

## X-ray scattering reveals molecular tilt is an order parameter for the main phase transition in a model biomembrane

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Synchrotron diffuse x-ray scattering data reveal a dramatic softening of the molecular tilt modulus  $K_\theta$  of the model biomembrane composed of DMPC lipids as the temperature is lowered towards the main phase transition temperature at  $T_M = 24^\circ\text{C}$ . Spontaneous tilt occurs below  $T_M$ , suggesting that tilt is a symmetry breaking order parameter. Consistent with this hypothesis, it is also found that a different lipid POPS has no spontaneous tilt below its  $T_M$  at  $14^\circ\text{C}$  and correspondingly its tilt modulus did not soften as  $T_M$  was approached from above. As previously known, the bending modulus  $K_C$  of DMPC also softens close to  $T_M$ , but unlike the tilt modulus,  $K_C$  has a maximum  $3^\circ$  above  $T_M$ , which also marks the limit of the well-known anomalous swelling regime. Tilt adds a different perspective to our previous understanding of the main phase transition in lipid bilayers.

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Biomembranes in living cells must be flexible, which is one reason that they are nanoscopic materials only a few nanometers thick. Mechanical descriptors are widely recognized to be central for assessing membrane properties. For transverse flexibility, the bending modulus is the main mechanical descriptor. The Helfrich-Canham (HC) theory [1] is the simplest theoretical framework that incorporates a mechanical modulus for bending. For symmetric lipid bilayers of fixed topology this is a one-parameter continuum theory in which the curvature energy is proportional to the bending modulus  $K_C$  times the membrane curvature squared. Although a recent modification has recently been proposed for gel phase bilayers [2], the HC theory is generally deemed valid for the liquid-crystalline fluid phase at long length scales.

However, simulations have made it clear that the HC theory is not valid at shorter length scales [3–5]. A growing consensus is that the continuum theory can be significantly improved by including a molecular tilt degree of freedom. Such a theory was proposed by Hamm and Kozlov (HK) [6]. Importantly for biological relevance, it was shown that the HK theory alleviated the concern [7] that the HC theory had predicted an impossibly large activation energy for the biologically essential function of the fusion of membranes [8]. Subsequently, it was shown that the HK theory also quantitatively accounts for the observed deviations in the simulated fluctuation spectra [4]. Further development of the tilt theory has been made [9], including methods for extracting both the tilt modulus and the bending modulus from simulations [5]; this further theory has also passed an additional recent test regarding how the length of the hydrocarbon tails depends on tilt [10].

For many years my laboratory has used the traditional HC model for fluctuations to analyze diffuse x-ray scattering in order to obtain the bending modulus  $K_C$ , which measures the stiffness of membranes, and the bulk modulus  $B$ , which measures the interaction between neighboring membranes in our systems consisting of stacks of membranes. Recently, we have learned how to include the molecular tilt modulus in our analysis of x-ray data [11]. We found that the data are

fit better for the tilt-dependent model, and we have reported the first experimental value of the tilt modulus  $K_\theta$  for bilayers of one type of lipid 1,2-dioleoyl-sn-glycero-3-phosphocholine (DOPC) [12]. Reanalysis that includes the tilt modulus has recently been reported for earlier x-ray scattering data collected in this laboratory [13]. This Rapid Communication presents results that have special physics significance.

Results for the tilt modulus of bilayers of the lipid 1,2-dimyristoyl-sn-glycero-3-phosphocholine (DMPC) in the fluid, chain-melted phase are shown in Fig. 1. It is noteworthy that there is good agreement with values recently obtained from simulations, namely,  $K_\theta = 40.2 \pm 2$  mN/m at  $T = 30^\circ\text{C}$  from fluctuation spectra using the CHARMM36 all-atom force field [5] and  $K_\theta = 38.8 \pm 2$  mN/m at  $T = 27^\circ\text{C}$  by studying buckling with a united atom force field [14].

However, the most striking aspect of Fig. 1, not yet investigated by simulations, is the rapid decrease in  $K_\theta$  as the main transition in DMPC is approached by lowering temperature within the fluid phase. This opens a different perspective for the main phase transition. In critical phenomena spontaneous symmetry breaking of a quantity below the critical point identifies an order parameter, and a modulus (the inverse susceptibility) of that order parameter then vanishes at a critical point. It is well known that there is spontaneous tilt in the gel phase [15,16]. Although the main transition is into the ripple phase, recent structural work [17] shows that the hydrocarbon chains in the DMPC ripple phase also have spontaneous tilt (along with other interesting features not comprehended by existing theory). The fact that the transition is ultimately first order could just mean that the usual thermodynamic trajectory only goes close to but not exactly through a critical point. It is therefore a reasonable hypothesis that tilt is an order parameter for the main phase transition in DMPC.

