

## Double Core-Hole Production in N<sub>2</sub>: Beating the Auger Clock

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We investigate the creation of double *K*-shell holes in N<sub>2</sub> molecules via sequential absorption of two photons on a time scale shorter than the core-hole lifetime by using intense x-ray pulses from the Linac Coherent Light Source free electron laser. The production and decay of these states is characterized by photoelectron spectroscopy and Auger electron spectroscopy. In molecules, two types of double core holes are expected, the first with two core holes on the same N atom, and the second with one core hole on each N atom. We report the first direct observations of the former type of core hole in a molecule, in good agreement with theory, and provide an experimental upper bound for the relative contribution of the latter type.

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The development of intense, short pulse, x-ray free electron lasers (FELs) will allow the exploration of novel states of atoms, molecules and clusters, with potential impact on applications ranging from single-pulse imaging of biomolecules to high-energy density materials [1–5]. One intriguing possibility enabled by these sources is the ability to produce atoms and molecules with multiple electron vacancies in core orbitals through the sequential absorption of multiple photons on a time scale faster than Auger decay, and it has been suggested that double core holes (DCHs) could provide the basis for richer, sensitive, spectroscopies than conventional inner-shell photoelectron spectroscopy [6,7]. The different atomic sites in molecules introduce multiple possibilities for the DCH configurations, e.g., DCHs with both vacancies on a single site (DCHSS) and DCHs with single vacancies on two different sites (DCHTS). For such states, the presence and location of the first hole is predicted to affect the energy to produce the second hole, as well as the decay mechanisms and fragmentation patterns of the resulting DCH states [6]. The magnitude of the energy shifts of these states will provide unique spectral signatures as well as new information on the chemical environment of the core holes [7,8]. While the intensity of conventional x-ray sources is too low to produce observable DCHs through sequential absorption processes, DCH states have been observed as a

result of electron correlation in single-photon absorption. Unfortunately, these processes have a very low yield [9,10], and only DCHSS states have been observed in this manner. At the intensities of the new x-ray FELs, photoabsorption should compete effectively with Auger decay [11,12], allowing the production of both types of DCH states through sequential two-photon processes.

In this Letter, we report the first experimental attempt to characterize the DCH states of a molecule, N<sub>2</sub>, by sequential two-photon absorption from an x-ray FEL, the LCLS. The production and decay of these states are characterized by using photoelectron spectroscopy and Auger electron spectroscopy. The experimental results are interpreted with the help of *ab initio* calculations [13,14] of the singly, doubly, and triply ionized states and their corresponding Auger spectra [15]. Both Green's function calculations and multireference configuration interaction (aug-cc-pCVQZ basis set) calculations were performed. These results will serve as a basis for understanding double- and multiple-core-hole states in more complex molecules. The experiment was conducted using the LCLS Atomic Molecular and Optical (AMO) physics instrument. We used photon energies of 1.0 keV and 1.1 keV, a pulse duration of 280 fs, and a fluence of  $\sim 7 \times 10^4$  photon  $\cdot \text{Å}^{-2}$  (the fluence at the beam focus is estimated to be  $\sim 1 \times 10^4$  photon  $\cdot \text{Å}^{-2}$  per pulse due to the photon beam line losses). The details of



(shakeup) SCH states. Above  $\sim 400$  eV, no Auger electrons of  $N_2$  have been observed with conventional x-ray sources. In contrast, the new spectra shown in Fig. 2 exhibit two new features lying 50–80 eV above the main SCH Auger peaks and with kinetic energies of 413 eV and 442 eV. These signals are clearly associated with Auger processes since the kinetic energies of the peaks do not vary with the photon energy. Contributions from photoelectrons observed in Fig. 1 are strongly suppressed in Fig. 2 due to the selected electron emission direction perpendicular to the light polarization. Our calculation shown in Fig. 2 (yellow [light grey] shaded curve) predicts peaks at  $\sim 410$  eV, which are produced by the Auger decay of the main DCHSS state of  $N_2^{2+}$  involving mainly  $2\sigma_u$  and  $1\pi_u$  electrons. In light of the expected underestimate in theory of 2–4 eV due to partially unaccounted electron relaxation in the final states [23], the prediction agrees very well with the observed peak position. We conclude that the newly observed Auger electron signal at  $\sim 410$  eV kinetic energy stems from the decay of DCHSS states and is associated with so-called Auger hypersatellites [9,10]. We estimate the DCHSS signal intensity to be  $\sim 1\%$  of the main Auger peak signal between 355 eV–370 eV ( $\pm 5$  eV). Experiments have also been performed at the LCLS to characterize the angular distribution of these hypersatellite electrons [24].

