

Bond Stretching Phonon Softening and Kinks in the Angle-Resolved Photoemission Spectra of Optimally Doped $\text{Bi}_2\text{Sr}_{1.6}\text{La}_{0.4}\text{Cu}_2\text{O}_{6+\delta}$ Superconductors

J. Graf,¹ M. d'Astuto,² C. Jozwiak,³ D. R. Garcia,^{1,3} N. L. Saini,⁴ M. Krisch,⁵ K. Ikeuchi,⁶
A. Q. R. Baron,⁷ H. Eisaki,⁸ and A. Lanzara^{1,3,*}

¹Materials Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA

²Institut de Minéralogie et de Physique des Milieux Condensés (IMPMC), Université Pierre et Marie Curie–Paris 6, case 115, 4, place Jussieu, 75252 Paris cedex 05, France[†]

³Department of Physics, University of California Berkeley, California 94720, USA

⁴Dipartimento di Fisica, Università di Roma La Sapienza, 00185 Roma, Italy

⁵European Synchrotron Radiation Facility, F-38043 Grenoble Cédex, France

⁶Spring-8/JAEA, 1-1-1 Kouto, Sayo, Hyogo, 670 Japan

⁷Spring-8/RIKEN&JASRI, 1-1-1 Kouto, Sayo, Hyogo, 670 Japan

⁸AIST Tsukuba Central 2, Umezono, Tsukuba, Ibaraki, 305-8568, Japan

(Received 28 September 2007; published 4 June 2008)

We report the first measurement of the Cu-O bond stretching phonon dispersion in optimally doped $\text{Bi}_2\text{Sr}_{1.6}\text{La}_{0.4}\text{Cu}_2\text{O}_{6+\delta}$ using inelastic x-ray scattering. We found a softening of this phonon at $\mathbf{q} = (\approx 0.25, 0, 0)$ from 76 to 60 meV, similar to the one reported in other cuprates. A comparison with angle-resolved photoemission data on the same sample revealed an excellent agreement in terms of energy and momentum between the angle-resolved photoemission nodal kink and the *soft* part of the bond stretching phonon. Indeed, we find that the momentum space where a 63 ± 5 meV kink is observed can be connected with a vector $\mathbf{q} = (\xi, 0, 0)$ with $\xi \geq 0.22$, corresponding exactly to the soft part of the bond stretching phonon.

DOI: [10.1103/PhysRevLett.100.227002](https://doi.org/10.1103/PhysRevLett.100.227002)

PACS numbers: 74.72.Hs, 74.25.Jb, 74.25.Kc

To fully uncover the nature of the electron-phonon interaction in a material, it is important to directly map the electron and phonon dispersion. This is possible by performing angle-resolved photoemission spectroscopy (ARPES) and inelastic neutron scattering (INS) or x-ray scattering (IXS) experiments. For example, when a phonon wave vector matches $2\mathbf{k}_F$, where \mathbf{k}_F is the Fermi momentum, a Kohn anomaly or a charge density wave can result from the enhancement of the electron-phonon coupling due to the nesting of the Fermi surface. The resulting renormalization of the phonon dispersions can be studied using INS or IXS. Similarly, the coupling of electrons to bosons is manifested in ARPES by a kink in the electronic dispersion.

INS and IXS data have shown an anomalous softening of the Cu-O bond stretching (BS) mode at the metal-insulator phase transition in cuprates superconductors [1–6] as well as nonsuperconducting perovskites [7,8]. Whether this softening is related to superconductivity or to the strong charge order instabilities present in some of these compounds is not clear, since neither of these two phenomena are ubiquitous to all of these compounds.

ARPES data have shown a kink in the electronic dispersion at 60–70 meV along the nodal direction in various cuprates [9], which shifts to lower binding energy (30–40 meV) as the Brillouin zone (BZ) face is approached [10–12]. Because of the coinciding energy, the nodal kink has been associated with the BS phonon [9]. However, an alternative explanation has been suggested [13,14]. Therefore, what is the nature of the nodal kink, what determines the momentum region where it exists, and

what are the signatures of the phonon softening on the electronic structure are still open questions.

