

Luminescence from cavitation bubbles deformed in uniform pressure gradientsOtti Supponen,^{1,*} Danail Obreschkow,² Philippe Kobel,¹ and Mohamed Farhat¹¹Laboratory for Hydraulic Machines, Ecole Polytechnique Fédérale de Lausanne, Avenue de Cour 33 Bis, 1007 Lausanne, Switzerland²International Centre for Radio Astronomy Research, University of Western Australia, M468 7 Fairway, Crawley, Western Australia 6009, Australia

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Presented here are observations that demonstrate how the deformation of millimetric cavitation bubbles by a uniform pressure gradient quenches single-collapse luminescence. Our innovative measurement system captures a broad luminescence spectrum (wavelength range, 300–900 nm) from the individual collapses of laser-induced bubbles in water. By varying the bubble size, driving pressure, and perceived gravity level aboard parabolic flights, we probed the limit from aspherical to highly spherical bubble collapses. Luminescence was detected for bubbles of maximum radii within the previously uncovered range, $R_0 = 1.5\text{--}6$ mm, for laser-induced bubbles. The relative luminescence energy was found to rapidly decrease as a function of the bubble asymmetry quantified by the anisotropy parameter ζ , which is the dimensionless equivalent of the Kelvin impulse. As established previously, ζ also dictates the characteristic parameters of bubble-driven microjets. The threshold of ζ beyond which no luminescence is observed in our experiment closely coincides with the threshold where the microjets visibly pierce the bubble and drive a vapor jet during the rebound. The individual fitted blackbody temperatures range between $T_{\text{lum}} = 7000$ and $T_{\text{lum}} = 11\,500$ K but do not show any clear trend as a function of ζ . Time-resolved measurements using a high-speed photodetector disclose multiple luminescence events at each bubble collapse. The averaged full width at half-maximum of the pulse is found to scale with R_0 and to range between 10 and 20 ns.

DOI: [10.1103/PhysRevE.96.033114](https://doi.org/10.1103/PhysRevE.96.033114)**I. INTRODUCTION**

As a cavitation bubble undergoes a spherical collapse, it compresses its enclosed gaseous contents and—presumably—adiabatically heats them to temperatures of several thousands of degrees, which results in light emission called luminescence [1]. The drive to investigate luminescence comes from the intense energy focusing at a bubble collapse, which provides a catalytic host for unique chemical reactions [2,3], offering potential for cancer therapy [4,5], environmental remediation [6,7], and fabrication of nanomaterials [8,9]. While most past studies have researched sonoluminescence, that is, luminescence from acoustically driven bubbles, light emission has also been detected from hydrodynamic cavitation in engineering flows [10,11].

Due to the occurrence at the last instant of the collapse, the redistribution of the bubble's energy into luminescence, as well as shock waves, microjets, and elastic rebounds (see introduction in Ref. [12]), must be highly sensitive to topological changes in the cavity volume during the final collapse stage. This represents an important feature, considering that any anisotropy in the pressure field of the surrounding liquid will result in the deformation of an initially spherical bubble, inducing a microjet that pierces the bubble and therefore causes it to undergo a toroidal collapse [13,14]. The level of compression of the bubble gases is reduced for even slight bubble deformations, manifested in the weakening of the collapse shock-wave emissions [15,16]. Indeed, luminescence has been shown to vary with the proximity of near surfaces that break the spherical symmetry of the bubble [17–19]. It has also been shown that the lack of buoyancy enhances the energy

concentration at the final stage of bubble collapse [20], even for bubbles that are highly spherical and generally assumed not to be subject to deformation by gravity (maximum bubble radius, $R_0 \sim 40$ μm at atmospheric pressure). Bubbles collapsing with pronounced microjets in multibubble fields have been shown to emit less light (or none) compared to spherically collapsing bubbles [21].

Spectral analyses of luminescence have proposed a wide range of temperatures at a bubble collapse in water, depending on whether the bubble is trapped in an oscillating acoustic field (bubble temperatures $T > 10^4$ K) [1], induced by a laser pulse ($T \sim 7000\text{--}8000$ K) [22,23], induced by a spark ($T \sim 6700$ K) [24], or within a bubble cloud ($T < 5000$ K) [2,25]. Recent studies reached 1.4×10^4 K for an energetic bubble collapse provoked by piezoelectric transducers [26] and over 2×10^4 K for a centimetric bubble expanded by a chemical reaction in a free-fall tower [27]. Moreover, luminescence spectra from small bubbles (maximum radius $R_0 < 1$ mm) show a smooth continuum similar to a blackbody, while spectra of luminescence from large, laser-induced bubbles ($R_0 > 1$ mm) and multibubble sonoluminescence have shown emission lines of excited hydroxyl (OH^-) bands at 310 nm [19,28] that have been associated with aspherical bubble collapses. It is unclear, however, to what extent the spectral differences in these distinct scenarios are caused by physical or experimental factors, and a systematic picture of the role of pressure-field anisotropies—and the resulting bubble deformation—on luminescence is still lacking.

This work presents observations on the luminescence of initially highly spherical, millimetric bubbles collapsing at different levels of deformation caused by the gravity-induced uniform pressure gradient. We probe the transition from toroidal jetting bubbles in controlled pressure gradients to highly spherical bubbles in microgravity and cover a broad

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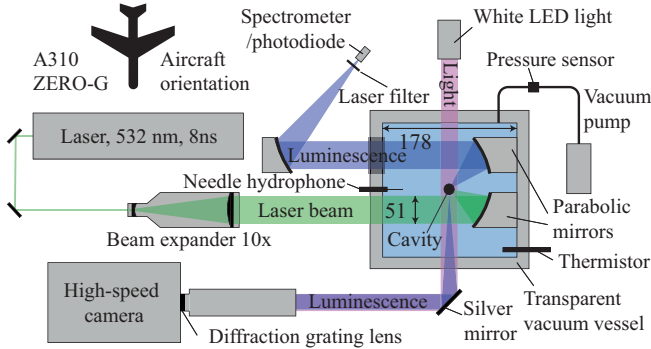


FIG. 1. Schematic top view of the experimental setup. Dimensions are given in millimeters.

parameter space. Spectral and time-resolved measurements are made on single-cavitation-bubble luminescence from individual collapses of transient, laser-induced vapor bubbles in water, contrasting with the established single-bubble sonoluminescence, which is normally understood as the time-averaged light emitted by an oscillating bubble trapped in an acoustic field. It also differs from the averaged single-cavitation-bubble luminescence, from luminescence of gas bubbles, and from luminescence of bubbles in liquids doped with noble gas.

II. EXPERIMENTAL SETUP

Figure 1 shows a schematic of our experiment. We generate highly spherical bubbles by using an immersed parabolic mirror to focus a Q-switched Nd:YAG laser (532 nm, 8 ns) in the middle of a cubic test chamber filled with demineralized water. The water is initially partially degassed to remove large bubbles from the container boundaries, but we presume the water to be mostly air saturated. The bubbles are so spherical that the dominant pressure-field anisotropy deforming the bubble is the gravity-induced pressure gradient [12]. Furthermore, owing to their high sphericity, these bubbles do not suffer a fission instability, i.e., bubble splitting [23,29], during their collapse, allowing very large bubbles to compress their enclosed gases efficiently and luminesce in the absence of external perturbations. We obtain the bubble's maximum radius R_0 by measuring its collapse time T_c (i.e., half oscillation time) of the bubble with a needle hydrophone, which detects the passage of the shock waves emitted at the generation and the collapse of the bubble. The maximum bubble radius is then obtained via $R_0 = 1.093T_c(\Delta p/\rho)^{1/2}$ [30], where $\Delta p = p_0 - p_v$ is the driving pressure (p_0 being the static pressure at the height of the bubble and p_v the liquid vapor pressure) and ρ is the liquid density. It is considered unnecessary to correct this relation for the bubble's asphericity, as the deformations in this work remain weak. The temperature of the water is recorded with a thermistor and kept at room temperature (294.2 ± 1 K), and p_v is computed for each bubble individually using the Antoine equation. Simultaneous visualizations of the luminescence, radial evolution of the bubble, and shock-wave emission are made with an ultrahigh-speed CMOS camera (Shimadzu HPV-X2) filming at 10×10^6 frames/s (fps) with an exposure time of 50 ± 10 ns and a back-light LED.

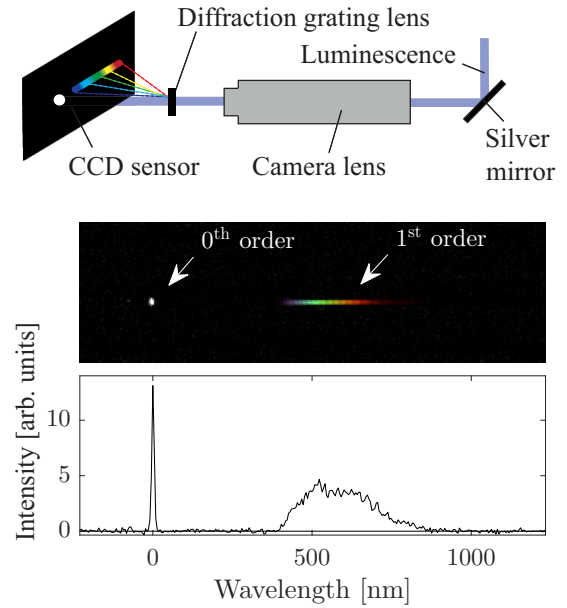


FIG. 2. Top: Schematic of the CCD luminescence detection system coupled with a diffraction grating lens. Middle: Typical luminescence signal (zeroth- and first-order spectra) as recorded by the CCD sensor. Bottom: Corresponding raw spectrum obtained from the pixel intensities of the image.