Although calculations of the Auger decay of DCHSS shakeup states have not been performed explicitly, the energy splitting between the DCHSS main Auger peak and the DCHSS shakeup Auger peaks can be estimated by the energy difference between the calculated DCHSS photoline and the DCHSS shakeup photoline. The predicted energies of the DCHSS shakeup Auger electrons are indicated by black stem lines in Fig. 2. The good agreement of the calculated positions with the peak at 442 eV suggests that this peak might correspond to the Auger decay of the DCHSS shakeup states.

The detection and assignment of electrons associated with the production of the uniquely molecular DCHTS states requires improved energy resolution, which is achieved by increasing the retardation voltage in the photoelectron spectrometer to 480 V, as shown in Fig. 3. The spectrum is analyzed by taking into account the contributions from single-photon shakeup or shakeoff (SUO) processes as previously determined by Svensson *et al.* [25] and Kaneyasu *et al.* [26], and convoluting them with the experimental energy resolution of the current study. The dashed lines and blue (dark grey) shaded area in Fig. 3 indicate the shape and the range of uncertainty of the SUO contributions, respectively, based on the uncertainties in assigning exact peak areas in the spectrum of Svensson *et al.* [25]. Note that in single-photon ionization, SUO processes would generate the only contributions to the photoelectron spectrum in the kinetic energy range between  $\sim 400$  eV and  $\sim 580$  eV. Therefore, all photoelec-

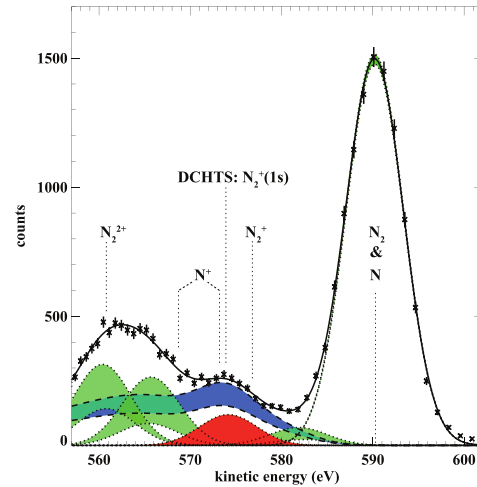


FIG. 3 (color online). Photoelectron spectrum (crosses with error bars) recorded at 1.0 keV photon energy and 480 V spectrometer retardation voltage. The blue, green, and red shaded areas mark the correlated uncertainties for various spectral components (see text). The constant sum of all components is marked by the solid black curve. Calculated line positions are marked by vertical dashed lines. The red [medium grey] shaded area allows for an upper limit estimate of DCHTS contributions.

tron signal shown in Fig. 3 that is not contained in the main photoline at 590 eV or the SUO contributions is due to multiple photon processes that are enabled by the high intensity of the LCLS. To disentangle the major contributions generated by multiple photon ionization, a nonlinear least-squares fit is performed based on five spectral components with identical peak shapes and independent amplitudes and center energies (green [light grey] and red [medium grey] shaded areas). The common peak shape is derived by a description of the well-defined main photoline by a Gaussian function, which gives a good approximation of the light-source- and apparatus-induced line broadening effects.

The range of possible SUO intensities results in a corresponding uncertainty range for the intensities of all other spectral contributions, as indicated by the green (light grey) and red (medium grey) shaded areas, the borders of which indicate the upper and lower limits. However, all peak center energies and the overall fit quality are virtually unaffected by the SUO uncertainty. Within the model of identical, apparatus-limited peak shapes, the fit result is globally unique as determined by a Monte Carlo based sampling of the available fit starting parameter space by  $\sim 10^5$  independent fit cycles. We therefore base our spectral analysis on the comparison of the fit results with calculated line positions, as marked in Fig. 3 by vertical dashed lines. The theoretical peak positions mark the kinetic energies expected from  $1s$  inner-shell ionization of neutral and valence-ionized atoms and molecules as indicated. Also marked is the expected position of the DCHTS photoelectron line that results from inner-shell

ionization of an  $N_2$  ( $1s^{-1}$ ) molecular core-hole state. We note that subtleties such as dynamic line shifts during molecular fragmentation are not included in the fit model. However, these are considered minor effects with respect to the conclusions that are derived here.

