401 (2011).
- [17] J. Larson, *Phys. Scr.* **76**, 146 (2007).
- [18] A. Sell, A. Leitenstorfer, and R. Huber, *Opt. Lett.* **33**, 2767 (2008).
- [19] F. Junginger, A. Sell, O. Schubert, B. Mayer, D. Brida, M. Marangoni, G. Cerullo, A. Leitenstorfer, and R. Huber, *Opt. Lett.* **35**, 2645 (2010).
- [20] S. Adachi, *Phys. Rev. B* **35**, 7454 (1987).
- [21] P. D. Olszak, C. M. Cirloganu, S. Webster, L. A. Padilha, S. Guha, L. P. Gonzalez, S. Krishnamurthy, D. J. Hagan, and E. W. Van Stryland, *Phys. Rev. B* **82**, 235207 (2010).
- [22] M. C. Hoffmann, J. Hebling, H. Y. Hwang, K.-L. Yeh, and K. A. Nelson, *Phys. Rev. B* **79**, 161201 (2009).
- [23] R. W. Ziolkowski, J. M. Arnold, and D. M. Gogny, *Phys. Rev. A* **52**, 3082 (1995).
- [24] J. T. Andrews and P. Sen, *Quantum Semiclass. Opt.* **10**, 663 (1998).
- [25] T. Tritschler, O. D. Mücke, and M. Wegener, *Phys. Rev. A* **68**, 033404 (2003).
- [26] D. Golde, T. Meier, and S. W. Koch, *Phys. Rev. B* **77**, 075330 (2008).
- [27] B. E. Cole, J. B. Williams, B. T. King, M. S. Sherwin, and C. R. Stanley, *Nature (London)* **410**, 60 (2001).