The time-averaged luminescence spectrum from a single bubble collapse is captured in the dark by a spectrometer (Ocean Optics QEPro; exposure time, 8 ms). The light emitted during the bubble collapse is collected using a second, aluminum-coated, immersed parabolic mirror that reflects it through a fused silica window (for UV transparency) onto another parabolic mirror outside the test chamber. We chose aluminum-coated mirrors for their good UV reflection quality. The external mirror focuses the light through a laser-blocking filter onto the entrance of the optical fiber that leads to the spectrometer. Without the filter the laser would saturate the measured spectrum despite the spectrometer's being triggered only after the bubble generation.

The luminescence spectrum is simultaneously measured with a second high-speed camera (Photron SA1.1) that has a CCD sensor (in place of the CMOS camera). It is equipped with an astronomy-quality diffraction grating lens (RSpec Star Analyzer SA-100) and films at 10^5 fps with an exposure time of $10 \mu\text{s}$. The reason for using the CCD instead of the CMOS camera to measure the spectrum is that it guarantees the luminescence to be fully contained in its exposure time, which the latter cannot. The grating lens, placed between the camera objective and the CCD sensor, splits and deviates the light one or more diffraction orders located in a plane perpendicular to the grating lines, thus providing a spectrum on the sensor. A schematic of the CCD light detection system is shown in Fig. 2 along with a typical measured luminescence signal. The reasons behind measuring the spectrum additionally with the camera are that it fills in the spectral gap in the spectrometer ($\sim 500\text{--}700$ nm) caused by the laser-blocking filter and, more importantly, corrects the intensity of the spectrum recorded by the spectrometer, which is affected by the bubble's migration away from the parabolic mirror's

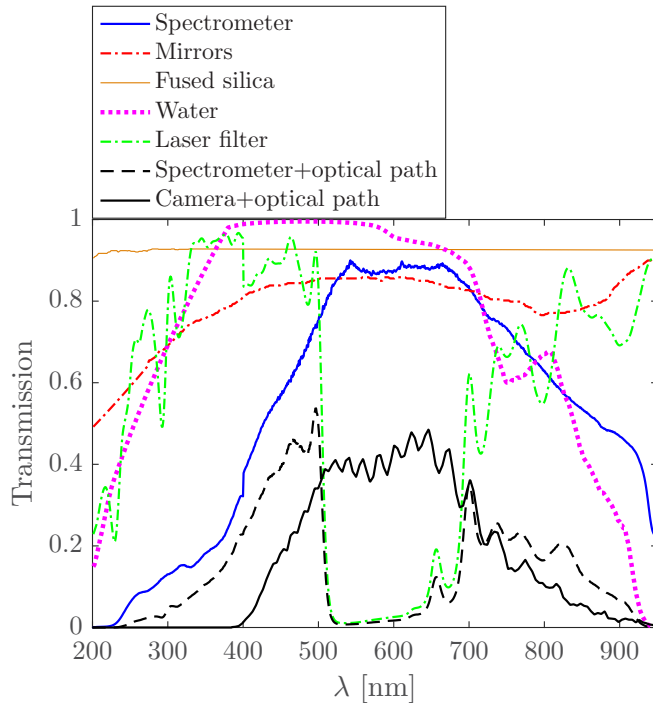


FIG. 3. Transmission of light as a function of the wavelength λ for the various elements on the optical path from the luminescence emission point to the detectors.

focal point. The bubble's displacement becomes important, in particular, at higher gravity levels for large bubbles that experience a strong Kelvin impulse [14] (i.e., the integrated momentum of the liquid during the growth and the collapse of the bubble [31]). This displacement can weaken the signal measured with the spectrometer, and therefore it is corrected using the spectrum recorded by the CCD. The CCD spectrum measurement is unaffected by the bubble's displacement, as the luminescence spot stays within the image plane.

The optical path from the luminescence to the spectrometer includes 194 mm water, 6 mm fused silica, two aluminum-coated parabolic mirrors and the laser filter. To reach the camera's CCD sensor, the luminescent light travels through water, acrylic glass, a silver mirror, the camera lens and the grating lens. The wavelength-dependent transmissions of the various elements in the optical paths are shown in Fig. 3. The calibration of the spectrometer detector and the absorption and transmission spectra of the various optical components were provided by their respective manufacturers. Water's absorption spectrum in the wavelength range of interest is found in the literature [32]. The spectrum measured by the high-speed camera with the grating filter was calibrated in-house combining the transfer functions of the camera and the optical path using a thermal light source placed inside the test chamber at the location where the bubble was generated. This innovative luminescence measurement system allows for (i) the collection of a substantial amount of light from the rapid, small, and weak luminescence of a single bubble collapse, (ii) the capture of a wide spectrum from a single bubble collapse,

including the UV, and (iii) staying far from the bubble so as not to disturb its dynamics.

Time-resolved measurements of the luminescence pulse are made using the same optical path as described above for the spectrometer, but by focusing the light onto a high-speed photodetector (Thorlabs DET10A/M Si detector) without a laser-blocking filter. The detector has a 1-ns rise time and is sensitive in the 200- to 1100-nm wavelength range. The photodetector signal is recorded by an oscilloscope (4-GHz bandwidth), which is triggered using the hydrophone signal of the collapse shock wave and applying a 25- μ s negative delay to account for the shock-wave propagation over a distance of ~ 37 mm to reach the hydrophone after the bubble collapse.

Three parameters influencing the bubble luminescence can be independently varied in our experiment: (i) the driving pressure, $\Delta p \equiv p_0 - p_v$ (0.06–1 b), where p_0 is adjusted using a vacuum pump; (ii) the bubble energy, $E_0 = (4\pi/3)R_0^3\Delta p$ (0.4–28 mJ), adjusted by the laser pulse energy; and (iii) the constant, uniform pressure gradient, $\nabla p (= \rho \mathbf{g}$, with the perceived gravitational acceleration $|\mathbf{g}|$ varied between 0 and 2 g , where $g = 9.81 \text{ ms}^{-2}$), modulated aboard European Space Agency parabolic flights (the 58th, 60th, and 62nd parabolic flight campaigns) and on the first Swiss parabolic flight. The interest in using the hydrostatic pressure gradient to deform bubbles is based on its uniformity in space and time, in contrast to near boundaries. This is an advantage in particular as it probes the influence of pressure gradients induced by any other inertial forces in addition to gravity. Moreover, any practical instance of a smooth pressure field can be approximated to first order by such a uniform pressure gradient, thus extending the scope of this study to any situation involving bubbles in anisotropic pressure fields [14,33]. These variables yield a wide range of maximum bubble radii, $R_0 \sim 1.5$ –10 mm. Such large bubbles present the advantage of easier resolution of the time and space scales associated with their collapse, in contrast to, e.g., single-bubble sonoluminescence experiments. Additional details on the experiment and the parabolic flights may be found in Ref. [12].

We account for the effect of bubble asphericity due to the gravity-induced pressure gradient through the anisotropy parameter $\zeta \equiv |\nabla p|R_0\Delta p^{-1}$, which is the dimensionless equivalent of the Kelvin impulse [14,31,33]. Here ζ is varied by adjusting the maximum bubble radius R_0 , the driving pressure Δp , and the pressure gradient $|\nabla p|$ (through variable gravity). Measuring at variable gravity allows for the decoupling of the roles of the driving pressure (Δp) and bubble deformation (ζ), which is important because the expression of ζ for gravity-induced deformation includes Δp . The pressure-field anisotropy caused by the nearest boundary in our experiment is considered with $\zeta = -0.195\gamma^{-2}$ (which represents the dimensionless Kelvin impulse for bubbles near boundaries [14]), where γ is the standoff parameter $\gamma = s/R_0$ and where $s = 55$ mm is the distance between the bubble center and the parabolic mirror. The resultant ζ is given by the vector sum of the respective directional ζ . We expect luminescence to vary with ζ , since an increasing ζ implies stronger bubble deformation, which, in turn, affects the different events associated with the bubble collapse, such as microjets [14,33] and shock waves [15,16].

III. SPECTRAL ANALYSIS IN VARIABLE GRAVITY

Selected images of high-speed movies visualizing luminescing bubbles of the same energy E_0 collapsing at different levels of ζ at normal gravity are shown in Fig. 4. The bubble interface, the luminescence, and the sharp shock waves are captured in the same movie, owing to the short exposure time (50 ns). We observe a weakening of the luminescent flash with increasing ζ . One may also see a pronounced deflection of light near the bubble wall in the frames preceding the luminescence, which is due to the pressure rise in the surrounding liquid predicted by Lord Rayleigh a century ago [30]. At $\zeta = 3.8 \times 10^{-3}$ there is no visible luminescence and the bubble's deformation is clearly manifested by the emitted shock wave(s) no longer being spherically symmetric.