This hypothesis has been further tested by examining a different lipid that does not have spontaneous tilt below the main transition. As is well understood, ordered lipid hydrocarbon chains spontaneously tilt when the head group steric area is larger than the lateral area of parallel ordered chains because cooperative tilting maximizes the cohesive van der Waals energy between parallel chains [18]. Phosphoserine (PS) lipids have smaller head group volumes than

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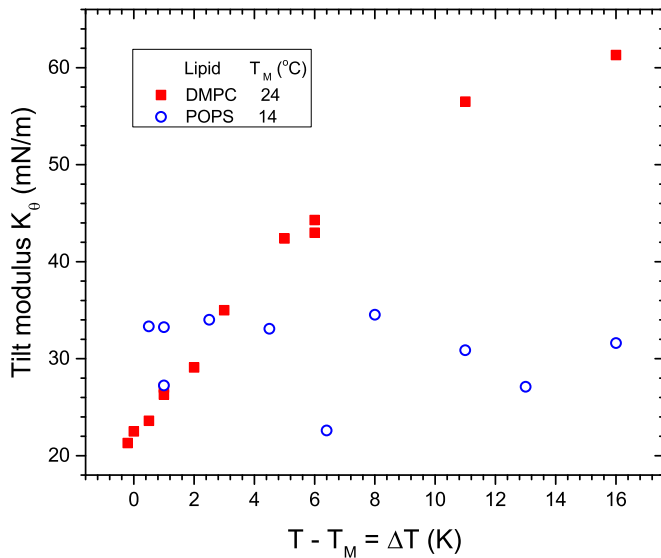


FIG. 1. Tilt modulus  $K_\theta$  vs temperature for bilayers of DMPC (solid squares) and POPS (open circles) with phase transition temperatures  $T_M$  shown in the legend. Results were obtained from analysis of diffuse x-ray scattering data [11,12].

phosphocholine (PC) lipids [19], so the chains in such lipids are less likely to tilt. For this study we used wide angle x-ray scattering shown in Fig. 2 to verify that the hydrocarbon chains of the 1-palmitoyl-2-oleoyl-sn-glycero-3-phospho-L-serine (POPS) lipid are untilted below its main phase transition temperature at  $T_M = 14^\circ\text{C}$ . The electrostatic repulsion of the singly charged POPS lipids leads to large repeat spacings ( $D \sim 150\text{--}190 \text{ \AA}$ ) in samples consisting of stacks of bilayers, and this provides enough diffuse scattering intensity for analysis even with the smaller intensity provided by an in-house rotating anode instead of the CHESS synchrotron that was used to obtain the DMPC results. Figure 1 shows that the tilt modulus of POPS has little temperature dependence upon approaching the transition. This is consistent with the hypothesis that the tilt modulus can exhibit critical behavior when it is a symmetry breaking order parameter in the low-temperature phase, but that it is unlikely to when tilt symmetry is not broken.

Figure 3 shows that our tilt-dependent analysis continues to find that the bending modulus  $K_C$  also decreases as the main transition is approached from the fluid phase [20], although not as dramatically as the tilt modulus. A decrease in  $K_C$  as  $T_M$  is approached has been correlated with an anomalous swelling in the repeat  $D$  spacing of bilayers in a stack [21]. When tilt is included in the analysis, the maximum in  $K_C$  occurs at a lower temperature ( $27^\circ\text{C}$ ) than previously ( $30^\circ\text{C}$ ) when the tilt-independent analysis was used. This maximum actually agrees better with the onset temperature of the anomalous swelling regime [20]. In contrast to DMPC,  $K_C$  for POPS (not shown) does not exhibit a maximum, but slowly increases as  $T$  approaches  $T_M$ ; gradual stiffening with reducing temperature is the ordinary behavior that one would expect *a priori* for conventional materials.

The older tilt-independent analysis essentially assumed that the tilt modulus was infinite. Allowing a finite tilt modulus in the analysis increases the obtained values of  $K_C$  as is

