Centrality Dependence of the Charged-Particle Multiplicity Density at Midrapidity in Pb-Pb Collisions at $\sqrt{s_{\mathrm{NN}}} = 2.76$ TeV

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The centrality dependence of the charged-particle multiplicity density at midrapidity in Pb-Pb collisions at $\sqrt{s_{\mathrm{NN}}} = 2.76$ TeV is presented. The charged-particle density normalized per participating nucleon pair increases by about a factor of 2 from peripheral (70%–80%) to central (0%–5%) collisions. The centrality dependence is found to be similar to that observed at lower collision energies. The data are compared with models based on different mechanisms for particle production in nuclear collisions.

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Quantum chromodynamics (QCD), the theory of the strong interaction, predicts a phase transition at high temperature between hadronic and deconfined matter (the quark-gluon plasma). Strongly interacting matter under such extreme conditions can be studied experimentally using ultrarelativistic collisions of heavy nuclei. The field entered a new era in November 2010 when the Large Hadron Collider (LHC) at CERN produced the first Pb-Pb collisions at a center-of-mass energy per nucleon pair $\sqrt{s_{\mathrm{NN}}} = 2.76$ TeV. This represents an increase of more than 1 order of magnitude over the highest energy nuclear collisions previously obtained in the laboratory.

The multiplicity of charged particles produced in the central rapidity region is a key observable to characterize the properties of the matter created in these collisions [1]. Nuclei are extended objects, and their collisions can be characterized by centrality, related to the collision impact parameter. The study of the dependence of the charged-particle density on colliding system, center-of-mass energy and collision geometry is important to understand the relative contributions to particle production of hard scattering and soft processes, and may provide insight into the partonic structure of the projectiles.

The ALICE Collaboration recently reported the measurement of the charged-particle pseudorapidity density at midrapidity for the most central (head-on) Pb-Pb collisions at $\sqrt{s_{\mathrm{NN}}} = 2.76$ TeV [2]. In this Letter, we extend that study to noncentral collisions, presenting the measurement of the centrality dependence of the multiplicity density of charged primary particles $dN_{\mathrm{ch}}/d\eta$ in the pseudorapidity interval $|\eta| < 0.5$. The pseudorapidity is defined as $\eta \equiv -\ln\tan(\theta/2)$, where θ is the angle between the charged-particle direction and the beam axis (z). Primary particles are defined as all prompt particles produced in the collision, including decay products, except those from weak decays of strange particles.

We report the charged-particle density per participant pair, $(dN_{\mathrm{ch}}/d\eta)/(\langle N_{\mathrm{part}}\rangle/2)$, for nine centrality classes, covering the most central 80% of the hadronic cross section. The average number of nucleons participating in the collision in a given centrality class, $\langle N_{\mathrm{part}}\rangle$, reflects the collision geometry and is obtained using Glauber modeling [3]. The results are compared with measurements at lower collision energy [4–9] and with theoretical calculations [10–14].

The data for this measurement were collected with the ALICE detector [15]. The data sample is the same as in [2] and the analysis techniques are similar. The main detector utilized in the analysis is the silicon pixel detector (SPD), the innermost part of the inner tracking system (ITS). The SPD consists of two cylindrical layers of hybrid silicon pixel assemblies covering $|\eta| < 2.0$ and $|\eta| < 1.4$ for the inner and outer layers, respectively. A total of 9.8×10^6 pixels of size 50×425 $\mu\mathrm{m}^2$ are read out by 1200 electronic chips. Each chip also provides a fast signal when at least one of its pixels is hit. These signals are combined in a programmable logic unit which supplies a trigger signal. A trigger signal is also provided by the VZERO counters, two arrays of 32 scintillator tiles covering the full azimuth within $2.8 < \eta < 5.1$ (VZERO-A) and $-3.7 < \eta < -1.7$ (VZERO-C). The trigger was configured for high efficiency for hadronic events, requiring at least two out of the following three conditions: (i) two pixel chips hit in the outer layer of the SPD, (ii) a signal in VZERO-A, (iii) a signal in VZERO-C. The threshold in the VZERO detector corresponds approximately to the energy deposition of a minimum ionizing particle. This trigger configuration led to a rate of about 50 Hz, with 4 Hz from nuclear interactions, 45 Hz from electromagnetic processes, and 1 Hz arising from beam background. In addition, in the offline event selection, we also use the information from two neutron zero degree calorimeters (ZDCs) positioned at ± 114 m from the interaction point. Beam background

