



## Optical Spectroscopy of Molecular Positronium

D. B. Cassidy, T. H. Hisakado, H. W. K. Tom, and A. P. Mills, Jr.

*Department of Physics and Astronomy, University of California, Riverside, California 92521-0413, USA*

(Received 11 January 2012; published 30 March 2012)

We report optical spectroscopic measurements of molecular positronium ( $\text{Ps}_2$ ), performed via a previously unobserved  $L = 1$  excited state.  $\text{Ps}_2$  molecules created in a porous silica film, and also in vacuum from an Al(111) crystal, were resonantly excited and then photoionized by pulsed lasers, providing conclusive evidence for the production of this molecular matter-antimatter system and its excited state. Future experiments making use of the photoionized vacuum  $L = 1$   $\text{Ps}_2$  could provide a source of  $\text{Ps}^+$  ions, as well as other multipositronic systems, such as  $\text{Ps}_2\text{H}^-$  or  $\text{Ps}_2\text{O}$ .

DOI: 10.1103/PhysRevLett.108.133402

PACS numbers: 36.10.Dr, 34.80.Lx, 42.62.Fi

The realization that the negative energy electron states of Dirac's relativistic quantum theory [1] actually exist [2] constitutes one of the most remarkable scientific discoveries of the last century. We now know that all particles have a corresponding negative energy state, or antiparticle, and according to the *CPT* theorem antiparticles and particles have opposite additive quantum numbers (such as electric charge or lepton number) but are otherwise identical.

Combinations of particle-antiparticle pairs may form exotic atomic states, the most well-known example being positronium ( $\text{Ps}$ ), the bound state between an electron and a positron [3]. Although  $\text{Ps}$  will decay via self-annihilation, its lifetime is long enough that it has a well defined atomic structure, and it is even possible for two  $\text{Ps}$  atoms to form a molecular state known as  $\text{Ps}_2$  [4]. This molecule is stable [5] but has a shorter lifetime against annihilation than  $\text{Ps}$ , so its existence may be inferred from changes in the  $\text{Ps}$  decay rate that depend on the  $\text{Ps}$  density. This was done in previous experiments [6,7], but such observations are necessarily indirect since other processes can also affect the  $\text{Ps}$  decay rate in a similar way.

In this Letter we report observations of optically induced changes in the annihilation yield of dense  $\text{Ps}$  over a narrow range of wavelengths near 251 nm consistent with what would be expected for electric dipole transitions between the ground state of  $\text{Ps}_2$  and a predicted but previously unobserved  $L = 1$  excited state [8]. These observations provide conclusive evidence for the production of the di-positronium molecule under conditions allowing  $\text{Ps}$ - $\text{Ps}$  interactions and verify the existence of one of its excited states. Using the methods we describe here it should be possible to measure the transition frequency accurately enough to test relativistic energy corrections [9], and photoionizing excited state  $\text{Ps}_2$  provides a way to produce  $\text{Ps}^+$  ions that could be used in the formation of exotic two-positron systems such as positronic water,  $\text{Ps}_2\text{O}$  [10].

The experiments were performed by implanting intense pulses of 1–4 keV positrons [11] from an accumulator [12] into either a porous silica film [6,13] or an Al(111) crystal

[7], resulting in the creation of  $\sim 10 \times 10^6$   $\text{Ps}$  atoms around 1 ns. A 2.3 T magnetic field at the target compresses the positron pulse to an areal density of up to  $10^{11} \text{ cm}^{-2}$ , which produces interacting  $\text{Ps}$  atoms [14]. We note that at this magnetic field,  $m = 0$  triplet  $\text{Ps}$  atoms have a lifetime  $< 2$  ns, and so in the following when we refer to long-lived triplet states we mean only  $|m| = 1$  atoms, whose vacuum lifetime is 142 ns, regardless of the magnetic field [3].

