

Comment on “Statistical Mechanics of Developable Ribbons”

Giomi and Mahadevan [1] claim to provide evidence for “an underlying helical structure” of ribbonlike polymers even in the absence of intrinsic curvature (or, in the authors’ words, “of a preferential zero-temperature twist”). They find a persistence length that is over 3 times that of a wormlike chain having the same bending rigidity. We show that these results are an artefact of the authors’ flawed formulation, which produces a helical bias, and that no support for the described effect can be claimed.

The bias stems from the authors’ use of the Frenet frame (FF) in place of a continuous material frame (MF). The FF is known to be an improper choice for ribbons [2]. Although the FF seems natural for developable strips, as the principal normal to the centerline coincides with the normal to the surface, it has the property that it flips at inflection points of the centerline, i.e., the principal normal and binormal change sign, becoming opposite to the continuous vectors of the MF (with ϕ jumping $\phi \rightarrow \phi + \pi$). The centerline of the ribbon is modeled in [1] as a polygonal line with neighboring vertices \mathbf{x}_i at distance a . The shape of the ribbon between \mathbf{x}_{i-1} and \mathbf{x}_i may be inflected or not. If it is, then the MF, if aligned with the FF at \mathbf{x}_{i-1} , will be opposite to the FF at \mathbf{x}_i and its discrete torsion is $\tau^2 = a^{-2}|\mathbf{b}_i + \mathbf{b}_{i-1}|^2 = 2a^{-2}(1 + \cos\phi_i)$. Consequently, the energy of the inflected shape is lower than the energy of the noninflected shape if $\phi_i > \frac{\pi}{2}$. However, this option is ignored in [1], where the MF is always coincident with the FF at \mathbf{x}_i , so it is always assumed that the ribbon segment has no inflection (purely by construction of the discrete chain). Instead, it must be highly bent with a high localized torsion to match the boundary conditions at \mathbf{x}_{i-1} and \mathbf{x}_i . It means that in the continuum limit $a \rightarrow 0$ the torsion diverges. The employed discretization prevents inflected solutions to be approximated.

Inflected conformations are therefore excluded in the Monte Carlo simulations, which generate an ensemble of ribbon shapes with an artificial bias in favor of solutions without inflection points. One should expect the absence of inflected conformations (see Fig. 4) to affect the ensemble’s statistical properties: experimenting with a strip of paper it is easy to produce inflected shapes; they do not cost much energy. Consider two deformations of the strip: a strip locally held in a shape close to a ring and a strip held in an approximate S -shaped form. The center line of the former has no inflection point, contrary to the latter. Clearly the elastic energy is about the same in these two configurations, yet the procedure of [1] completely ignores the inflected S shape.

The bias can be eliminated by using a MF throughout. For instance, sticking with the present angles, rather than through a jump in ϕ one can describe inflection points by allowing θ to take on negative values. The range of the angles would then change to $\theta \in (-\pi, \pi]$, $\phi \in (-\frac{\pi}{2}, \frac{\pi}{2}]$, also used for discrete chains in [3]. This eliminates a possible jump by π in ϕ and hence amounts to using a nonflipping, material, frame. Following the authors’ reasoning one would then assume that $\langle \sin\theta \rangle = 0 = \langle \sin\phi \rangle$, leading to three real eigenvalues of the averaged transfer matrix \mathbf{R} and thus exponential decay of the tangent-tangent correlation function $\langle \mathbf{t}_0 \cdot \mathbf{t}_n \rangle$, in contradiction to the oscillatory decay shown in Fig. 3.

The authors also present an analytical calculation that confirms their numerical results by producing a persistence length l_p for developable ribbons that is more than 3 times that of a wormlike chain of the same bending rigidity. However, for the redefined angles, with area element $d\Omega = |\sin\theta|d\theta d\phi$, following the computations in [1] we arrive at a modified partition function, Eq. (9), in which $I_0(\beta\mu)$ is replaced by $I_0(\beta\mu) + L_0(\beta\mu)$, where L_0 is the modified Struve function of 0th order. For $\zeta \gg 1$, $L_0(\zeta) = I_0(\zeta) - \frac{2}{\pi\zeta} + \mathcal{O}(\zeta^{-3})$, which implies that the asymptotic expression for $\langle \cos\theta \rangle$ on p. 4 (left, middle) of [1] remains valid. Therefore, we find $l_p = -a(\log\langle \cos\theta \rangle)^{-1} \simeq \frac{32}{35} \frac{Dw}{k_B T}$, which is slightly lower than for the wormlike chain model.

We conclude that there is no ground for the claim in [1] that developable ribbons possess an “underlying helical structure”. This conclusion is also in agreement with calculations in [4], where a properly discretized developable ribbon is found to yield an exponentially decaying tangent-tangent correlation function, consistent with our result above.

E. L. Starostin and G. H. M. van der Heijden
Centre for Nonlinear Dynamics
University College London
Gower Street, London WC1E 6BT, United Kingdom

Received 28 June 2010; published 2 December 2011

DOI: 10.1103/PhysRevLett.107.239801

PACS numbers: 87.10.Pq, 82.35.Lr, 82.35.Pq

- [1] L. Giomi and L. Mahadevan, *Phys. Rev. Lett.* **104**, 238104 (2010).
- [2] S. M. Rappaport and Y. Rabin, *J. Phys. A* **40**, 4455 (2007).
- [3] S. Hu, M. Lundgren, and A. J. Niemi, *Phys. Rev. E* **83**, 061908 (2011).
- [4] B. Mergell, M. R. Ejtehadi, and R. Everaers, *Phys. Rev. E* **66**, 011903 (2002).