The luminescence spectrum is well approximated by the blackbody model [23,34], and since the bubble temperature cannot be directly measured, a fitted blackbody provides a reasonable estimation for it. The effective blackbody temperature and energy of luminescence can be inferred by fitting the spectra with a Planckian function of the form

$$L(\lambda, I, T_{\text{lum}}) = A \frac{I}{\lambda^5} \frac{1}{\exp\left(\frac{hc}{\lambda k_B T_{\text{lum}}}\right) - 1} \quad (\text{J/nm}), \quad (1)$$

where λ is the wavelength, h and k_B are the Planck and Boltzmann constants, respectively, c is the speed of light, A is a constant prefactor determined from calibration, T_{lum} is the blackbody temperature, and I stands for the product of the luminescence pulse duration and the projected emitting surface (which cannot be disentangled with the spatial and temporal

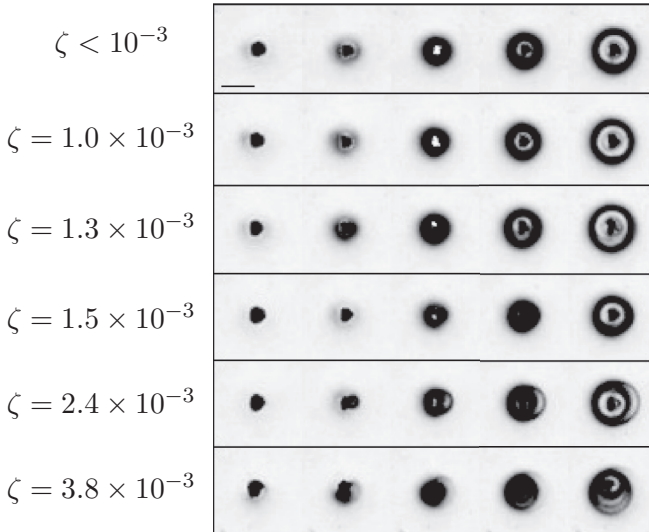


FIG. 4. Visualization of luminescence emitted at the final collapse stage of a single cavitation bubble at various ζ . The luminescent flash is visible in the middle frame and followed by the rebound. The interframe time is 100 ns, the exposure time is 50 ns, and the black line in the top panel shows the 1-mm scale. The bubble energy is the same in all cases ($E_0 \approx 27$ mJ) and ζ is varied by adjusting the driving pressure, from top to bottom, as $\Delta p = 98, 78, 58, 48, 28$, and 18 kPa, yielding maximum bubble radii of $R_0 = 4.1, 4.3, 4.8, 5.1, 6.1$, and 7.1 mm. These bubbles were imaged on-ground at normal gravity.

resolution of our apparatus). The best-fit values are obtained by fitting Eq. (1), after correcting it for the absorption losses in Fig. 3, with the measured raw spectra through maximum likelihood for the pair $(E_{\text{lum}}, T_{\text{lum}})$, where $E_{\text{lum}} = IT_{\text{lum}}^4$ is the luminescence energy through the Stefan-Boltzmann law. The estimated standard error of the maximum likelihood fit is obtained from the covariance matrix (estimated via the inverse of the Hessian matrix) representing the goodness of fit to the data. Figure 5 displays a typical measured luminescence spectrum from a single bubble collapse.

We estimate the total luminescence energy E_{lum} by assuming a uniform light emission in the solid angle of 4π . In this way, 6.7% of all the photons are expected to reach the calibrated spectrometer detector. We use as a reference a highly spherical bubble collapsing in microgravity, which is assumed to undergo no displacement from the focal point of the parabolic mirror.

Figure 6(a) shows the luminescence energy E_{lum} , obtained through the best Planckian fit, as a function of the maximum bubble radius R_0 for three different ranges of driving pressure Δp . Only bubbles collapsing highly spherically ($\zeta < 7 \times 10^{-4}$) have been selected in order to exclude deformation-induced hindering of the luminescence, and the data include bubbles collapsing in microgravity. The maximum radii are within the range $R_0 = 1.5$ –3.5 mm, which, to our knowledge, extend to the largest reported laser-induced luminescing bubbles collapsing freely and spherically in water. As expected, one may observe an increase in E_{lum} with increasing R_0 for a fixed Δp , the tendency being consistent

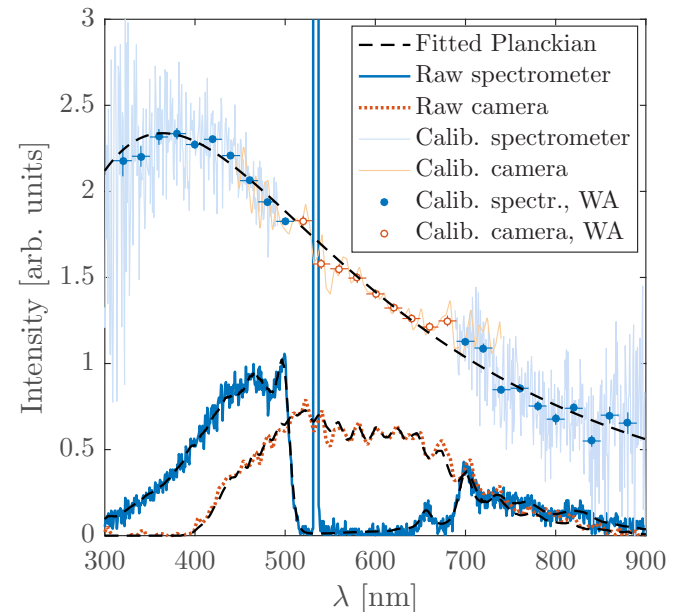


FIG. 5. Typical luminescence spectrum from a single bubble collapse, measured with a spectrometer with an exposure time of 8 ms and a high-speed CCD camera with an exposure time of 10 μs . Both raw and calibrated spectra are shown, together with the fitted Planckians. Window averages (WAs) of 20-nm windows are also displayed. The peak around 532 nm is caused by the strong laser pulse despite the $>99\%$ attenuation of the filter. Here $R_0 = 3.0$ mm, $\Delta p = 78$ kPa, and $|g| = 1$ g.

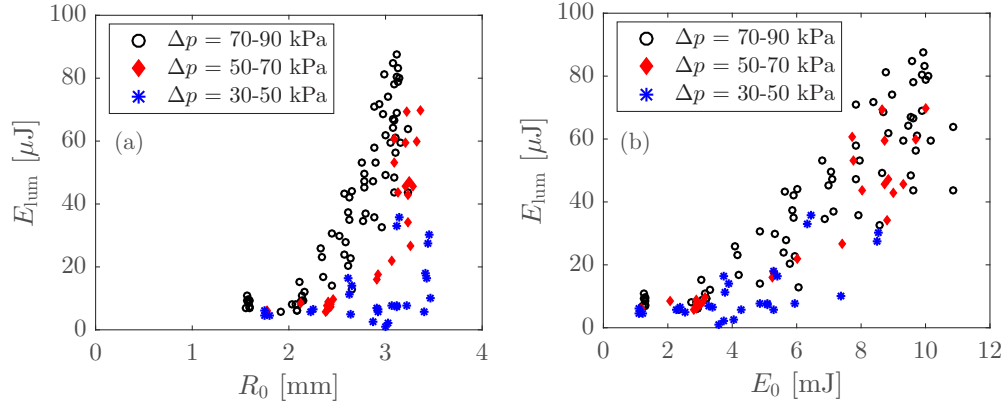


FIG. 6. Luminescence energy E_{lum} as a function of (a) the maximum bubble radius R_0 and (b) the bubble energy E_0 for three ranges of driving pressures Δp . Each point corresponds to a measurement from a single, spherical collapse ($\zeta < 7 \times 10^{-4}$).

with the literature [17,18,29]. In the literature, however, a decrease in luminescence energy for laser-induced bubbles with increasing maximum radii beyond $R_0 \approx 1.5$ mm has also been reported [18]. This is likely attributed to the use of less pointlike focusing methods (e.g., converging lens) that yield bubbles that are more disturbed in the collapse phase and cause, e.g., bubble splitting [22,29], such disturbances being enhanced for increasing bubble radius. Bubbles with $R_0 > 3.5$ mm in our experiment are affected by the nearest surface, i.e., the parabolic mirror at a distance of 55 mm from the bubble center, which is accounted for in ζ .

For a given R_0 , a lower Δp yields weaker luminescence, which is expected since $E_{\text{lum}} \propto E_0 = (4\pi/3)R_0^3\Delta p$ [22,35]. Figure 6(b) verifies this relation but still suggests slightly weaker luminescence energies for bubbles collapsing with a lower Δp . This result is consistent with the past observation of more energetic luminescence from bubbles collapsing at higher static pressures for a fixed E_0 [35]. Bubbles at a low Δp have a longer collapse time and thereby an increased surface area and interaction time, possibly yielding increased energy loss by thermal conduction or mass flow by nonequilibrium evaporation or condensation at the bubble wall [35].

The important scatter of our results is due to the limited reproducibility of the luminescence. We find the spectral intensities between individual bubbles under the same conditions to vary by approximately 45%, while the maximum bubble radii vary by less than 1%. These brightness fluctuations are likely related to the microscopic size of the luminescent plasma, which makes it highly sensitive to minor perturbations and easily obscured by nuclei and impurities in the water.