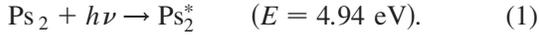
We measure the amount of  $\text{Ps}$  produced using single shot positron annihilation lifetime spectroscopy [15]. The fraction of our measured lifetime spectra in the interval 50–300 ns after the positron pulse yields the parameter  $f_d$ , a measure of the amount of long-lived triplet  $\text{Ps}$  present [6]. When a high density positron beam is implanted into certain targets, we observe density dependent changes in the  $\text{Ps}$  decay rate that we ascribe either to spin exchanging scattering events that change triplet  $\text{Ps}$  atoms into short lived ( $m = 0$ ) states [14], or to the formation of  $\text{Ps}_2$  (whose lifetime in the ground state is  $\sim 0.22$  ns). Distinguishing between different quenching signals is not possible using only lifetime spectroscopy with our present experimental arrangement.

Fortunately, there is thought to be at least one excited state of the  $\text{Ps}_2$  molecule [8], which means that this system may be studied via optical spectroscopy. The  $L = 1$  excited molecular state can, in simplistic terms, be thought of as a  $2P$  atom bound to a  $1S$  atom, with the latter acting as though it were a dielectric medium [9], changing the wavelength of the  $1S$ - $2P$  transition from  $\sim 243$  nm for free  $\text{Ps}$  to  $\sim 251$  nm for  $\text{Ps}_2$ . This difference is large enough that there will be no confusion between atomic and molecular  $\text{Ps}$  excitation, which is important given that any experimental conditions that result in the production of  $\text{Ps}_2$  will invariably also produce atomic  $\text{Ps}$  in much larger quantities.

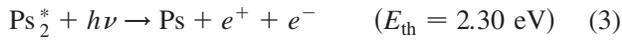
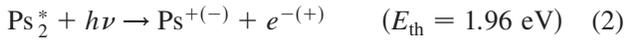
Spectroscopy of  $\text{Ps}_2$  molecules was performed with essentially the same methods as our experiments with  $\text{Ps}$ , as described elsewhere [13]; we use a Nd-YAG-pumped dye laser to generate tunable ultraviolet (uv) light pulses around 251 nm, with a bandwidth of  $\sim 100$  GHz. Residual green light from the pump laser is used to ionize the excited

state molecules, with pulse energies of  $\sim 1$  mJ and 30 mJ for the uv and green lasers, respectively. The  $\sim 2$  mm FWHM uv-laser beam and the  $\sim 0.15$  mm FWHM positron beams were aligned using a phosphor screen and CCD camera, and the 6 mm FWHM green laser was adjusted to overlap the uv light.

We excite  $\text{Ps}_2$  molecules to the first excited state [8] using  $\sim 251$  nm light:



Excited state molecules created by the uv light were subsequently ionized with green light (532 nm; 2.33 eV) obtained from the (doubled) output of the pump laser via the following mechanisms:



These processes have different energy thresholds,  $E_{\text{th}}$ , so that in principle one could select process 2 or 2 & 3 using a tunable ionization laser. Process 3 may be thought of as the photoionization of a  $2P$  Ps atom in the presence of (i.e., weakly bound to) a spectator ground state Ps atom. A third possibility, photodissociation into two Ps atoms:  $\text{Ps}_2^* + h\nu \rightarrow \text{Ps} + \text{Ps}^*$ , will likely have a small probability for occurring but could include a rich series of resonances [16] corresponding to Rydberg Ps states that would be interesting to examine experimentally.

The  $\text{Ps}_2$  detection mechanism relies on observing changes in the Ps decay rate that are correlated with the laser wavelength. That is, when long-lived atoms combine at high density to form  $\text{Ps}_2$  the decay rate increases, and  $f_d$  will decrease. If these molecules are resonantly excited and photoionized, as in Eq. (3), this effect will be counteracted due to the production of some long-lived triplet Ps, but there will be no observable effect if Ps ions [17] are produced as shown in Eq. (2) as these will also annihilate rapidly [18]. However, a green light pulse sufficient to ionize  $\text{Ps}_2^*$  will also ionize  $\text{Ps}^\pm$  ions, so regardless of whether process (2) or (3) occurs, we shall end up with a free positron and electron and a Ps atom, the last being in a long-lived ( $|m| = 1$ ) triplet state 50% of the time [19]. Thus, of all the positronium molecules that are excited to the  $L = 1$  state, we should expect to observe an increase in  $f_d$  due to, at most, half of them (assuming the liberated positrons annihilate quickly with bound electrons in the target, which is likely since they will not have sufficient energy to make further Ps atoms).

