Resilience of antagonistic networks with regard to the effects of initial failures and degree-degree correlations

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In this study we investigate the resilience of duplex networked layers α and β coupled with antagonistic interlinks, each layer of which inhibits its counterpart at the microscopic level, changing the following factors: whether the influence of the initial failures in α remains [quenched (case Q)] or not [free (case F)]; the effect of intralayer degree-degree correlations in each layer and interlayer degree-degree correlations; and the type of the initial failures, such as random failures or targeted attacks (TAs). We illustrate that the percolation processes repeat in both cases Q and F, although only in case F are nodes that initially failed reactivated. To analytically evaluate the resilience of each layer, we develop a methodology based on the cavity method for deriving the size of a giant component (GC). Strong hysteresis, which is ignored in the standard cavity analysis, is observed in the repetition of the percolation processes particularly in case F. To handle this, we heuristically modify interlayer messages for macroscopic analysis, the utility of which is verified by numerical experiments. The percolation transition in each layer is continuous in both cases Q and F. We also analyze the influences of degree-degree correlations on the robustness of layer α , in particular for the case of TAs. The analysis indicates that the critical fraction of initial failures that makes the GC size in layer α vanish depends only on its intralayer degree-degree correlations. Although our model is defined in a somewhat abstract manner, it may have relevance to ecological systems that are composed of endangered species (layer α) and invaders (layer β), the former of which are damaged by the latter whereas the latter are exterminated in the areas where the former are active.

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I. INTRODUCTION

Our real world is composed of a huge variety of systems, which function in various layers, such as technology, society, and biology, and are continuously growing in unstable environments. It is thus of great importance to capture the essence of these complex critical systems. The graph [1,2] or network is one of the most powerful tools, where the constituents of the systems are regarded as nodes and the interactions between the nodes as links. Since it has been detected that networks representing real-world systems exhibit small-world properties [3,4] and scale-free (heterogeneous) properties [5] in general, various topological characteristics have been demonstrated and salient results have been reported [6–9]. One of the most important properties of a network is its robustness, that is, its tolerance to the malfunction of some nodes and/or links, which is frequently evaluated as an aggregated property, characterized as the structural phase transition of the emergence of a giant component (GC) [10]. Although vast studies have been conducted in this field, research remains insufficient, because in most studies network patterns were projected as a single layer and the effect exerted by the fact that real-world systems couple with one another was not realized.

For the purposes of analyzing real-world networks more essentially, the concept of multilayer networks was developed and is considered a new paradigm of complex network science [11–28]. The seminal work on multilayer networks is the analysis of the robustness of interdependent networks presented in [29,30]. A mutually connected GC consisting of nodes that belong to a GC in each of the layers collapses even if only a portion of the nodes has initially failed in one layer, triggering a chain of failures (called the cascade phenomenon) that spreads over all networks. This type of model may in fact be the most dependable because real-world systems in general are becoming increasingly dependent on one another [31].

A different class of multiplex networks consists of those that couple with each other with antagonistic interlayer interactions, which are called antagonistic networks. Several papers were published with regard to the robustness of antagonistic networks with neutral degree-degree correlations [32–34]. The theoretical framework was presented to analyze the robustness on duplex antagonistic networks without initial failures in [32] and later extended to include the failures in [11]. Although it was surprising that these models exhibited the first-order transition in the GC size, the definition of the antagonistic property of interlinks was artificial to some extent, in particular for numerical experiments. In our model the property of interlinks is defined simply at the microscopic level: Nodes that belong to the GC deactivate their replica nodes, while the other nodes, which do not belong to the GC, activate their replica nodes. In addition, we explicitly define that initial failures occur in layer α .

Although our model is defined in a somewhat abstract manner, one may be able to regard it as a family of graph models for ecological systems [35–39]. Employing duplex networks instead of a single network, we represent habitat patches of two categories of species (endangered species and

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the invasive ones) and interactions in and between them; the habitats of endangered species and those of invaders are projected on layer α and layer β , respectively, and the GC in each layer represents the largest and most significant habitat of the relevant layer. Each interlayer link represents the antagonistic interaction because the invaders prey on the endangered species, while the latter are conserved, thus the eradication program expels the former.

In this paper we analyze the resilience of antagonistic duplex networks that suffer from failures, in terms of the following three factors: (i) the type of the initial failures, (ii) the remaining effect of the initial failures, and (iii) degree-degree correlations. Two scenarios are examined featuring the two types of initial node failures in the the first layer α : Nodes suffer random failures (RFs) or high degree nodes selectively fail, called targeted attacks (TAs). The result of the initial failures propagates to the confronting layer β , which causes node failures at the second stage and the outcome return to the layer α . At the third stage, two possibilities are considered for the remaining effect of the initial damage. In one scenario, which is referred to as the quenched setting (case Q), the effect of the initial damage remains such that failed nodes cannot become active again. In contrast, in the second scenario, termed the free setting (case F), all the nodes are free of the initial damage, which also implies that the nodes can be reactivated with the aid of replica nodes.