Figure 7 displays three examples of typical spectra of single-cavitation-bubble luminescence, with the only varying parameter being the perceived gravity level (0g, 1g, and 1.8g). It is evident that the gravity-induced pressure gradient quenches the single-cavitation-bubble luminescence energy. Surprisingly, on none of the raw spectra do we observe a prominent peak corresponding to OH^- or other emission lines at any wavelength, even for the most deformed luminescing bubbles. This could, however, be due to the limited wavelength resolution of our apparatus.

To quantify the fraction of the bubble energy dissipated into luminescence, we normalize the luminescence energy

E_{lum} to the bubble energy E_0 . We only retain cases where luminescence is detected by both the spectrometer and the CCD camera. Note that the CCD signal helps correcting the spectrum of the spectrometer if the bubble moves out of the focal point of the parabolic mirror during its collapse. The dependence of the relative luminescence energy on the anisotropy parameter ζ is displayed in Fig. 8(a). Here ζ is altered by a wide range of R_0 , g , and Δp in order to disentangle their respective effects on luminescence from that of the bubble deformation. The maximum Δp was achieved when the test vessel reached the aircraft cabin pressure, i.e., $p_0 \approx 80$ kPa. The results show a rapid quenching of the relative luminescence energy with increasing ζ . Luminescence takes up to approximately 1% of the bubble's initial energy. The rest of the bubble's energy is distributed predominantly into shock-wave emission and the formation of a rebound bubble for spherically collapsing bubbles [36]. Owing to microgravity, we are able to create large bubbles, which in normal gravity would be deformed, that collapse highly spherically at low Δp

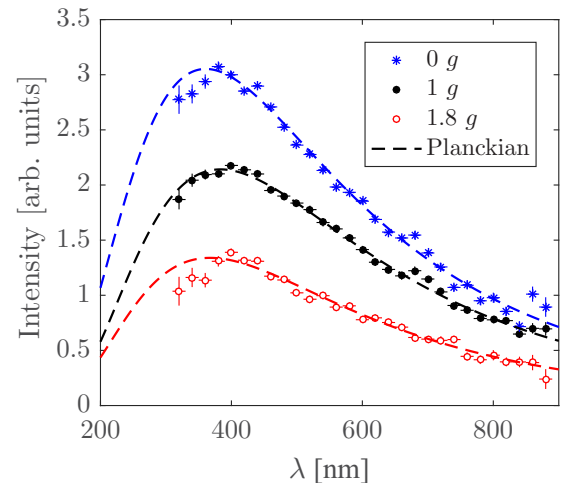


FIG. 7. Single-cavitation-bubble luminescence spectra at three gravity levels for the same laser pulse energy ($R_0 = 3 \pm 0.1$ mm) and static pressure of the water ($p_0 = 81 \pm 1$ kPa). Each spectrum is measured at a single bubble collapse.

and emit luminescence. Correspondingly, higher gravity levels allow us to stretch the range of ζ to higher values for a given Δp . Up to the scatter, the data points exhibit a linear trend on a logarithmic scale as a function of ζ regardless of the gravity level. Luminescence is not detected by the spectrometer for anisotropy levels beyond $\zeta \approx 3.5 \times 10^{-3}$, which corresponds to the same Kelvin impulse at $\gamma \approx 7.5$ for bubbles deformed by neighboring surfaces [14]. Note that we only obtain reliable fitted blackbody energies, which require the CCD signal, up to $\zeta \approx 1.8 \times 10^{-3}$ [in Fig. 8(a)], due to the poor signal-to-noise ratio of luminescence from more deformed bubbles.

Figure 8(b) displays our best-fit estimates of the bubble's blackbody temperatures as a function of ζ . We obtain reliable fitted blackbody temperatures, which only require the spectrometer signal, up to $\zeta \approx 2.5 \times 10^{-3}$. The temperatures fall in the range $T_{\text{lum}} = 7000\text{--}11\,500$ K, which is in good agreement with previous laser-induced bubble luminescence studies, in which the temperatures from averaged spectra varied between 7680 K (close to a solid surface) and 9150 K (at elevated ambient pressure) [19,22]. This range, however, is attributed to the important scatter (which is expected owing to the experimental and fitting errors) rather than a clear relationship with the governing parameters. The highly spherical bubbles with the highest luminescence energies do not exhibit higher blackbody temperatures than the luminescing deformed bubbles. This result is in disagreement with the observations of Brujan and Williams [22], who found the temperatures

(estimated from averaged spectra) to decrease with decreasing distance between the bubble and a rigid boundary, that is, with increasing bubble deformation.

IV. TIME-RESOLVED MEASUREMENTS

The luminescence pulse duration for spherically collapsing laser-induced cavitation bubbles has been shown to be on the nanosecond scale and to scale with the maximum bubble radius R_0 . For example, for $R_0 = 0.3$ mm, the full width at half-maximum (FWHM) has been measured as $\tau \approx 3$ ns [22,29]; for $R_0 = 1$ mm, $\tau \approx 6\text{--}8$ ns [18,22]; and for $R_0 = 1.8$ mm, $\tau \approx 10$ ns [18]. Centimetric bubbles generated by a spark or expanded through a chemical reaction may luminesce for tens of microseconds [24,27]. Owing to the high sphericity of the initial plasma generating the bubble, large bubbles in our experiment ($R_0 > 2$ mm) are able to collapse spherically without bubble splitting decreasing the efficiency of the final gas compression. We therefore expect the luminescence pulse durations here to exceed those reported in the literature for laser-induced bubbles.

Figures 9(a)–9(f) show waterfall plots of 20 photodetector signals measured from single bubble collapses with a fixed bubble energy $E_0 \approx 22$ mJ and at different driving pressures Δp . The signals are sorted so that the peak amplitudes are in descending order from bottom to top. Here $t = 0$ μs corresponds to the instant at which the hydrophone detects the collapse shock, which has propagated a distance of 37 mm from the bubble. It should be noted that the amplitudes of the photodetector signals are not corrected for the bubble displacement. All photodetector measurements are made on-ground at normal gravity. The standard deviation of the maximum peak timing with respect to $t = 0$ μs ranges from 8 to 12 ns. Consistent with the spectral analysis in Sec. III, the energy of the luminescence signals decreases with increasing ζ . The number of peaks in the photodetector signals varies between one and four, suggesting multiple events yielding light emission. Similar peaks have been observed in the past in photomultiplier tube measurements for both single and multiple bubble collapses [18,37]. Such multiple peaks are often randomly distributed in time with respect to the strongest peak, which, for the majority of cases, is the last event. Figures 9(a)–9(c) show signals with up to two peaks, and at lower driving pressures [Figs. 9(d) and 9(f)], where the amplitudes have substantially decreased, even three or four peaks may be observed. The luminescence events occur within a time frame of approximately 200 ns.

Figures 10(a)–10(l) show the averages of the photodetector signals at three driving pressures ($\Delta p = 98, 58,$ and 38 kPa) and at four bubble energies ($E_0 = 22, 15, 9,$ and 5 mJ). Each maximum peak is set to $t = 0$ ns when the averaging is performed. The range covered by the individual signals and the standard deviations are also displayed. The more energetic bubbles show multiple peaks [Figs. 10(a)–10(f)], while at lower energies luminescence is measured as a single peak [Figs. 10(g)–10(l)]. Figures 10(c) and 10(j) display signals with similar peak amplitudes, yet the high-energy bubble collapsing at low pressure yields multiple peaks while the low-energy bubble collapsing at atmospheric pressure yields a single peak. Figures 10(d) and 10(h) display signals for bubbles

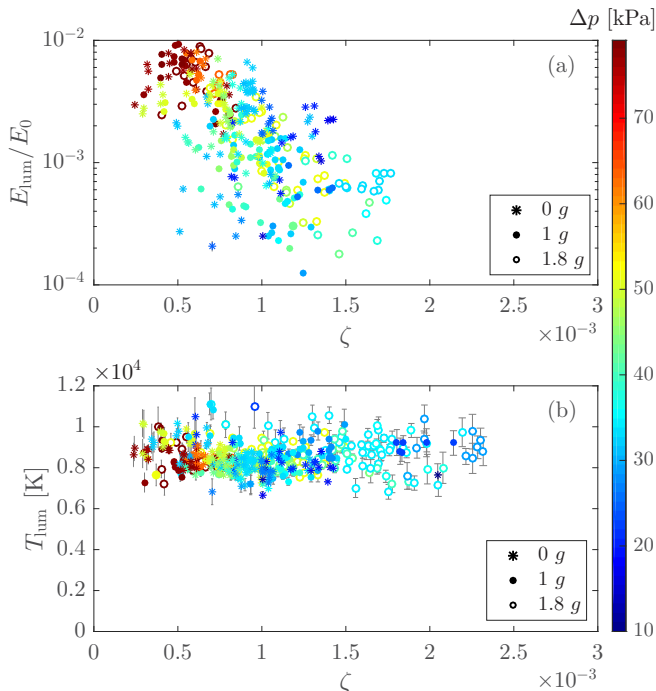


FIG. 8. Single-cavitation-bubble luminescence (a) relative energy E_{lum}/E_0 and (b) blackbody temperature T_{lum} as a function of the anisotropy parameter ζ . Each data point represents a single bubble measurement. Colors indicate driving pressures and symbols indicate different levels of gravity. Error bars indicate the $\pm\sigma$ uncertainty of the best-fit estimate of the blackbody temperature, while the error for the E_{lum}/E_0 estimate is small ($\sigma \sim 10^{-5}$).