$\text{Ps}_2$  molecules produced inside an open porous silica matrix will not, for the most part, be able to escape into vacuum owing to their short lifetime. Thus, laser spectroscopy of such molecules has to occur inside the pores. Experiments with confined Ps atoms [20] have demonstrated that it is feasible to make such measurements. Figure 1(a) shows the laser induced change in  $f_d$  as a

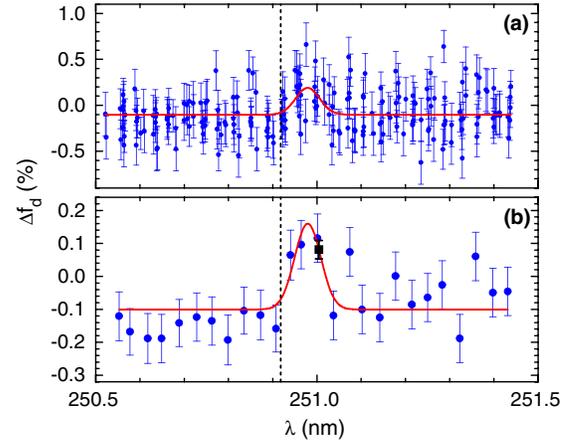


FIG. 1 (color online).  $\text{Ps}_2$  resonance measured inside a porous silica film. All data recorded in four separate runs as a function of the laser (vacuum) wavelength (a) and rebinned data (b). The vertical scale is the change in the delayed fraction in percent. The measurement for the single square point in (b) was obtained as described in the text. The solid lines are simple Gaussian fits to the data from which we determine a line center of  $\lambda_0 = (250.979 \pm 0.006)$  nm, amplitude  $(0.027 \pm 0.008)\%$ , and a full width at half maximum (FWHM) of  $(0.069 \pm 0.026)$  nm. The dashed vertical line indicates the theoretical resonant wavelength for the  $\text{Ps}_2$  excitation in vacuum.

function of the wavelength, measured with the high density positron beam implanted into a porous silica target. The measured signal is  $\Delta f_d = f_d(\text{on}) - f_d(\text{off})$  (on and off refer to the lasers) for four runs, each made up of 52 pairs of single shot measurements of  $f_d$  with laser-on and laser-off at various wavelengths. For each run the subtraction of the laser-off measurements was done using a straight line fit to the laser-off measurements versus measurement number. The errors were assigned so as to make the straight line fits have unity reduced chi square. Figure 1(b) shows the same data as 1(a), ordered by uv wavelength and averaged over independent sets of 8 neighboring measurements with the wavelength and signal both weighted by the inverse square of the individual errors. The square point in Fig. 1(b) is from a separate run made up of 36 pairs of measurements of  $\Delta f_d = f_d(\text{on}) - f_d(\text{off})$ , where on and off refer to uv light on and off resonance, at 251.0 and 250.55 nm. The  $\Delta f_d$  value for this run,  $\Delta f_d = (0.18 \pm 0.02)\%$ , was shifted to include the  $-(0.10 \pm 0.02)\%$  laser-off background and its error determined from the Gaussian fit to the measurements in Fig. 1(a). A peak is observed at a wavelength close to the calculated value [8,9] of  $\lambda_0 = (250.9179 \pm 0.0011)$  nm. The data of Fig. 1 were obtained using an arrangement in which the wavelength was automatically tuned, and was measured using a wave meter, giving an accuracy of  $\pm 0.005$  nm.