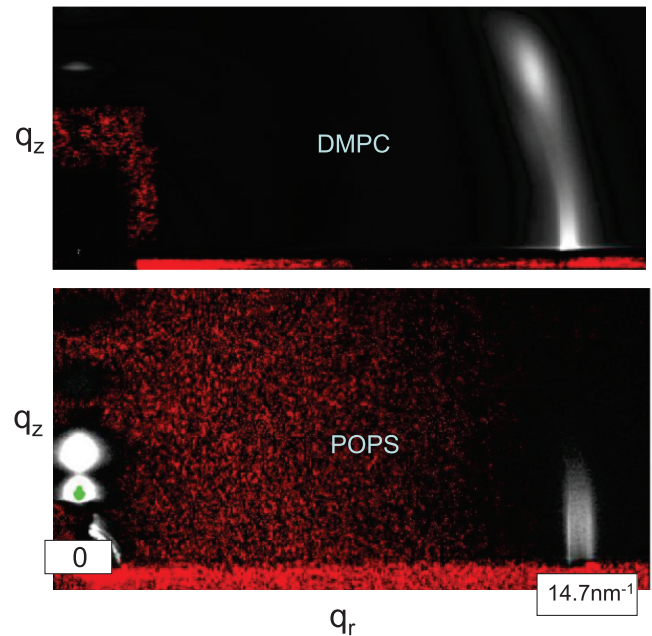


FIG. 2. Grazing incidence background subtracted scattering intensity from oriented bilayers of DMPC at  $10^\circ\text{C}$  (top) and POPS at  $5^\circ\text{C}$  (bottom). The beam is located at 0, lamellar diffraction from the stack of membranes appears along the  $q_r = 0$  meridian, blocked for DMPC in (a). Wide angle scattering from aligned hydrocarbon chains of POPS is centered at  $q_r = 1.47 \text{ \AA}^{-1}$  and  $q_z = 0$ , with negative  $q_z$  intensity cut off by the substrate. The nonzero extent in the  $q_z$  direction is due to the finite thickness of the chain region. Lipids with tilted chains such as DMPC exhibit additional wide angle peaks not centered on the  $q_z = 0$  equator [15,16].

shown in Fig. 3. This is roughly analogous to having two springs in series: Decreasing the tilt spring constant requires an increase in the bending spring constant in order to retain

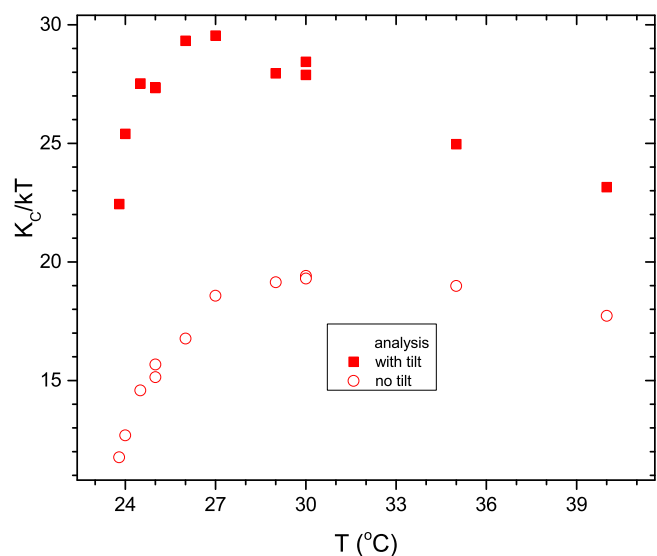


FIG. 3. Bending modulus  $K_C$  vs temperature ( $T_M = 24^\circ\text{C}$ ) for DMPC bilayers with tilt (solid squares) and without tilt (open circles) included in the analysis of the diffuse x-ray scattering data [12].

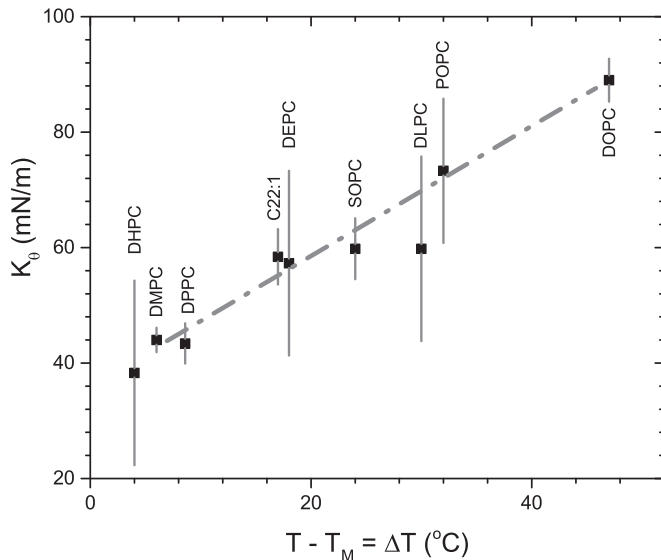


FIG. 4. Tilt modulus  $K_\theta$  vs temperature difference  $\Delta T$  between the experimental temperature and the main chain melting transition temperature  $T_M$  of each of nine common lipids with the same large PC head group that induces spontaneous tilt in gel phases [13]. See <https://avantlipids.com/> for the full chemical names, structures, and  $T_M$  of these lipids.