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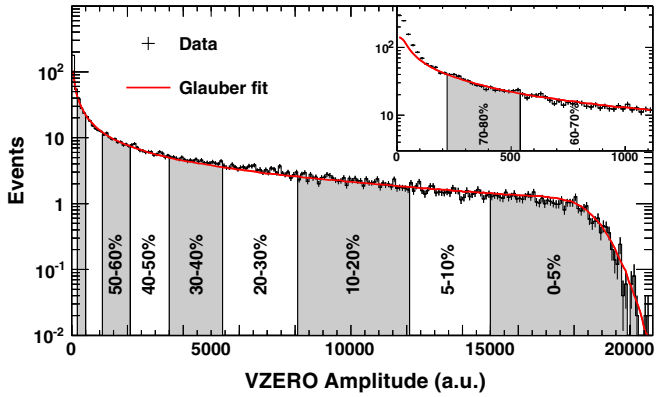

FIG. 1 (color online). Distribution of the summed amplitudes in the VZERO scintillator tiles (histogram); inset shows the low amplitude part of the distribution. The curve shows the result of the Glauber model fit to the measurement. The vertical lines separate the centrality classes used in the analysis, which in total correspond to the most central 80% of hadronic collisions.

events are removed using the VZERO and ZDC timing information. Electromagnetically induced interactions are reduced by requiring an energy deposition above 500 GeV in each of the neutron ZDCs.

After event selection, the sample consists of about blue 65 000 events. Figure 1 shows the distribution of the summed amplitudes in the VZERO scintillator tiles together with the distribution obtained with a model of particle production based on a Glauber description of nuclear collisions [3]. We use a two-component model assuming that the number of particle-producing sources is given by $f \times N_{\rm part} + (1-f) \times N_{\rm coll}$, where $N_{\rm part}$ is the number of participating nucleons, $N_{\rm coll}$ is the number of binary nucleon-nucleon collisions and f quantifies their relative contributions. The number of particles produced by each source is distributed according to a negative binomial distribution, parametrized with μ and κ, where μ is the mean multiplicity per source and κ controls the large multiplicity tail.

In the Glauber calculation [16], the nuclear density for ^{208}Pb is modeled by a Woods-Saxon distribution for a spherical nucleus with a radius of 6.62 fm and a skin depth of 0.546 fm, based on data from low energy electron-nucleus scattering experiments [17]. A hard-sphere exclusion distance of 0.4 fm between nucleons is employed. Nuclear collisions are modeled by randomly displacing the two colliding nuclei in the transverse plane. Nucleons from each nucleus are assumed to collide if the transverse distance between them is less than the distance corresponding to the inelastic nucleon-nucleon cross section, estimated from interpolating data at different center-of-mass energies [18] to be 64 ± 5 mb at $\sqrt{s} = 2.76$ TeV

The values of f, μ, and κ are obtained from a fit to the measured VZERO amplitude distribution. The fit is restricted to amplitudes above a value corresponding to 88% of the hadronic cross section. In this region the trigger and event selection are fully efficient, and the contamination by electromagnetic processes is negligible. Centrality classes are determined by integrating the measured distribution above the cut, as shown in Fig. 1.

The determination of $dN_{\rm ch}/d\eta$ is performed for each centrality class. The primary vertex position is extracted by correlating hits in the two SPD layers. All events in the sample corresponding to 0%–80% of the hadronic cross section are found to have a well-defined primary vertex. To minimize edge effects at the limit of the SPD acceptance, we require $|z_{\rm vtx}| < 7$ cm for the reconstructed vertex, leading to a sample of about 49 000 events.