The observed line center is  $(250.979 \pm 0.006)$  nm, which is  $(0.061 \pm 0.006)$  nm longer than expected for  $\text{Ps}_2$  in vacuum. In fact, because the  $\text{Ps}_2$  was observed inside a confining pore we had thought that the resonance

frequency might be shifted from the vacuum level, but to a *shorter* wavelength, as was observed for atomic Ps in a similar target material, where a shift of  $-0.05$  nm was measured [20]. This was attributed to the effect of the cavity on the larger  $2P$  state of the atom; the opposite effect observed for  $\text{Ps}_2$  is not presently understood theoretically. It can, however, be explained qualitatively if the Ps and  $\text{Ps}_2$  internal energies are reduced by similar amounts due to van der Waals interactions with the cavity walls, whereas, due to its higher mass, the positive confinement energy shift for  $\text{Ps}_2$  is only half that of Ps.

The measured linewidth is  $(0.069 \pm 0.026)$  nm; that this is the same as the  $0.070$  nm Doppler width expected for free  $\text{Ps}_2$  at a temperature of  $300$  K is undoubtedly fortuitous since the thermalization time for Ps in porous silica is of the order of  $10$  ns, far longer than the  $0.22$  ns  $\text{Ps}_2$  lifetime. Thus, after the  $\text{Ps}_2$  is formed, with release of its  $\sim 0.4$  eV binding energy [8] to its center of mass, it will not have time to fully thermalize. However, we expect that  $\text{Ps}_2$  optical transitions in the porous material will be line narrowed [21], and that the observed linewidth actually has nothing to do with the usual Doppler effect, and is more likely to be related to some property of the sample, such as the distribution of cavity sizes [20].

Although porous silica films are very useful for confining Ps atoms, and thus also for observing Ps-Ps interactions [14], for spectroscopic investigations it is preferable to produce  $\text{Ps}_2$  in vacuum. Moreover, using silica films has the disadvantage that intense laser light can cause significant and permanent damage to the target material, reducing the amount of Ps (and  $\text{Ps}_2$ ) formation. The accumulation of laser damage means that one cannot collect data indefinitely, limiting the available statistics. In addition to this, temporary damage sites produced during the laser pulse [20] further reduce the amount of Ps formation, which is why the off resonance background of Fig. 1(b) is negative.

We also conducted experiments using an Al(111) target, from which it is possible to produce molecular positronium in vacuum [7]. The Al(111) surface was prepared by  $\text{Ar}^+$  ion bombardment for  $1$  hr at  $549^\circ\text{C}$ , but we found that even after cleaning these surfaces were rather unstable, possibly due in part to light induced chemical effects when there is some small amount of oxide present on the surface [22]. The sample instability is evident in the data of Fig. 2(a) which shows the time sequence of  $f_d$  measurements with and without the uv laser. The effect of the increasing amount of Ps formation with subsequent pairs of shots is removed from the data by subtracting the laser-off data, represented by its best fitting quadratic curve, from the uv-laser-on data. The errors are taken to be those that give unity reduced chi square for the quadratic fit to the laser-off data. The  $\text{Ps}_2$  resonance is clearly observed in Fig. 2(b), which shows the measurements with the  $532$  nm (green)  $\text{Ps}_2^*$  photoionizing laser on. The measurements with the green laser off in Fig. 2(c) are consistent with no effect,

as expected, since the unionized  $\text{Ps}_2^*$  molecules quickly annihilate in  $\sim 0.44$  ns or decay back to the ground state [8] (with a 19% branching ratio [9]), resulting in no signal. These data were taken using our initial laser setup that required hand tuning of the wavelength, which was measured to an accuracy of  $\pm 0.01$  nm using a grating spectrometer and a helium  $3^1P-2^1S$  reference line.