A combination of IXS and ARPES measurements on the same material is therefore very valuable in establishing the relation, if any, between these two behaviors. Such a study has been missing thus far. The best system to perform this cross analysis is the single layer $\text{Bi}_2\text{Sr}_{1.6}\text{La}_{0.4}\text{Cu}_2\text{O}_{6+\delta}$ superconductor (Bi2201) as no report of a magnetic resonance mode exists, allowing one to draw a direct comparison between the electronic structure and the phonon spectrum. Also, the Bi-based cuprates are among the most studied materials by ARPES because of the good quality of their cleaved surface. On the other hand, the lack of large single crystals has made INS experiment difficult and, though IXS can probe submillimeter crystals, this is a challenging experiment due to the very low inelastic cross section of the BS mode [15].

We report the first measurements of the dispersion of the longitudinal phonons in optimally doped Bi2201. We find a strong renormalization of the BS phonon dispersion along the $\mathbf{q} = (\xi, 0, 0)$ direction, similar to the one reported in several cuprates. ARPES measurements revealed the presence of two distinct kink structures near the node and the BZ face at 63 and 35 meV, respectively, similar to the two kink structures reported in double layer $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ [10,11] and $\text{La}_{2-x}\text{Sr}_x\text{Cu}_2\text{O}_4$ [12].

From the comparison between ARPES and IXS measurements on the same sample, we show that the *soft* part of the BS phonon matches the energy and momentum of the 63 meV kink.

The single crystals of Bi2201 ($T_c = 33$ K) were grown using the traveling-solvent floating-zone technique [16]. The IXS experiment was performed at beam line BL35XU at SPring-8 in Japan and at beam line ID28 at the European Synchrotron Radiation Facility (ESRF) in France. The momentum and energy resolutions were better than 0.08 \AA^{-1} and 3.2 meV , respectively. The measurements were done in the transmission geometry at low temperature (≈ 10 K). The sample thickness was $\approx 30 \text{ \mu m}$, allowing $\approx 30\%$ of x-ray transmission.

IXS energy scans were carried out at $\mathbf{Q} = \mathbf{G} + \mathbf{q}$ points of the reciprocal lattice, where \mathbf{G} is the zone center vector, and \mathbf{q} is the reduced vector. The exact \mathbf{Q} points are $\mathbf{Q} = (3.02, 0, 0.18)$, $(3.09, 0, 0.06)$, $(3.15, 0, -0.06)$, $(3.22, 0, -0.19)$ at SPring8 and $\mathbf{Q} = (3.25, 0, 0)$, $(3.45, 0, 0)$ at ESRF. We will neglect the c -axis component of the momentum transfer since the phonon anomaly was found to be independent of it [17]. The six \mathbf{Q} points will be described by the reduced vector $\mathbf{q} = (\xi, 0, 0)$.

The ARPES experiments were carried out at beam line V-4 of the Stanford Synchrotron Radiation Laboratory with 22.4 eV photon. The ARPES data along the node were reproduced using an in-house 6 eV laser-based setup similar to the one described in Ref. [18]. In both cases the total energy and angular resolution were better than 14 meV and 0.4° , respectively, and the sample temperature was below 15 K .

Figure 1 shows a typical IXS energy scan in longitudinal geometry at $\mathbf{Q} = (3.02, 0, 0)$. A harmonic oscillator convoluted with the instrumental resolution function was used for the fit. Three features can be clearly distinguished in 1(a): (i) the elastic peak at 0 meV is the static diffuse scattering coming from disorder, (ii) the first longitudinal acoustic phonon around 6 meV , and (iii) two low energy optical modes at 14 and 21 meV . Other phonon modes are not resolved due to the dominating contribution from the tails of the elastic and acoustic signals. Figure 1(b) shows the same energy scan at higher energy. Two peaks corresponding to the last two longitudinal optical (LO) phonon modes are clearly resolved in the data.