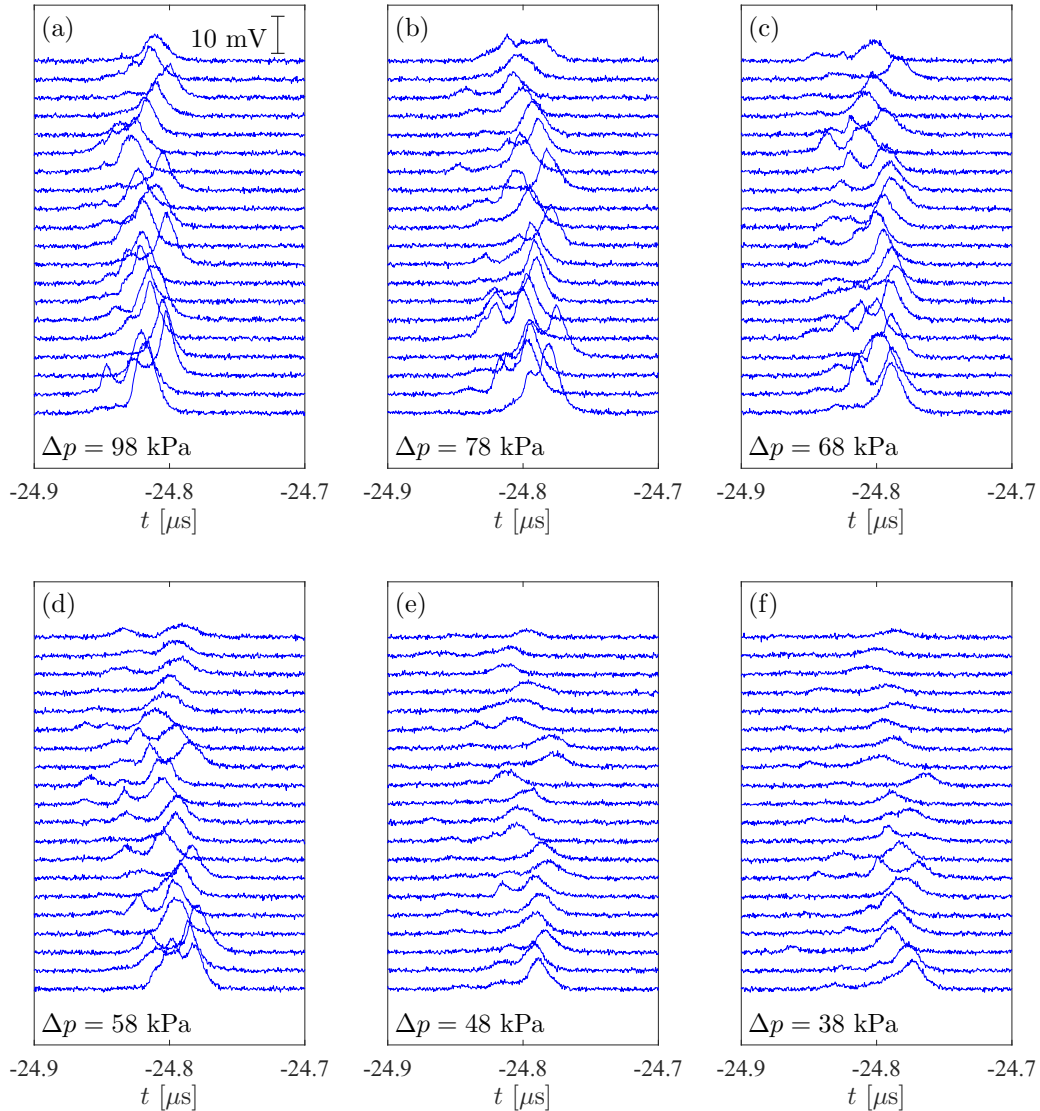


FIG. 9. Waterfall plots of the luminescence signals measured by the photodetector for different driving pressures: (a) $\Delta p = 98$ kPa ($R_0 = 3.8$ mm, $\zeta = 7.8 \times 10^{-4}$), (b) $\Delta p = 78$ kPa ($R_0 = 4.0$ mm, $\zeta = 9.0 \times 10^{-4}$), (c) $\Delta p = 68$ kPa ($R_0 = 4.2$ mm, $\zeta = 9.9 \times 10^{-4}$), (d) $\Delta p = 58$ kPa ($R_0 = 4.5$ mm, $\zeta = 1.1 \times 10^{-3}$), (e) $\Delta p = 48$ kPa ($R_0 = 4.7$ mm, $\zeta = 1.3 \times 10^{-3}$), and (f) $\Delta p = 38$ kPa ($R_0 = 5.1$ mm, $\zeta = 1.6 \times 10^{-3}$). $E_0 \approx 22$ mJ. Each plot contains 20 signals. The scaling shown in (a) is the same in all plots. Here $t = 0$ μ s corresponds to the instant at which the hydrophone detects the collapse shock. The standard deviations for R_0 and ζ are $\sigma_{R_0} \approx 0.03$ mm and $\sigma_\zeta \approx 1.5 \times 10^{-5}$, respectively. Measurements were made at normal gravity.

with the same maximum radius but with different energies, and, again, the higher-energy bubble yields more prominent additional peaks than the low-energy one. However, we find no clear correlation between the number, amplitudes, or timings of the peaks and the bubble's asphericity.

Figures 11(a) and 11(b) show the measured luminescence durations as the FWHM and the full width, which are extracted directly from the average of 20 individual photodetector signals. The full width here is defined as the duration of the averaged signal above 1% of its peak amplitude (the noise in the averaged signals has been smoothed out sufficiently not to affect this low threshold). In order to complete the graph for previously measured luminescence durations for smaller laser-induced bubbles, FWHM data from Baghdassarian *et al.* (1999) [29] and from Ohl (2002) [18] are included for purposes

of comparison. The trend for the duration of these large bubbles remains similar to that for the previously reported smaller bubbles, that is, approximately linear as a function of R_0 . While past research has suggested that the pulse duration increases for bubbles collapsing at higher pressures [22] and for bubbles deformed by a neighboring surface [18], the direct roles of Δp and ζ on the pulse duration in our results are unclear. In particular, luminescence durations at $\Delta p = 38$ kPa seem to be outliers from the general trend, with the FWHM remaining almost constant for $R_0 = 3$ –5 mm.

Finally, a typical example of an ultrahigh-speed CMOS camera recording of the luminescence is shown in Fig. 12, where luminescence events are visible in the visualization with backlight illumination [Fig. 12(a)] and in the dark [Fig. 12(b)] for a relatively deformed bubble. Figure 12(b)