A single Gaussian fit to the data of Fig. 2(b) [dashed (red) line] gives a line center  $\lambda_0 = (250.97 \pm 0.02)$  nm, significantly longer than the calculated vacuum wavelength. However,  $\text{Ps}_2$  molecules produced at the Al(111) surface will only travel  $\sim 0.1$  mm before annihilation. To obtain maximal overlap with the  $\text{Ps}_2$  cloud in vacuum the uv laser was directed into the target at an angle of  $\sim 25^\circ$ . As a result there will be a Doppler shift of the transition to a longer wavelength for incident light, and to a shorter wavelength for reflected light, as has been observed using Ps formed on Si surfaces [23]. As we expect there to be equal and opposite Doppler shifts associated with the incident and reflected light, we fitted a single-width double

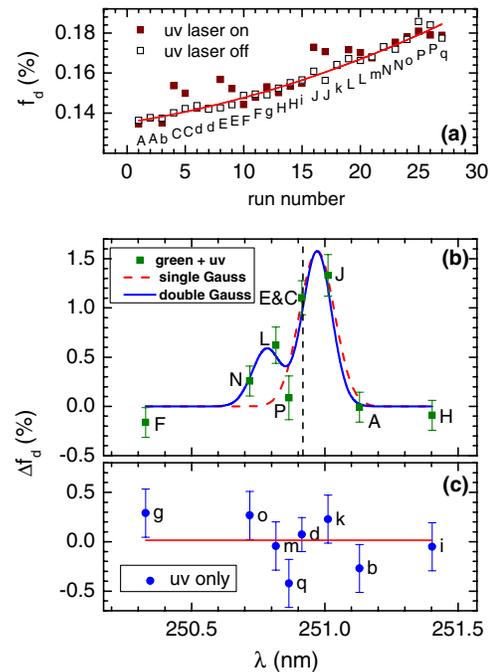


FIG. 2 (color online).  $\text{Ps}_2$  resonance measured in vacuum with an Al(111) crystal at room temperature: (a) data with uv laser-on and off showing variations in  $f_d$  during the measurement and the quadratic fit to the laser-off data. Measurements are labeled sequentially using upper (lower) case letters for green photoionizing laser on (off). (b) lineshape measured with the uv and green lasers on. The dashed line in (b) is a single Gaussian fit to the data and the solid line is a Double Gaussian fit (with a single width), as explained in the text. The dashed vertical line indicates the theoretical resonant wavelength for the  $\text{Ps}_2$  excitation in vacuum. (c) Measurement with uv laser on and green laser off, showing the best fitting horizontal line,  $\Delta f_d = (0.016 \pm 0.087)\%$ , consistent with no effect.

Gaussian function to the data in Fig. 2(b) [solid (blue) line]. To obtain convergence, we constrained the amplitude of the main peak to be the same as that of the single Gaussian fit. The amplitude of the peak corresponding to the reflected light depends on the reflection coefficient, which is unknown, but will be  $<1$ . The center of the two peak positions is  $(250.88 \pm 0.02)$  nm, which is  $\sim 2\sigma$  shorter than the calculated vacuum wavelength. However, even though the improved reduced chi square with respect to the single Gaussian fit ( $\chi^2/\nu = 5.46/4$  instead of  $15.10/5$ ), supports the conclusion that light reflection should be taken into account, the shape of the assumed double Gaussian line shape cannot be correct; the exact value of the transition wavelength for these data is therefore not known as accurately as the double Gaussian fit suggests, and the measurement is broadly consistent with the theoretical value.

The full width at half maximum (FWHM) of the double and single Gauss fits is  $\sim 0.15$  nm, indicating  $\text{Ps}_2$  emission energies of  $\sim 0.2$  eV, consistent with the 0.435 eV binding energy and the energetics of  $\text{Ps}_2$  formation at an Al(111) surface [24]. The  $(1.5 \pm 0.5)\%$  amplitude implies that more than 5% of the  $\sim 2 \times 10^7$  slow positrons implanted into the Al target end up forming at least half a million  $\text{Ps}_2$  molecules (where we assume that the  $\sim 100$  GHz bandwidth of the laser is power broadened so as to cover  $\sim 30\%$  of the  $\text{Ps}_2 \rightarrow \text{Ps}_2^*$  Doppler-broadened linewidth, as has been observed in experiments with Ps [25]).