Figure 2 shows the evolution of these two LO phonons across the BZ. Figure 2(a) shows the IXS spectra for ξ spanning from the BZ center to the BZ face. At the BZ center, two peaks are resolved. The first peak is at 58 meV and the second is at 76 meV . Both are almost resolution limited. As we approach the BZ face, both peaks disperse toward 65 meV and as ξ reaches 0.22 – 0.25 the two modes are indistinguishable. The $\xi = 0.25$ spectrum shows actually a rather ill-defined broad peak, extending far in the low energy side. For $\xi = 0.45$ we can resolve again two sharp peaks.

In Figs. 2(b) and 2(c) we show the dispersion of the two branches across the BZ. There are two likely scenarios depending on the symmetry of the two branches. If they have the same symmetry, they anticross; otherwise they

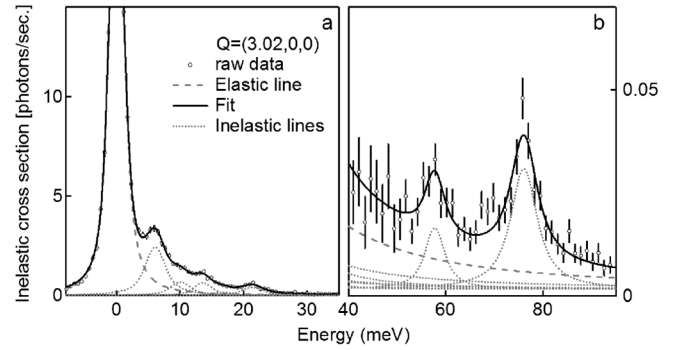


FIG. 1. IXS phonon spectrum of Bi2201 at 10 K . The markers are the raw data, the dashed line is the resolution function, and the plain line is a fit. (a) Low energy part of the spectrum. (b) High energy part with a smaller y scale.

simply cross. The dashed lines are guides for the eyes used to illustrate the crossing 2(b) and the anticrossing 2(c) scenarios. In the case of an anticrossing, the character of the two modes are exchanged between the two branches at the anticrossing point. Therefore in either case, the 60 meV peak at $\xi = 0.45$ has the character of the Cu-O BS mode, leading to a total softening of 16 meV . A similar anticrossing phenomenon is observed in the n -doped cuprate

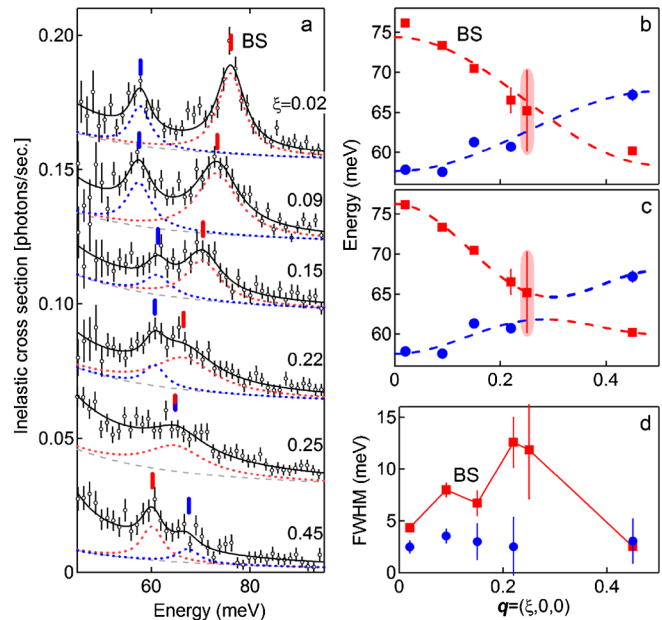


FIG. 2 (color online). LO phonons dispersion. (a) IXS spectra for $\mathbf{Q} = (3 + \xi, 0, 0)$ with ξ spanning from the BZ center (top spectrum, $\xi = 0.02$) to the BZ face (bottom spectrum, $\xi = 0.45$). The spectra are vertically shifted. The plain lines show the fit, the dashed lines show the elastic tail, and the dotted lines show the last two modes used in the fit. (b),(c) Phonon dispersions. The cosine dashed lines are guides for the eyes illustrating the crossing (b) and anticrossing (c) scenarios. (d) FWHM. The error bars are an estimate of the standard deviation of the fit coefficients.