the same overall stiffness. However, this analogy is not perfect because, compared to  $K_C$ , smaller  $K_\theta$  softens preferentially at smaller length scales; this is fortunate because otherwise separate values for the two moduli could not be extracted from the experimental data. The increase in the value of  $K_C$  when tilt is included is greater for DMPC than for the DOPC lipid because  $K_\theta$  is smaller for DMPC than for DOPC [12]. Values obtained from this tilt-dependent analysis for ten types of lipids mostly measured at 30 °C are reported elsewhere where they are compared to values obtained from other methods and from simulations [13]. Our analysis reduces the previous differences [22] between the x-ray values of  $K_C$  and the generally larger values obtained by classical experimental methods that examine the shapes of giant unilamellar vesicles on large length scales where the Helfrich-Canham model should be accurate [23]. In contrast, neutron spin echo (NSE), as the x-ray method, measures at smaller length scales than the classical methods. NSE values of  $K_C/kT$  for DMPC range from 15 [24] to about 20 [25–27]. These agree better with the no tilt values in Fig. 3 than with our tilt-dependent values. A similar comparison also holds for the 1-palmitoyl-2-oleoyl-sn-glycero-3-phosphocholine (POPC) lipid [25,28]. This is not surprising as all the NSE analyses use the original Helfrich-Canham model; it would be interesting to have those analyses repeated using the tilt-dependent model.

Let us return to the main theme of this Rapid Communication, namely, the hypothesis that tilt is an order parameter whose modulus decreases dramatically as the main transition is approached for lipids that have spontaneous breaking of tilt symmetry below the main transition. Figure 4 shows the values of the tilt modulus for nine phosphocholine (PC) lipids versus the temperature difference  $\Delta T$  between the temperature

$T$  of measurement and the main transition temperature  $T_M$  of that lipid. The systematic trend fully supports the hypothesis.

It is well recognized that the main phase transition in single-component lipid bilayers such as DMPC is a first-order transition. However, it is well recognized in the field of critical phenomena that first-order transitions may occur in the vicinity of a critical point, in which case critical phenomena are still evident as the transition is approached. For example, in fluids, there could be a physical constraint on the pressure that makes the temperature trajectory pass close to the critical point while intersecting the line of first-order transitions that leads to the critical point along a different thermodynamic trajectory. Then, instead of the modulus coupled to the order parameter going to zero, the modulus becomes small but remains nonzero at the ultimate first-order transition temperature.

A more specific explanation is that it has been widely recognized for a long time that the main transition in lipid bilayers is primarily driven by a cooperative collapse of the disordered hydrocarbon lipid chains into the all-*trans* aligned state, as evidenced by the large transition enthalpy [29]. This chain freezing transition could occur before  $K_\theta$  vanishes. It may be noted that a toy model of that chain melting transition indicates an unusual sort of order parameter and critical point [30,31]. That toy model (which does not include tilt) also shows how other parameters can cause the system to avoid its critical point under the usual experimental conditions and undergo a first-order transition instead [32]. Furthermore, unlike typical fluid or magnetic critical points, the toy model predicts asymmetrical thermodynamic behavior above and below the first-order transition.

There was some previous experimental evidence that the main transition in DMPC, while definitely first order, was in the vicinity of a critical point. The thermal coefficient of volume expansion  $(\partial V/\partial T)_P$  increased as  $T$  approached  $T_M$  from above [33], but volume is rather weakly coupled to the driving forces for the transition, so the data were not as dramatic as they are for  $K_\theta$  in Fig. 1. Nevertheless, it would be surprising that the criticality signal for DMPC evidenced in Fig. 1 is so strong if the main driver of the transition is only chain melting. Of course, the results in this study emphasize that chain tilting is also intimately involved in the main transition in DMPC.

This brings one to theoretical modeling. The above discussion suggests that a model should definitely involve chain melting, with its area-dependent order parameter, for both DMPC and POPS. For DMPC it should now also involve chain tilting, which would have an  $XY$ /planar order parameter with “spin” dimension two and spatial dimension two in the parlance of critical phenomenology [34]. But there is even more. While the low-temperature gel phase has spontaneous tilt and frozen chains, DMPC only enters this phase 10° below the main transition. The intervening phase between the fluid phase and the gel phase is the unusual ripple phase in which DMPC has a well-defined out-of-plane asymmetric ripple; that defines yet a third-order parameter that becomes nonzero below the main transition and then vanishes again 10° lower in the gel phase. Interestingly, lipids such as POPS that do not have tilt in their low-temperature phase also do not have the intervening ripple phase [35], which suggests that there is coupling between tilting and rippling. What

is needed for PC lipids such as DMPC is a grand unified theory that accommodates this rich and varied phase behavior and that includes the recently discovered softening of the tilt modulus reported in this study as well as the recently discovered staggered monolayer structure of the DMPC ripple phase [17]. Perhaps previously proposed Landau-Ginzburg continuum models [36–38] may lead the way.

The inclusion of tilt has become important in the analysis of simulations [4,5,9]. The present experimental results for the temperature dependence of the tilt modulus are now providing

a different link between the mechanical properties and the thermodynamic phase behavior of lipid bilayers.

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