The statistical significance of the 251 nm resonances in Figs. 1 and 2 can be judged from  $P$ , the probability that a random sample of data from a population representing the null hypothesis gives a worse fit to a wavelength independent constant than the actual measured data. For Fig. 1(b),  $P = 3 \times 10^{-5}$  or  $6 \times 10^{-8}$  for null hypotheses of an unconstrained constant background or a fixed background, respectively, (the latter determined by the Gaussian fit). For the data of Fig. 2(b),  $P = 2 \times 10^{-5}$  or  $6 \times 10^{-14}$  for the null hypothesis of zero background (as determined from the data of Fig. 2(c)), taking into account the chi square per degree of freedom of the Gaussian or double Gaussian fits, respectively. Thus, the spectroscopic data presented in Figs. 1 and 2 demonstrate the existence of the  $L = 1$  excited state predicted by Varga *et al.* [8], confirm our earlier experiment [6] and set the stage for future work in which more detailed measurements will be made.

$\text{Ps}_2$  formed in vacuum by two Ps atoms interacting on a surface will initially carry away some fraction of its binding energy, and so the Doppler width of the excitation line is relatively broad, since the  $\text{Ps}_2$  mass is so low. This limits the accuracy with which measurements of the ( $L = 1$ ) line center can be made to  $\sim 100$  ppm. A measurement to  $\sim 10$  ppm would be necessary to test calculations of relativistic contributions to both the ground and excited state  $\text{Ps}_2$  energies [9]. Improvements to the experimental method, such as using a different or modified formation

target to obtain colder  $\text{Ps}_2$  or employing spectral hole burning techniques [26], might make it possible to reach the 10 ppm level.

An excited state with  $L = 0$  and even C parity has been predicted [27] that suggests the possibility of a Doppler-free two photon excitation that could provide a  $\sim 1$  ppm measurement (determined by the lifetimes of the ground and excited  $\text{Ps}_2$  states). We note, however, that the properties of this state are in question because the group representation [28] used in the calculation may not properly take into account the symmetry of the  $\text{Ps}_2$  system [29].

As indicated in Eq. (2), ionization of the excited state  $\text{Ps}_2$  molecule can result in the production of either positive or negative Ps ions. Although  $\text{Ps}^-$  ions have been observed before [17,18] and can now be produced quite efficiently [19],  $\text{Ps}^+$  ions have not yet been observed experimentally. However, it is likely that some fraction of the signal in the present work is due to  $\text{Ps}^+$  production.

In order to optimize and identify the formation of Ps ions via excited molecular states, one would use a laser just sufficient to ionize  $\text{Ps}_2^*$  (1.96 eV, 633 nm), and electric fields to control the resulting ions [17]. However, 633 nm light will photoionize  $\text{Ps}^\pm$  ions, with a cross section similar to that of  $\text{Ps}_2^*$  molecules [19], so the total Ps ion yield would then be approximately half of the  $\text{Ps}_2^*$  yield. Thus, it should be possible to form a  $\text{Ps}^+$  ion beam with  $\sim 1\%$  efficiency. An interesting prospect is that collisions using slow  $\text{Ps}^+$  ions could provide a means of transferring more than one antiparticle to ordinary matter atoms or molecules. This might be a way to produce and study new exotic multi-positronic systems [30] such as di-positronium hydride,  $\text{Ps}_2\text{H}^-$  [31], or positronic water,  $\text{Ps}_2\text{O}$  [10].

We are grateful to C. W. Clark, D. M. Schrader, and Y. Suzuki for discussions, and L. Liskay and co-workers for providing porous silica samples. This work was supported in part by the U.S. National Science Foundation under Grant No. PHY 0900919 and by the U.S. Air Force Research Office.