$\text{Nd}_{1.86}\text{Ce}_{0.14}\text{CuO}_{4+\delta}$ [3,19] where the Cu-O BS mode anticrosses with the Nd-O BS mode. In the case of optimally doped $\text{La}_{2-x}\text{Sr}_x\text{Cu}_2\text{O}_4$ (LSCO), a mode with little dispersion is seen at the same energy, 58 meV, and has been attributed to the O bond bending mode [4].

In order to distinguish between the two scenarios, we undertook phonon calculations using a classical shell model. The parameters used are based on the model from Ref. [20]. This model did not reproduce the low and high energy modes observed experimentally in a reliable way. This indicates that spectral ellipsometry data alone could not provide enough constraints for the parameters of the model. All our attempts to find a better set of potentials compatible with previous established potential sets [21] failed. Several factors might be responsible for such disagreement, including (i) the charge distribution and bonds of the high Z atoms containing f electrons (Bi) are very difficult to model accurately [22] and (ii) the one-dimensional incommensurate superstructure along the b axis [23] which is neglected in the calculation.

Figure 2(d) shows the corresponding FWHM of each peak. The two modes are sharp and almost resolution limited across the whole BZ except for $\xi = 0.22\text{--}0.25$, where only one broad peak is resolved.

Figure 3 shows a comparison of the Cu-O BS phonon dispersion [3(a)] and FWHM [3(b)] for various optimally p -doped single layer cuprates. For Bi2201, an additional data point at $\xi \approx 0.3$ is still needed to establish the exact momentum value of the maximal softening, though these data already suggest that the maximum FWHM is observed around $\xi \approx 0.22\text{--}0.25$ in agreement with other cuprates [19]. We note however that the softening in the Bi2201 dispersion is smoother and observable already at very low ξ in contrast with other cuprates.

In Fig. 4 we compare the IXS and ARPES results for the phonon and electron dispersions in Bi2201. Figure 4(a) shows the dispersion relation near the Fermi energy measured by ARPES along the nodal direction from the node to the BZ face. In agreement with previous reports we observe a kink at 63 ± 5 meV along the nodal direction (cut 1) which shifts toward lower energy, 35 meV, as the tip of the Fermi arc, measured from the data in the pseudogap phase [24], is approached (cut 3). A similar shift of the kink toward lower binding energy has been reported for other cuprates [10–12]. The kink positions is extracted using a two straight lines fit [9], the first derivative of the dispersion [4(b)] and a self-energy type of analysis [25].

The direct comparison with the dispersion of the BS phonon clearly shows that the energy of the 63 meV kink coincides with the energy of the soft part of the BS phonon. To extend the comparison also to the momentum space we show in 4(c) a constant energy map of ARPES spectral intensity versus momentum at the Fermi energy (Fermi surface) in the superconducting state and a constant energy contour (gray line) at the nodal kink energy (63 meV).

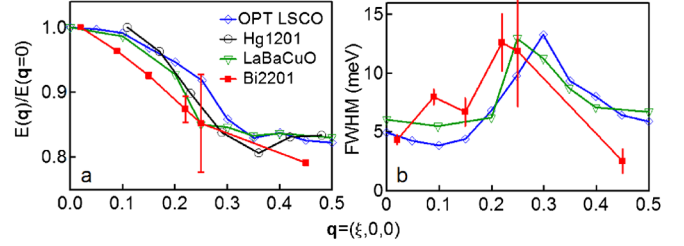


FIG. 3 (color online). (a) Dispersion of the BS mode for various optimally p -doped single layer cuprates in comparison with the Bi2201 data. The frequencies are normalized to their value at $\xi = 0$. The error bars are an estimate of the standard deviation of the fit coefficients. The optimally doped LSCO and LaBaCuO data are from [17] and the Hg1201 data are from [1]. (b) FWHM when available.