The measurement of the charged-particle multiplicity is based on the reconstruction of tracklets [2]. A tracklet candidate is defined as a pair of hits, one in each SPD layer. Using the reconstructed vertex as the origin, differences in azimuthal ($\Delta\varphi$, bending plane) and polar ($\Delta\theta$, nonbending direction) angles for pairs of hits are calculated [19]. Tracklets are defined by hit combinations that satisfy a selection on the sum of the squares (δ^2) of $\Delta\varphi$ and $\Delta\theta$, each normalized to its estimated resolution (60 mrad for $\Delta\varphi$ and $25\sin^2\theta$ $m$rad for $\Delta\theta$). The tolerance in $\Delta\varphi$ for tracklet reconstruction effectively selects charged particles with transverse momentum above 50 MeV/c. If multiple tracklet candidates share a hit, only the combination with the smallest δ^2 is kept.

The charged-particle pseudorapidity density $dN_{\rm ch}/d\eta$ in $|\eta| < 0.5$ is obtained from the number of tracklets by applying a correction $\alpha \times (1-\beta)$ in bins of pseudorapidity and z position of the primary vertex. The factor α corrects for the acceptance and efficiency of a primary track to form a tracklet, and β reflects the fraction of background tracklets from uncorrelated hits. The fraction β is estimated by matching the tails of the data and background δ^2 distributions. The latter is obtained by selecting combinatorial tracklets from a sample of simulated events with similar SPD hit multiplicities generated with HIJING [20] and a GEANT3 [21] model of the detector response. The estimated background fraction varies from 1% in the most peripheral to 14% in the most central class.

The correction α is obtained as the ratio of the number of generated primary charged particles and the number of reconstructed tracklets, after subtraction of the combinatorial background. Thus, α includes the corrections for the geometrical acceptance, detector and reconstruction inefficiencies, contamination by weak decay products of strange particles, photon conversions, secondary interactions, and undetected particles with transverse momentum below 50 MeV/c. The correction is about 1.8 and varies little with centrality. Its magnitude is dominated by the effect of tracklet acceptance: the fraction of SPD channels active during data taking was 70% for the inner and 78% for the outer layer.

Systematic uncertainties on $dN_{\rm ch}/d\eta$ are estimated as follows: for *background subtraction*, from 0.1% in the

most peripheral to 2.0% in the most central class, by using an alternative method where fake hits are injected into real events; for *particle composition*, 1%, by changing the relative abundances of protons, pions, kaons by up to a factor of 2; for contamination by weak decays, 1%, by changing the relative contribution of the yield of strange particles by a factor of 2; for extrapolation to zero transverse momentum, 2%, by varying the estimated yield of particles at low transverse momentum by a factor of 2; for dependence on event generator, 2%, by using quenched and unquenched versions of HIJING [20], as well as DPMJET [22] for calculating the corrections. The systematic uncertainty on $dN_{\rm ch}/d\eta$ due to the centrality class definition is estimated as 6.2% for the most peripheral and 0.4% for the most central class, by using alternative centrality definitions based on track or SPD hit multiplicities, by using different ranges for the Glauber model fit, by defining cross-section classes integrating over the fit rather than directly over the data distributions, by changing the $N_{\rm part}$ dependence of the particle production model to a power law, and by changing the nucleon—nucleon cross section and the parameters of the Woods—Saxon distribution within their estimated uncertainties and by changing the internucleon exclusion distance by ±100%. All other sources of systematic errors considered (tracklet cuts, vertex cuts, material budget, detector efficiency, background events) were found to be negligible. The total systematic uncertainty on $dN_{\rm ch}/d\eta$ amounts to 7.0% in the most peripheral and 3.8% in the most central class. A large part of this uncertainty, about 5.0% for the most peripheral and 2.5% for the most central class, is correlated among the different centrality classes. The $dN_{\rm ch}/d\eta$ values obtained for nine centrality classes together with their systematic uncertainties are given in Table I. As a cross check of the centrality selection the $dN_{\rm ch}/d\eta$ analysis was repeated using centrality cuts defined by slicing perpendicularly to the correlation between the energy deposited in the ZDC and the VZERO amplitude. The resulting $dN_{\rm ch}/d\eta$ values differ by 3.5% in the most peripheral (70%–80%) and by less than 2% in all the other classes from those obtained by using the VZERO selection alone, which is well within the systematic uncertainty. Independent cross checks performed using tracks reconstructed in the TPC and ITS instead of tracklets yield compatible results.