In the above-described realistic scenario, case F corresponds to the situation in which endangered species can recover, while in case Q, they cannot even though invaders disappear in the area. In both cases percolation processes exhibit periodic phenomena, which were also reported in a different model [34]. In addition, we consider the effects of two types of degree-degree correlations, those between nodes within a layer (intralayer degree-degree correlations) and those between replica nodes (interlayer degree-degree correlations). In general, degrees in real-world networks are correlated [40–42] and therefore the influence of degree-degree correlations is considered one of the most important topics in the research on multilayer networks [11,43,44].

As the main part of this paper we address the development of an analytical framework based on the cavity method developed in statistical mechanics [45-48], which is categorized as a mean-field approach [11,49,50], supposing a locally treelike structure and utilizing the Bethe-Peierls approximation. In our framework, we first describe the flow for computing GC size from a microscopic viewpoint, the formulation of which is extended to a macroscopic viewpoint, and the expected GC size is analytically evaluated solving a set of self-consistent equations numerically. Unfortunately, the results obtained in this fashion deviate from numerical ones, in particular in case F. The cause of the discrepancy lies in the assumption of self-averaging property that nodes of the same degree have equivalent statistical properties, though at each single instance local states of some nodes, which are affected by the hysteresis in the layer, deterministically contribute to global property of the relevant layer, namely, GC size. The technique to efface the influence of the hysteresis is implicitly used for robustness analysis of multilayer interdependent networks [11] and antagonistic networks [32]. However, this is not valid in the latter case, unlike the former case: Some inactive nodes

are regarded as active, which may makes some of them belong to the GC, if it exists. Although the fraction of these nodes is almost negligible [51], their existence may significantly influence the critical behavior of the system, namely, whether the manner of the percolation transition is continuous [34] or discontinuous [8,32]. In keeping with the periodic phenomena, we heuristically describe the GC size at the microscopic level and extend it to the macroscopic one, the utility of which is confirmed by comparing with the numerical ones. The percolation transition of each layer turns out to be continuous in both cases Q and F in our model; in particular, that of layer α depends on the first stage and the critical point is determined only by intralayer topologies in layer α . On the other hand, the GC size depends on both interlayer and intralayer correlations, in particular the GC size in layer α exceeds about half of the layer size.

The remainder of this paper is organized as follows. In Sec. II we present the problem setup and introduce the notation that is used in our analysis. In Sec. III we develop an analytical framework for evaluating the robustness on antagonistic networks. In Sec. IV we examine the accuracy in evaluating the GC size of our methodology. We find discrepancies between theory and experiment for the GC size evaluation, particularly in case F. To resolve this inconsistency, we heuristically improve the developed methodology, which is verified by numerical experiments in Sec. V. In Sec. VI we discuss the influence of interlayer and intralayer degree-degree correlations on the robustness of each layer and suggest the relevance of real-world ecological systems that have been reported recently. In Sec. VII we conclude the paper with a summary. The periodic phenomena are reconfirmed by the heuristic in the Appendix.

II. MODEL

In this section we present a brief outline of our model of antagonistic networks consisting of layers (networks) α and β , where the number of nodes in each layer is *N*. They are generated separately in some initial configuration, where no isolated node exists in either network prior to the failure process.

Our model is seeded by initial damages that destroy a portion of the nodes in layer α , chosen uniformly at random or targeted (selected degree-dependent randomly) with probability 1 - q. As a result of the first stage, a GC may remain in layer α , the size of which is typically O(N) [or $O(N^{2/3})$ at the critical point exactly [1]. We show that nodes that belong to the GC in layer α deactivate their replica nodes in layer β , while all other nodes that do not belong to the GC activate their replica nodes, which causes the failure of nodes in layer β at the second stage, resulting in a GC in layer β that differs from the GC in layer α . Similarly, all the nodes that belong to the GC in layer β deactivate their replica nodes, while the rest of the nodes activate their replica nodes. The difference between cases Q and F corresponds to whether or not the initial damages remain in layer α at the third stage: In case Q, nodes are affected by both the initial damages and layer β , while in case F, nodes are free from the initial damage and are only affected by layer β .

The organization of this section is as follows. In Sec. II A we show that the failure process oscillates in both cases Q and F. In Sec. II B we provide the bipartite graph representations of the original networks, which are necessary for microscopic analysis in Sec. III A. In Sec. II C the topologies of each networked layer, degree distribution, and interlayer and intralayer degree-degree correlations are introduced. They are used for macroscopic analysis in Sec. III C.

A. Percolation process

In Fig. 1 we categorize the nodes in each layer into three groups and depict them as the stage elapses, which does not depend on the initial failure type (RFs or TAs) or any degree-degree correlations in and/or between networks.