The comparison between the momentum of the BS phonon and the constant energy maps can provide unique information on the wave vectors needed for the electrons at the nodal kink to be scattered to the Fermi level by this particular phonon. The gray shadow area represents the range of $\mathbf{q} = (\xi, 0, 0)$ vectors that connect the electronic excitations at the 63 meV kink energy (gray closed contour) to the Fermi surface segments from the node to the BZ face at the locations where the nodal kink is observed.

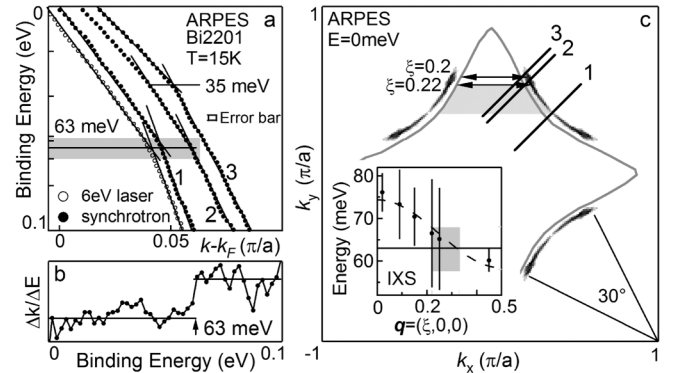


FIG. 4. (a) Electron dispersion measured by ARPES for three different momentum cuts. Cut 1 is along the node and was measured with 22.4 eV synchrotron radiation (solid circles) and with 6 eV radiation using a laser setup (open circles). Cuts 2 and 3 are further toward the BZ face, close to the edge of the Fermi arcs reported in pseudogap state [24]. Cuts 2 and 3 were measured with synchrotron light only. The spectra are shifted horizontally and a ≈ 4 meV gap for cuts 2 and 3 is subtracted. (b) First derivative of cut 1. (c) Experimental Fermi surface with the momentum location of the three cuts. The apparent finite Fermi arcs are due to the experimental resolution. The solid line shows a constant energy contour at the kink energy, 63 meV. The inset of (c) shows the IXS dispersion and peak FWHM (error bars) of the BS mode discussed above. The shadow area represents the nodal kink energy and the \mathbf{Q} vectors connecting the Fermi surface segments where the nodal kink is observed. Note that $\xi = 0.2$ corresponds to $0.4\Delta k_x$ (π/a).

For clarity, the gray area is shown only for half of a Fermi arc, but the same construction can be applied to the entire Fermi arc by symmetry.

We find that the entire momentum region where we observe a 63 meV kink can be connected with a momentum transfer $\mathbf{q} = (\xi, 0, 0)$ or equivalently $\mathbf{q} = (0, \xi, 0)$ with $\xi \geq 0.22$ to an opposite, quasiparallel Fermi surface segment. This momentum transfer corresponds exactly to the momentum region where the Cu-O BS phonon shows a softening and an anomalously *broad* line shape [illustrated with the large error bars in the inset of 4(c)]. Note that the furthest tips of the Fermi surface arcs are connected by $\mathbf{q} = (0.5, 0, 0)$ which still corresponds to the soft BS mode. However, because of the relative orientation of these Fermi arcs (the momentum transfer is almost parallel to the Fermi arcs), the coupling with a perpendicular phonon [$\mathbf{q} = (0, 0.22, 0)$] is expected to be stronger. Interestingly, the maximum broadening corresponds to the case in which the momentum transfer is almost perpendicular to the Fermi arcs, suggesting that the anomalous softening and broadening observed in the phonon spectra might be a consequence of this Fermi surface topology. Note that although the weak doping dependence of the phonon softening seems to be in contrast with the strong doping dependence of the Fermi arcs, the broadening of the phonon peak at $q = 0.25$ is strongly doping dependent in line with the proposed scenario. Similar IXS and ARPES studies are needed to shed light on the nature of the 35 meV mode.

Finally we note that the BS mode is supposed to be nondispersive at about 80 meV along the other directions and, in particular, along [110] for the full breathing mode [17,19]. The absence of any strong feature above 63 meV in the ARPES data shows that the nodal charge carriers are preferentially coupled with the soft Cu-O half-breathing BS mode dispersing along the [100] direction as suggested by recent local spin-density approximation + U results [26].