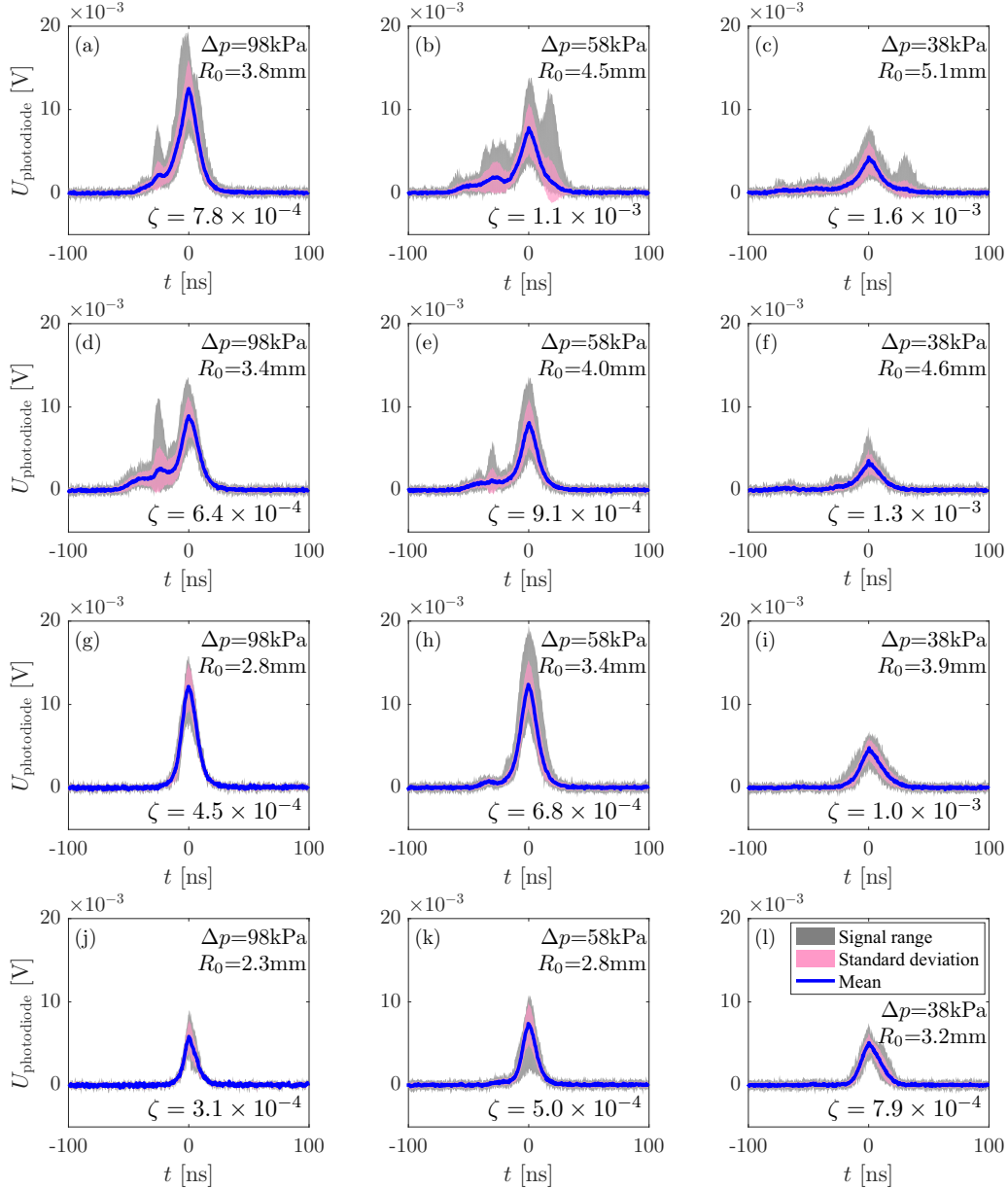


FIG. 10. Mean of 20 luminescence signals measured by the photodetector for different bubble energies—(a–c) $E_0 = 22$ mJ, (d–f) $E_0 = 15$ mJ, (g–i) $E_0 = 9$ mJ, and (j–l) $E_0 = 5$ mJ—and three driving pressures: $\Delta p = 98, 58,$ and 38 kPa. The ranges covered by the individual signals and the standard deviations are also displayed. Here $t = 0$ ns corresponds to the maximum peak.

shows the luminescent flash in the dark in two frames and thereby implies that the total luminescence event duration here exceeds the interframe time of 100 ns, consistent with the photodetector measurements [see Fig. 11(b) for $R_0 \approx 5$ mm]. The images here likely only capture the beginning and the end of the light emission, while the peak intensity occurs between the images (the exposure time 50 ns covers only half of the interframe time). In fact, the CMOS camera systematically captures the luminescent flash in two or even three consecutive frames and occasionally gets saturated. We also observe an upward shift of the light spot in the images in Fig. 12(b). This might be expected, because according to momentum conservation, most of the bubble's translational motion upon its nonspherical collapse occurs during its last collapse and

early rebound stages, when the luminescence is emitted. The bubble centroid's upward displacement during the collapse is clearly visible in Fig. 12(a).

V. DISCUSSION

The results presented here provide insight into how the topological changes in the cavity volume from a spherical to a jetting bubble affect the degree of adiabatic heating. Luminescence has an appreciable sensitivity to even the finest pressure-field anisotropies in the liquid caused by the gravity-induced pressure gradient. The threshold beyond which luminescence is no longer observed, $\zeta \approx 3.5 \times 10^{-3}$, is close to the limit where we start observing jetting bubbles in our experiment

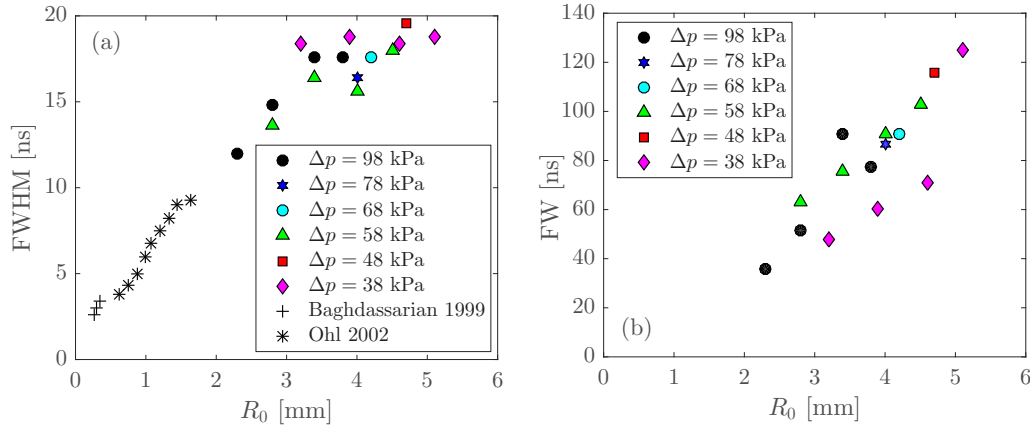


FIG. 11. (a) Full width at half-maximum (FWHM) and (b) full width (FW; with 1% of the peak amplitude as threshold) of the luminescence as a function of R_0 . The durations were extracted directly from the averaged photodetector signals of 20 bubbles. FWHM data at atmospheric pressure from Baghdadassarian *et al.* (1999) [29] and from Ohl (2002) [18] are shown for reference.

($\zeta \sim 10^{-3}$), the latter, however, being a limit that is difficult to define with precision. Considering a bubble deformation by nearby boundaries yielding an identical Kelvin impulse [$\gamma = (0.915\zeta^{-1})^{1/2}$ [14], where, equivalently, $\gamma = s/R_0$, s being the distance between the bubble and the boundary], the threshold at which we no longer detect luminescence here would be equivalent to a bubble collapsing at a distance of 7.5 times its maximum radius from the boundary. This limit disagrees with previous studies on luminescence from laser-induced bubbles deformed by near boundaries, where the equivalent limit is much lower, e.g., $\gamma \sim 3.5$ in Refs. [17] and [22] (corresponding to $\zeta \sim 0.016$). This discrepancy is possibly attributed either to different sources of deformations yielding different levels of gas compression or to the sensitivity of luminescence to the initial bubble sphericity. The latter hypothesis is supported by our previous observation that the level of deformation at which a microjet visibly pierces the bubble and drives a vapor jet during the rebound for bubbles deformed by near surfaces in our experiment ($\zeta \approx 10^{-3}$ or $\gamma \approx 14$) is also significantly lower compared to the literature (typically $\gamma \approx 5$) [14]. Likewise, we have recently measured the shock-wave energy to start being sensitive to ζ at longer distances away from surfaces ($\gamma \approx 8$ [16]) compared to the literature ($\gamma \approx 3$ [15]). As mentioned earlier, lens-based bubble generation systems, in contrast to the use of a high-convergence parabolic mirror, produce bubbles with higher surface perturbations that are amplified during the last collapse stage [38]. Consequently, a potential microjet, which can be regarded as the lowest-order deviation from a sphere and is thus most effective at inhibiting the final gas compression,

may be masked by more important, higher-order perturbations. This could make the bubble experience a collapse that perhaps more effectively compresses the gas and that is less susceptible to external factors, possibly even appearing spherical. This hypothesis is illustrated in Fig. 13.

However, luminescence can also occur for jetting bubbles, as has previously been shown for bubbles deformed by a neighboring surface [17], for acoustic cavitation bubbles in multibubble fields in xenon-saturated phosphoric acid [21], and for xenon gas bubbles collapsed by a passing shock wave [39]. A possible reason for our not observing light emission for bubbles that produced clear gravity-driven “vapor jets” upon rebound could be linked to the characteristic shape that the bubble assumes at the moment of the jet piercing. We have previously shown that, according to potential flow theory, the gravity-induced deformation yields a broad jet whose shape is very similar to that of the bubble wall it pierces [14], and thereby the gas compression after the jet impact becomes particularly weak [see Fig. 13(b)]. In contrast, when the bubble is deformed by a neighboring rigid or a free surface, at certain ranges of ζ , potential flow theory predicts small vapor “pockets” remaining between the jet and the opposite bubble wall upon the first contact of the jet with it [14], such as in the illustration in Fig. 13(c). We have previously observed luminescence from the location at which the jet pierces the bubble wall for bubbles collapsing near a free surface, as shown in Fig. 14 (adapted from Ref. [16]). This is due to the contact between the jet and the opposing wall being more irregular, which is characteristic of bubbles near free surfaces. The jet thus divides the bubble into multiple

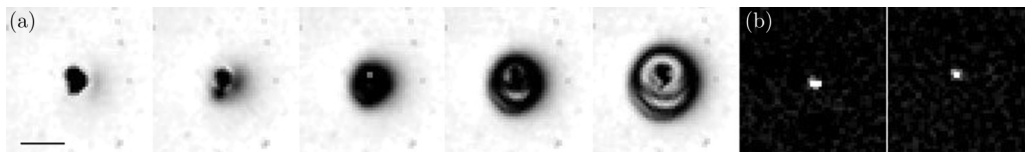


FIG. 12. Images of the luminescence emission of a bubble collapsing at $\Delta p = 38$ kPa and normal gravity ($R_0 = 5.1$ mm, $\zeta = 1.3 \times 10^{-3}$) captured with an ultrahigh-speed CMOS camera: (a) with a backlight LED and (b) in the dark. The black bar in the leftmost panel shows the 1-mm scale. The interframe time is 100 ns and the exposure time is 50 ns. The contrast and brightness of the images have been adjusted to optimize the visual clarity of events.