In order to compare bulk particle production in different collision systems and at different energies, the charged-particle density is divided by the average number of participating nucleon pairs, $\langle N_{\rm part}\rangle/2$, determined for each centrality class. The $\langle N_{\rm part}\rangle$ values are obtained using the Glauber calculation, by classifying events according to the impact parameter, without reference to a specific particle production model, and are listed in Table I. The systematic uncertainty in the $\langle N_{\rm part}\rangle$ values is obtained by varying the parameters entering the Glauber calculation as described above. The geometrical $\langle N_{\rm part}\rangle$ values are consistent within uncertainties with the values extracted from the Glauber fit in each centrality class, and agree to better than 1% except for the 70–80% class where the difference is 3.5%.

Figure 2 presents $(dN_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$ as a function of the number of participants. Point-to-point, uncorrelated uncertainties are indicated by the error bars, while correlated uncertainties are shown as the grey band. Statistical errors are negligible. The charged-particle density per participant pair increases with $\langle N_{\rm part}\rangle$, from 4.4 ± 0.4 for the most peripheral to 8.4 ± 0.3 for the most central class. The values for Au-Au collisions at $\sqrt{s_{\rm NN}} = 0.2$ TeV, averaged over the RHIC experiments [7], are shown in the same

TABLE I. $dN_{\rm ch}/d\eta$ and $(dN_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$ values measured in $|\eta| < 0.5$ for nine centrality classes. The $\langle N_{\rm part}\rangle$ obtained with the Glauber model are given.

Centrality	$dN_{\rm ch}/d\eta$	$\langle N_{\rm part}\rangle$	$(dN_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$
0%–5%	1601 ± 60	382.8 ± 3.1	8.4 ± 0.3
5%–10%	1294 ± 49	329.7 ± 4.6	7.9 ± 0.3
10%–20%	966 ± 37	260.5 ± 4.4	7.4 ± 0.3
20%–30%	649 ± 23	186.4 ± 3.9	7.0 ± 0.3
30%–40%	426 ± 15	128.9 ± 3.3	6.6 ± 0.3
40%–50%	261 ± 9	85.0 ± 2.6	6.1 ± 0.3
50%–60%	149 ± 6	52.8 ± 2.0	5.7 ± 0.3
60%–70%	76 ± 4	30.0 ± 1.3	5.1 ± 0.3
70%–80%	35 ± 2	15.8 ± 0.6	4.4 ± 0.4