In conclusion, we report the first evidence of an anomalous dispersion of the Cu-O bond stretching phonon mode in a Bi cuprate. A direct comparison with ARPES data shows that the energy and momentum where the strongest coupling to a 63 meV mode is observed corresponds exactly to the energy and momentum of the soft and broad Cu-O BS phonon. Considering that a magnetic resonance mode was never observed in single layer Bi2201 [27], this result provides a strong support for the lattice origin of the 60–70 meV ARPES nodal kink. Although the role of

electron-phonon coupling for superconductivity is not known yet, it is beyond any doubt that it is an important interaction linked with the Fermi surface topology, and it should not be neglected in any realistic theory of cuprates.

We thank D. Reznik and D.-H. Lee for useful discussions. The ARPES and the IXS measurements were supported by the Division of Materials Sciences and Engineering, Office of Basic Energy Sciences of the U.S. Department of Energy under Contract No. DE-AC03-76SF00098. The IXS measurements were also partially supported by the France-Berkeley Grant.

*ALanzara@lbl.gov

†Institut de Minéralogie et de Physique des Milieux Condensés (IMPMC), CNRS UMR 7590, Campus Boucicaut, 140 rue de Lourmel, 75015 Paris, France.

- [1] H. Uchiyama *et al.*, Phys. Rev. Lett. **92**, 197005 (2004).
- [2] N. Pyka *et al.*, Phys. Rev. Lett. **70**, 1457 (1993).
- [3] M. d'Astuto *et al.*, Phys. Rev. Lett. **88**, 167002 (2002).
- [4] R.J. McQueeney *et al.*, Phys. Rev. Lett. **82**, 628 (1999).
- [5] J.-H. Chung *et al.*, Phys. Rev. B **67**, 014517 (2003).
- [6] L. Pintschovius and M. Braden, Phys. Rev. B **60**, R15 039 (1999).
- [7] D. Reznik *et al.*, Nature (London) **440**, 1170 (2006).
- [8] W. Reichardt and M. Braden, Physica (Amsterdam) **263–264B**, 416 (1999).
- [9] A. Lanzara *et al.*, Nature (London) **412**, 510 (2001).
- [10] T. Cuk *et al.*, Phys. Rev. Lett. **93**, 117003 (2004).
- [11] A.D. Gromko *et al.*, Phys. Rev. B **68**, 174520 (2003).
- [12] K. Terashima *et al.*, Phys. Rev. Lett. **99**, 017003 (2007).
- [13] P.D. Johnson *et al.*, Phys. Rev. Lett. **87**, 177007 (2001).
- [14] A. Kaminski *et al.*, Phys. Rev. Lett. **86**, 1070 (2001).
- [15] M. d'Astuto *et al.*, Physica (Amsterdam) **316–317B**, 150 (2002).
- [16] H. Eisaki *et al.*, Phys. Rev. B **69**, 064512 (2004).
- [17] D. Reznik *et al.*, J. Low Temp. Phys. **147**, 353 (2007).
- [18] J.D. Koralek *et al.*, Phys. Rev. Lett. **96**, 017005 (2006).
- [19] M. Braden *et al.*, Phys. Rev. B **72**, 184517 (2005).
- [20] N.N. Kovaleva *et al.*, Phys. Rev. B **69**, 054511 (2004).
- [21] S.L. Chaplot *et al.*, Phys. Rev. B **52**, 7230 (1995).
- [22] B. Renker, H. Schober, and F. Gompf, J. Low Temp. Phys. **105**, 843 (1996).
- [23] M.D. Kirk *et al.*, Science **242**, 1673 (1988).
- [24] T. Kondo *et al.*, Phys. Rev. Lett. **98**, 267004 (2007).
- [25] S. Verga, A. Knigavko, and F. Marsiglio, Phys. Rev. B **67**, 054503 (2003).
- [26] P. Zhang, S.G. Louie, and M.L. Cohen, Phys. Rev. Lett. **98**, 067005 (2007).
- [27] T. Sato *et al.*, Phys. Rev. Lett. **91**, 157003 (2003).