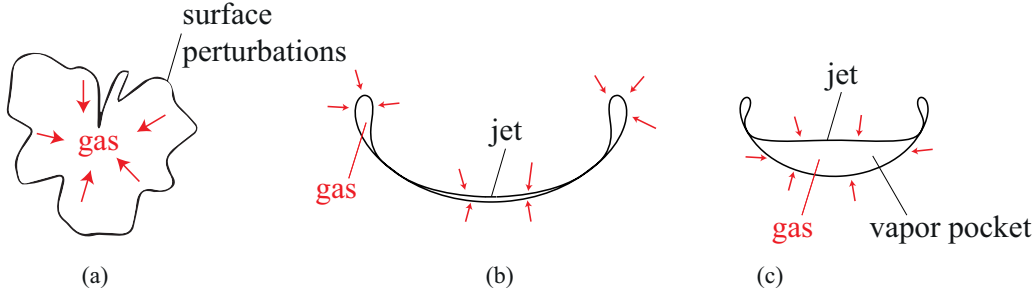


FIG. 13. Illustration of the possible effect of the bubble’s surface perturbations on its gas compression. Sketches of shapes at the final collapse stage for a bubble (a) with surface perturbations, (b) with a downward jet induced by a uniform pressure gradient, and (c) with a downward jet induced by a neighboring free surface.

separate segments, one of which is a vapor pocket between the jet and the opposite wall that is individually able to collapse in an almost spherical way, which, in turn, yields an effective compression. This hypothesis is supported by our previous observations where such vapor pockets emitted strong shocks for bubbles near a free surface [16]. However, we are unable to temporally distinguish the jet impact from the individual collapses of the remaining bubble segments at a low enough ζ for luminescence to still be visible. It would be interesting in the future to study more thoroughly the effect of the bubble shape on luminescence by varying this shape with different sources of deformation (e.g., comparing different surfaces and gravity) in a single setup.

A surprising finding is that the spectroscopically estimated blackbody temperatures of luminescence barely vary with the different levels of bubble deformation [Fig. 8(b)], while its energy varies by two orders of magnitude [Fig. 8(a)]. We do not exclude the possibility that the scatter of the data, partly caused by the fitting error, hides a possible weak variation of the blackbody temperature with ζ . However, it could also be due to the fact that as the radiation power scales as T_{lum}^4 , any attempt to increase T_{lum} immediately results in an accelerated loss of energy by radiation. Another potential physical reason could be the presence of water vapor, which increases the heat capacity ratio [1,40,41]. It has been shown numerically that for sonoluminescent bubbles that have compression ratios beyond $R_0/R_{\text{min}} \sim 20$, water vapor starts affecting the power-law

increase in the maximum temperature with the compression ratio, finally asymptoting to $T_{\text{lum}} \approx 10\,000\text{ K}$ [40]. It is difficult to measure the minimum bubble size in our experiment because the luminescence and the light deflection caused by the pressure rise in the surrounding liquid “hide” the bubble in the last stage of the collapse (see images in Fig. 4). However, when choosing the luminescent flash size as the minimum radius, we get compression ratios $R_0/R_{\text{min}} > 40$, which is already in the regime where vapor affects the heating.

The noncondensable gas trapped inside the bubble plays a key role in luminescence emission. We believe that the bubble contains (i) vapor, of which the partial pressure is assumed to stay at the liquid vapor pressure p_v during most of the bubble’s lifetime; (ii) the laser-generated gas (demonstrated in Ref. [42]), which we assume to depend on the energy deposited by the laser to generate the bubble, that is, to be proportional to $E_{\text{laser}} \propto E_0 \propto R_0^3 \Delta p$; and (iii) the diffused gas from the water to the bubble, which depends both on the total bubble surface during its lifetime, which is proportional to $R_0^3 \Delta p^{1/2}$, and on the diffusion-driving pressure Δp . Each of these likely contributes to the noncondensable gas, which is difficult to measure directly. The laser-generated and diffused gases are both proportional to the bubble’s maximum volume, while they may depend on Δp to a different extent. A method has been proposed by Tinguely *et al.* [36] to estimate the initial partial pressure of the noncondensable gas p_{g_0} by fitting the Keller-Miksis model [43] to the observed rebound. Applying this method to the observed radial evolution of spherically collapsing bubbles at various Δp values in microgravity, we can estimate the variation of p_{g_0} as a function of Δp . Our preliminary results find that p_{g_0} remains almost constant ($p_{g_0} \approx 4\text{ Pa}$) for the range of Δp covered here, differing less than the standard deviation, as illustrated in Fig. 15. Furthermore, the luminescence energy data obtained in this range of Δp in Fig. 8 suggest that the bubble’s deformation (ζ) is the dominant source of luminescence energy hindering rather than Δp , even though a weak dependence on the latter may exist. Figure 6, which shows luminescence energies as a function of the bubble energy at different Δp ’s, however, suggests some additional dependence of the luminescence energy on Δp . A systematic study with a controlled gas content of the water, preferably in microgravity to remove the effect of bubble deformation by gravity, would be useful to clarify the effect of noncondensable gas on luminescence and on other bubble collapse phenomena.

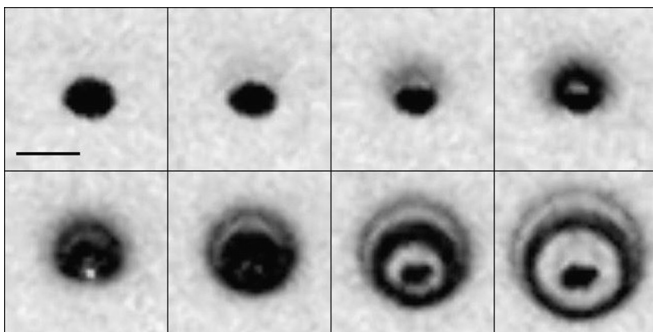


FIG. 14. Luminescence emission at the location of jet impact for a bubble collapsing near a free surface. $R_0 = 4.1\text{ mm}$, $\gamma = 7.2$. The interframe time is 100 ns and the black bar in the top-left panel shows the 1-mm scale. The microjet is directed downwards. Adapted from Ref. [16].

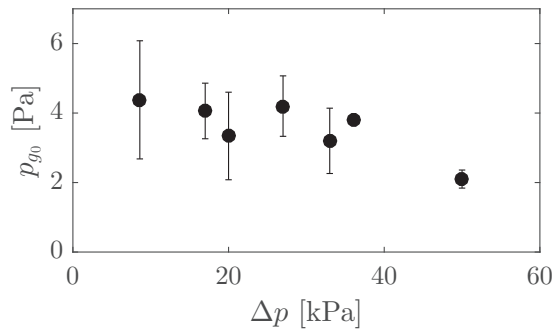


FIG. 15. Averaged initial partial pressure of the noncondensable gas, estimated by fitting the Keller-Miksis model to the observed rebound radial evolution, as a function of the driving pressure Δp . The data contain bubbles of different radii (1–3.5 mm) collapsing highly spherically ($\zeta < 0.0007$), and error bars show the standard deviation.

Finally, it would be interesting to understand the physics behind the multiple luminescence emission events that are measured by the photodetector (Figs. 9 and 10). These peaks show considerable fluctuations in their numbers, amplitudes, shapes, and timings. The timing of the strongest luminescence event with respect to the emission of the collapse shock wave is remarkably reproducible, varying by only ~ 10 ns (Fig. 9). The finding that larger bubbles emit more peaks than smaller ones is consistent with the literature, although the bubble sizes reported in the past were much smaller overall and multiple peaks were observed for bubbles with $R_0 < 2$ mm [18]. The discrepancy between our observations (single peak for $R_0 < 3$ mm) and the past literature is, again, likely due to the high initial sphericity of the bubble in our experiment. The multiple peaks could be associated with different hot spots, which could be the result of an inhomogeneous bubble interior or bubble splitting, as suggested by Ohl [18]; this would indeed be strongly affected by the initial bubble sphericity. They could also be linked to plasma instabilities, to minor impurities trapped within the bubble, or to the potential formation of

a “hidden” (nonpiercing) microjet, which is challenging to verify since the levels of deformations here are so weak.

VI. CONCLUSION

In this work, we have captured broad spectra of single-cavitation-bubble luminescence from individual collapses using an innovative measurement technique. We have measured luminescence from a previously uncovered range of maximum radii ($R_0 = 1.5\text{--}6$ mm) of laser-induced bubbles, thanks to their high initial sphericity. The bubbles were controllably deformed from highly spherical to jetting bubbles under the effect of the gravity-induced hydrostatic pressure gradient. The deformation was quantified with the dimensionless anisotropy parameter ζ , which was adjusted via the maximum bubble radius, driving pressure, and variable gravity aboard parabolic flights. We found a rapid decrease in the relative luminescence energy E_{lum}/E_0 with ζ . No clear variation of the fitted blackbody temperature, which ranged between $T_{\text{lum}} = 7000$ and $T_{\text{lum}} = 11\,500$ K, as a function of ζ or the driving pressure was found. The threshold of luminescence approximately coincides with the ζ at which we start observing vapor jets in our experiment. The light emission is found to be nonuniform in time for the most energetic bubbles, as multiple events are detected in the time-resolved measurements by a photodetector, while low-energy bubbles emit single luminescence peaks. The luminescence events were found to occur in a time frame of 200 ns. The FWHM of the averaged luminescence signal scales with R_0 and is generally of the order of 10–20 ns.