Taking into account the finite precision of the calculations ( $\sim 1$ – $2$  eV) and the finite energy resolution of the measurements, the agreement of the experimentally derived peak positions with the theoretical expectations is mostly very good. Independent, physically motivated cross-checks, such as the analysis of the ratio of valence-to inner-shell ionization probabilities and the rate of molecular dissociation compared to the pulse length, confirm that the measured peak intensities for inner-shell ionization of  $N_2^{2+}$  and  $N^+$  agree qualitatively with the picture of sequential multiple photon ionization by the LCLS at the peak intensity of the experiments. The weak contribution at  $\sim 582$  eV cannot readily be assigned to a calculated line. Slight deviations of the mainline peak shape from a perfect Gaussian, as can be seen at the high-energy side of the main photoline, may account for some part of this feature. Some signal from  $N_2^+$  ionization may also contribute to this part of the photoelectron spectrum.

One spectral component is consistently found exactly at the position of the expected DCHTS photoline (574 eV). Unfortunately, this spectral region is also marked by the strongest SUO contributions. Additionally, the  $1s^{-1}$  photoline from inner-shell ionization of excited  $N^+$  ( $1s^2 2s^1 2p^3$ ) fragments coincides with the DCHTS photoline within the theoretical and experimental uncertainties. As indicated by the red (medium grey) area, the estimate of the total multiple photon ionization signal at 574 eV varies between zero and 8% of the main photoline intensity. Given the uncertainty in the contributions from overlapping spectral components, we can provide only an upper bound for the DCHTS signal rather than an absolute value. The upper limit of the DCHTS contribution is mainly defined by the minimum contribution of the  $N^+$  ( $1s^2 2s^1 2p^3$ ) ionization signal. The lower bound of this signal is estimated by the peak at 566 eV, which is mainly attributed to inner-shell ionization of  $N^+$  ( $1s^2 2s^2 2p^2$ ), and by the branching ratio between ground and excited states of  $N^+$  fragments after inner-shell ionization of  $N_2$  [27]. This analysis leads to an upper bound for DCHTS contributions in the photoelectron spectrum of 4% relative to the intensity of the  $N_2$  ( $1s^{-1}$ ) main photoline. This value is in qualitative agreement with the  $\sim 1\%$  DCHSS signal described above and an estimated ratio of  $\sim 1.65$  between DCHTS and DCHSS signal intensities. An independent estimate of the intensity of the DCHTS process by examination of the Auger spectrum is not possible at this stage. The nature of these states marked by a single core hole per atom leads to Auger electron energies that are embedded in the SCH Auger spectrum, as indicated by the grey shaded structure near 350 eV kinetic energy in Fig. 2. Furthermore, Auger decay

of sequential multiple photoionization products, in addition to correlation satellites and the decay of SUO states, may also contribute to this range of the Auger electron spectrum, making it difficult to isolate the DCHTS Auger processes at this time. Note that in Fig. 2 only the high-energy, most intense, bands of the computed DCHTS and DCHSS Auger spectra are shown, involving ionization of the outer-valence electrons. The low-energy band of the DCHTS spectrum is calculated to fall outside the experimental range, while the DCHSS one falls in the 380–390 eV shakeup region where it is not appreciably detected by the experiment. The theoretical determination of the relative probabilities for SCH and DCH formation requires complex calculations that we hope to stimulate with the present work. Thus, no comparison between experiment and theory has been made on this point.

In conclusion, we have presented the first experimental results on the characterization of DCH states produced by sequential two-photon absorption in molecules. The observation of DCHSS states is supported by photoelectron and Auger electron spectra, as well as by theoretical calculations. In contrast, the observation of the uniquely molecular DCHTS states remains ambiguous, and only an upper bound to their production can be determined. Higher intensity (afforded by improved beam transport optics), improved stability of photon energy, higher electron energy resolution, and detailed studies of the intensity dependence of the electron spectra will thus be required to fulfill the potential of DCH spectroscopies.

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