- 
- [1] P. A. M. Dirac, *Proc. R. Soc. A* **117**, 610 (1928).
  - [2] C. D. Anderson, *Phys. Rev.* **43**, 491 (1933).
  - [3] A. Rich, *Rev. Mod. Phys.* **53**, 127 (1981).
  - [4] J. A. Wheeler, *Ann. N.Y. Acad. Sci.* **48**, 219 (1946).
  - [5] E. A. Hylleraas and A. Ore, *Phys. Rev.* **71**, 493 (1947).
  - [6] D. B. Cassidy and A. P. Mills, Jr., *Nature (London)* **449**, 195 (2007).
  - [7] D. B. Cassidy, S. H. M. Deng, and A. P. Mills, Jr., *Phys. Rev. A* **76**, 062511 (2007).
  - [8] K. Varga, J. Usukura, and Y. Suzuki, *Phys. Rev. Lett.* **80**, 1876 (1998).
  - [9] M. Puchalski and A. Czarnecki, *Phys. Rev. Lett.* **101**, 183001 (2008).
  - [10] N. Jiang and D. M. Schrader, *Phys. Rev. Lett.* **81**, 5113 (1998).

- [11] D.B. Cassidy, S.H.M. Deng, R.G. Greaves, and A.P. Mills, Jr., *Rev. Sci. Instrum.* **77**, 073106 (2006).
- [12] C.M. Surko and R.G. Greaves, *Phys. Plasmas* **11**, 2333 (2004).
- [13] D.B. Cassidy, T.H. Hisakado, V.E. Meline, H.W.K. Tom, and A.P. Mills, Jr., *Phys. Rev. A* **82**, 052511 (2010).
- [14] D.B. Cassidy and A.P. Mills, Jr., *Phys. Rev. Lett.* **107**, 213401 (2011).
- [15] D.B. Cassidy, S.H.M. Deng, H.K.M. Tanaka, and A.P. Mills, Jr., *Appl. Phys. Lett.* **88**, 194105 (2006).
- [16] J. DiRienzi and R.J. Drachman, *Can. J. Phys.* **88**, 877 (2010).
- [17] A.P. Mills, Jr., *Phys. Rev. Lett.* **46**, 717 (1981).
- [18] H. Ceeh, C. Hugenschmidt, K. Schreckenbach, S.A. Gärtner, P.G. Thirolf, F. Fleischer, and D. Schwalm, *Phys. Rev. A* **84**, 062508 (2011).
- [19] K. Michishio, T. Tachibana, H. Terabe, A. Igarashi, K. Wada, T. Kuga, A. Yagishita, T. Hyodo, and Y. Nagashima, *Phys. Rev. Lett.* **106**, 153401 (2011).
- [20] D.B. Cassidy, M.W.J. Bromley, L.C. Cota, T.H. Hisakado, H.W.K. Tom, and A.P. Mills, Jr., *Phys. Rev. Lett.* **106**, 023401 (2011).
- [21] R.H. Dicke, *Phys. Rev.* **89**, 472 (1953).
- [22] M.S. Fee S. Chu, A.P. Mills, Jr., R.J. Chichester, D.M. Zuckerman, E.D. Shaw, and K. Danzmann, *Phys. Rev. A* **48**, 192 (1993).
- [23] D.B. Cassidy, T.H. Hisakado, H.W.K. Tom, and A.P. Mills, Jr., *Phys. Rev. Lett.* **106**, 133401 (2011).
- [24] A.P. Mills, Jr., *Nucl. Instrum. Methods Phys. Res., Sect. B* **192**, 107 (2002).
- [25] D.B. Cassidy, T.H. Hisakado, H.W.K. Tom, and A.P. Mills, Jr., *Phys. Rev. Lett.* **108**, 043401 (2012).
- [26] E.g., C.J. Foot, *Atomic Physics* (Clarendon Press, Oxford, 2005).
- [27] Y. Suzuki and J. Usukura, *Nucl. Instrum. Methods Phys. Res., Sect. B* **221**, 195 (2004).
- [28] D.B. Kinghorn and R.D. Poshusta, *Phys. Rev. A* **47**, 3671 (1993).
- [29] D.M. Schrader, *Phys. Rev. Lett.* **92**, 043401 (2004).
- [30] D.M. Schrader, in *Physics with Many Positrons*, edited by A. Dupasquier, A.P. Mills, Jr. and R.S. Brusa (IOS, Amsterdam/SIF, Bologna, 2010), p. 337.
- [31] K. Varga, *Phys. Rev. Lett.* **83**, 5471 (1999).