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- [1] M. Brenner, S. Hilgenfeldt, and D. Lohse, Single-bubble sonoluminescence, *Rev. Modern Phys.* **74**, 425 (2002).
- [2] Y. T. Didenko, W. B. McNamara, and K. S. Suslick, Temperature of multibubble sonoluminescence in water, *J. Phys. Chem. A* **103**, 10783 (1999).
- [3] Y. T. Didenko and K. S. Suslick, The energy efficiency of formation of photons, radicals and ions during single-bubble cavitation, *Nature* **418**, 394 (2002).
- [4] I. Rosenthal, J. Z. Sostaric, and P. Riesz, Sonodynamic therapy—A review of the synergistic effects of drugs and ultrasound, *Ultrason. Sonochem.* **11**, 349 (2004).
- [5] T. Yu, Z. Wang, and T. J. Mason, A review of research into the uses of low level ultrasound in cancer therapy, *Ultrason. Sonochem.* **11**, 95 (2004).
- [6] T. J. Mason, Sonochemistry and the environment: Providing a green link between chemistry, physics and engineering, *Ultrason. Sonochem.* **14**, 476 (2007).
- [7] Y. G. Adewuyi, Sonochemistry in environmental remediation. 2. Heterogeneous sonophotocatalytic oxidation processes for the treatment of pollutants in water, *Environ. Sci. Technol.* **39**, 8557 (2005).
- [8] K. S. Suslick and K. J. Price, Applications of ultrasound to materials chemistry, *Annu. Rev. Mater. Sci.* **29**, 295 (1999).
- [9] A. Gedanken, Using sonochemistry for the fabrication of nanomaterials, *Ultrason. Sonochem.* **11**, 47 (2004).
- [10] T. G. Leighton, M. Farhat, J. E. Field, and F. Avellan, Cavitation luminescence from flow over a hydrofoil in a cavitation tunnel, *J. Fluid Mech.* **480**, 43 (2003).
- [11] M. Farhat, A. Chakravarty, and J. E. Field, Luminescence from hydrodynamic cavitation, *Proc. R. Soc. London A: Mathe., Phys. Eng. Sci.* **480**, rspa20100134 (2010).
- [12] D. Obreschkow, M. Tinguely, N. Dorsaz, P. Kobel, A. De Bosset, and M. Farhat, The quest for the most spherical bubble: Experimental setup and data overview, *Exp. Fluids* **54**, 1503 (2013).

- [13] O. Supponen, P. Kobel, D. Obreschkow, and M. Farhat, The inner world of a collapsing bubble, *Phys. Fluids* **27**, 091113 (2015).
- [14] O. Supponen, D. Obreschkow, M. Tinguely, P. Kobel, N. Dorsaz, and M. Farhat, Scaling laws for jets of single cavitation bubbles, *J. Fluid Mech.* **802**, 263 (2016).
- [15] A. Vogel and W. Lauterborn, Acoustic transient generation by laser-produced cavitation bubbles near solid boundaries, *J. Acoust. Soc. Am.* **84**, 719 (1988).
- [16] O. Supponen, D. Obreschkow, P. Kobel, M. Tinguely, N. Dorsaz, and M. Farhat, Shock waves from non-spherical cavitation bubbles, *Phys. Rev. Fluids* **2**, 093601 (2017).
- [17] C. D. Ohl, O. Lindau, and W. Lauterborn, Luminescence from Spherically and Aspherically Collapsing Laser Induced Bubbles, *Phys. Rev. Lett.* **80**, 393 (1998).
- [18] C.-D. Ohl, Probing luminescence from nonspherical bubble collapse, *Phys. Fluids* **14**, 2700 (2002).
- [19] E. A. Brujan and G. A. Williams, Luminescence spectra of laser-induced cavitation bubbles near rigid boundaries, *Phys. Rev. E* **72**, 016304 (2005).
- [20] T. Matula, Single-bubble sonoluminescence in microgravity, *Ultrasonics* **38**, 559 (2000).
- [21] C. Cairós and R. Mettin, Simultaneous High-Speed Recording of Sonoluminescence and Bubble Dynamics in Multibubble Fields, *Phys. Rev. Lett.* **118**, 064301 (2017).
- [22] E. A. Brujan, D. S. Hecht, F. Lee, and G. A. Williams, Properties of luminescence from laser-created bubbles in pressurized water, *Phys. Rev. E* **72**, 066310 (2005).
- [23] O. Baghdassarian, H.-C. Chu, B. Tabbert, and G. A. Williams, Spectrum of Luminescence from Laser-Created Bubbles in Water, *Phys. Rev. Lett.* **86**, 4934 (2001).
- [24] L. Zhang, X. Zhu, H. Yan, Y. Huang, Z. Liu, and K. Yan, Luminescence flash and temperature determination of the bubble generated by underwater pulsed discharge, *Appl. Phys. Lett.* **110**, 034101 (2017).
- [25] W. B. McNamara, Y. T. Didenko, and K. S. Suslick, Sonoluminescence temperatures during multi-bubble cavitation, *Nature* **401**, 772 (1999).
- [26] M. C. Ramsey and R. W. Pitz, Energetic Cavitation Collapse Generates 3.2 Mbar Plasma with a 1.4 J Driver, *Phys. Rev. Lett.* **110**, 154301 (2013).
- [27] J. Duplat and E. Villermaux, Luminescence from Collapsing Centimeter Bubbles Expanded by Chemical Reaction, *Phys. Rev. Lett.* **115**, 094501 (2015).
- [28] T. J. Matula, R. A. Roy, P. D. Mourad, W. B. McNamara III, and K. S. Suslick, Comparison of Multibubble and Single-Bubble Sonoluminescence Spectra, *Phys. Rev. Lett.* **75**, 2602 (1995).
- [29] O. Baghdassarian, B. Tabbert, and G. A. Williams, Luminescence Characteristics of Laser-Induced Bubbles in Water, *Phys. Rev. Lett.* **83**, 2437 (1999).
- [30] L. Rayleigh, On the pressure developed in a liquid during the collapse of a spherical cavity, *Philos. Mag.* **34**, 94 (1917).
- [31] J. R. Blake, The kelvin impulse: Application to cavitation bubble dynamics, *J. Austral. Math. Soc. Ser. B* **30**, 127 (1988).
- [32] G. M. Hale and M. R. Query, Optical constants of water in the 200-nm to 200- μ m wavelength region, *Appl. Opt.* **12**, 555 (1973).
- [33] D. Obreschkow, M. Tinguely, N. Dorsaz, P. Kobel, A. de Bosset, and M. Farhat, Universal Scaling law for Jets of Collapsing Bubbles, *Phys. Rev. Lett.* **107**, 204501 (2011).
- [34] R. A. Hiller, S. J. Putterman, and B. P. Barber, Spectrum of Synchronous Picosecond Sonoluminescence, *Phys. Rev. Lett.* **69**, 1182 (1992).
- [35] B. Wolfrum, T. Kurz, O. Lindau, and W. Lauterborn, Luminescence of transient bubbles at elevated ambient pressures, *Phys. Rev. E* **64**, 046306 (2001).
- [36] M. Tinguely, D. Obreschkow, P. Kobel, N. Dorsaz, A. de Bosset, and M. Farhat, Energy partition at the collapse of spherical cavitation bubbles, *Phys. Rev. E* **86**, 046315 (2012).
- [37] J. R. Sukovich, A. Sampathkumar, P. A. Anderson, R. G. Holt, Y. A. Pishchalnikov, and D. F. Gaitan, Temporally and spatially resolved imaging of laser-nucleated bubble cloud sonoluminescence, *Phys. Rev. E* **85**, 056605 (2012).
- [38] M. S. Plesset and T. P. Mitchell, On the stability of the spherical shape of a vapor cavity in a liquid, *Q. Appl. Math.* **13**, 419 (1956).
- [39] B. Kappus, S. Khalid, and S. Putterman, 100-watt sonoluminescence generated by 2.5-atmosphere-pressure pulses, *Phys. Rev. E* **83**, 056304 (2011).
- [40] B. D. Storey and A. J. Szeri, Water vapour, sonoluminescence and sonochemistry, *Proc. R. Soc. London A* **456**, 1685 (2000).
- [41] R. Toegel, B. Gompf, R. Pecha, and D. Lohse, Does Water Vapor Prevent Upscaling Sonoluminescence? *Phys. Rev. Lett.* **85**, 3165 (2000).
- [42] T. Sato, M. Tinguely, M. Oizumi, and M. Farhat, Evidence for hydrogen generation in laser- or spark-induced cavitation bubbles, *Appl. Phys. Lett.* **102**, 074105 (2013).
- [43] J. B. Keller and M. Miksis, Bubble oscillations of large amplitude, *J. Acoust. Soc. Am.* **68**, 628 (1980).