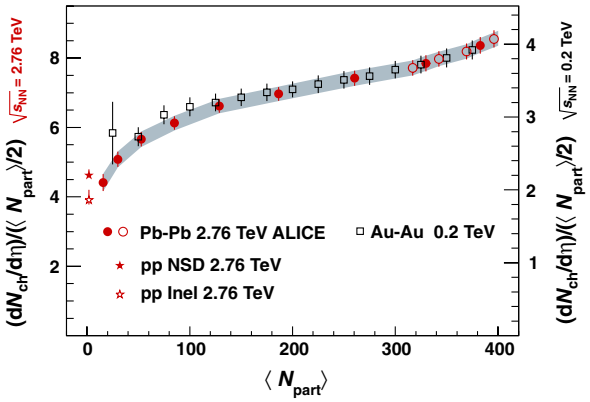

FIG. 2 (color online). Dependence of $(dN_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$ on the number of participants for Pb-Pb collisions at $\sqrt{s_{\rm NN}} = 2.76$ TeV and Au-Au collisions at $\sqrt{s_{\rm NN}} = 0.2$ TeV (RHIC average) [7]. The scale for the lower-energy data is shown on the right-hand side and differs from the scale for the higher-energy data on the left-hand side by a factor of 2.1. For the Pb-Pb data, uncorrelated uncertainties are indicated by the error bars, while correlated uncertainties are shown as the grey band. Statistical errors are negligible. The open circles show the values obtained for centrality classes obtained by dividing the 0%–10% most central collisions into four, rather than two classes. The values for non-single-diffractive and inelastic pp collisions are the results of interpolating between data at 2.36 [19,24] and 7 TeV [25].

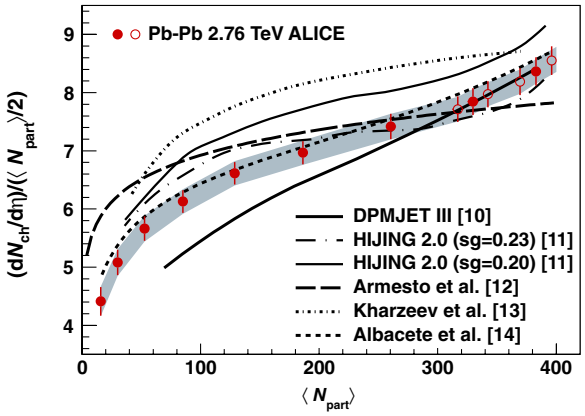

FIG. 3 (color online). Comparison of $(dN_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$ with model calculations for Pb-Pb at $\sqrt{s_{\rm NN}} = 2.76$ TeV. Uncertainties in the data are shown as in Fig. 2. The HIJING 2.0 curve is shown for two values of the gluon shadowing (s_g) parameter.

figure with a scale that differs by a factor of 2.1 on the right-hand side. The centrality dependence of the multiplicity is found to be very similar for $\sqrt{s_{\rm NN}} = 2.76$ TeV and $\sqrt{s_{\rm NN}} = 0.2$ TeV.

Theoretical descriptions of particle production in nuclear collisions fall into two broad categories: two-component models combining perturbative QCD processes (e.g., jets and mini jets) with soft interactions, and saturation models with various parametrizations for the energy and centrality dependence of the saturation scale. In Fig. 3 we compare the measured $(N_{\rm ch}/d\eta)/(\langle N_{\rm part}\rangle/2)$ with model predictions. A calculation based on the two-component Dual Parton Model (DPMJET [10], with string fusion) exhibits a stronger rise with centrality than observed. The two-component HIJING 2.0 model [23], which has been tuned [11] (Published after the most central $dN_{\rm ch}/d\eta$ value [2] was known.) to high-energy pp [19,24] and central Pb-Pb data [2], reasonably describes the data. This model includes a strong impact parameter dependent gluon shadowing (s_g) which limits the rise of particle production with centrality. The remaining models show a weak dependence of multiplicity on centrality. They are all different implementations of the saturation picture, where the number of soft gluons available for scattering and particle production is reduced by nonlinear interactions and parton recombination. A geometrical scaling model with a strong dependence of the saturation scale on nuclear mass and collision energy [12] predicts a rather weak variation with centrality. The centrality dependence is well reproduced by saturation models [13,14] (Published after the most central $dN_{\rm ch}/d\eta$ value [2] was known.), although the former overpredicts the magnitude.

In summary, the measurement of the centrality dependence of the charged-particle multiplicity density at midrapidity in Pb-Pb collisions at $\sqrt{s_{\rm NN}} = 2.76$ TeV has been presented. The charged-particle density normalized per participating nucleon pair increases by about a factor 2 from peripheral (70%–80%) to central (0%–5%) collisions. The dependence of the multiplicity on centrality is strikingly similar for the data at $\sqrt{s_{\rm NN}} = 2.76$ TeV and $\sqrt{s_{\rm NN}} = 0.2$ Theoretical descriptions that include a moderation of the multiplicity evolution with centrality are favored by the data.

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