(i) After the t = 1 percolation process, the nodes in layer α that constitute a GC make their replica nodes inactive at the start of stage t = 2. This guarantees that the nodes in layer α are active at the start of stage t = 3. In addition, the network topology is unchanged from stage t = 1. Therefore, it is ensured that the nodes belong to the GC after the t = 3 percolation process and repeating this argument concludes that the nodes continue to constitute the GC for ever at stages $t = 5,7,\ldots$. Accordingly, their replica nodes continue to be inactive at stages $t = 2,4,\ldots$.

(ii) The same argument guarantees that the active nodes in layer β at stage t = 2 are active at the stages t = 4, 6, ...and their replica nodes in α are never reactivated at stages t = 3, 5, ...

(iii) Statements (i) and (ii) may appear to guarantee that, when a node has been categorized as inactive, it cannot be reactivated later. However, this is not necessarily the case only in case F. This is because it is not ensured that the active nodes in layer β at stage t = 2, the replica nodes of which in layer α are inactive (damaged or isolated) at stage t = 1, form the GC, which allows a portion of the inactive nodes at t = 1 to be reactivated at stage t = 3.

Consideration of (i)–(iii) restricts possible state transitions to those depicted in Fig. 1. This figure indicates that we can terminate the repetition of the percolation at stage t = 3 in case Q and at stage t = 4 in case F, considering that the percolation processes converge.

B. Bipartite graph expression and notation

In Fig. 2 we provide the bipartite graph representations of the original networks, which help us consider the message passing scheme in the failure process graphically. Each original node is also a variable node and expressed as a circle. To indicate whether the variable node has initially failed or not without removing it, a function node is connected to each variable node, which is depicted as a closed square in the figure. For the purpose of passing messages, we append a function node that is depicted as an open slashed square expresses the role of the interlinks, whereas a function node that is depicted as an open slashed square that is depicted as an open slashed square expresses the role of the interlinks, whereas a function node that is depicted as an open square represents the role of the intralink. We now introduce the basic notation for antagonistic bipartite networks. We denote a variable node in layer α by i_{α} . The variable node i_{α} is directly connected with the set of



FIG. 1. Transitions of the states of nodes of the percolation process in antagonistic networks. Each stage is expressed as a large rounded rectangle, where red and blue represent networked layers α and β , respectively. Symbols on the left-hand side in a large rectangle express the condition of the set of nodes and those on the right-hand side represent the percolation result under the condition of the left-hand side. A cross represents the set of failed nodes at the stage, while a circle represents the set of nonfailed nodes, which are classified into two classes: (i) nodes belonging to the GC and (ii) nodes belonging to one of the small components, represented by a triangle. Dotted lines separate the groups of nodes that have different percolation results. In case Q, only the nodes that form the GC affect their replica nodes. In case F, all nodes influence their replica nodes. Possible transitions are shown for (a) case Q and (b) case F. (c) State transitions realized in the model examined in the case that node i_{β} is supposed to be the replica node of node i_{α} [11].

function nodes, which is denoted by ∂i_{α} . We denote a function node on each intralink in layer α by a_{α} and we denote a function node on interlinks by p. The function node a_{α} is directly connected with two variable nodes, denoted by ∂a_{α} .



FIG. 2. Diagram of antagonistic duplex networks.

C. Statistical expression of graph topologies

One of the most fundamental topologies of network (layer) is degree distribution, which is defined as the probability that a randomly chosen node has degree k_{α} , denoted by $p_{\alpha}(k_{\alpha})$. We also provide $r_{\alpha}(k_{\alpha})$, which denotes the degree distribution of a link computed as the probability that one terminal node of a randomly chosen link has degree k_{α} . We describe $r_{\alpha}(k_{\alpha})$ using $p_{\alpha}(k_{\alpha})$:

$$r_{\alpha}(k_{\alpha}) = \frac{k_{\alpha} p_{\alpha}(k_{\alpha})}{\sum_{l_{\alpha}} l_{\alpha} p_{\alpha}(l_{\alpha})}.$$
(1)

Related to this, the intralayer joint degree distribution (intralayer degree-degree correlations) is defined as the probability that, given an intralink is randomly chosen, one terminal node has degree k_{α} and the second terminal node has degree l_{α} , which is denoted by $r_{\alpha}(k_{\alpha}, l_{\alpha})$ and is described using $r_{\alpha}(k_{\alpha})$,

$$r_{\alpha}(k_{\alpha}) = \sum_{l_{\alpha}} r_{\alpha}(k_{\alpha}, l_{\alpha}).$$
⁽²⁾

Related to the intralayer degree-degree correlations, the intrajoint degree distribution (intralayer degree-degree correlations) is described as

$$r_{\alpha}(k_{\alpha}|l_{\alpha}) = \frac{r_{\alpha}(k_{\alpha},l_{\alpha})}{r_{\alpha}(l_{\alpha})}.$$
(3)

The interjoint degree distribution is denoted by $P(k_{\alpha},k_{\beta})$, defined as the probability that the degrees of a randomly chosen node pair are k_{α} and k_{β} and are called interlayer degree correlations. The relationship between $p_{\alpha}(k_{\alpha})$ and $P(k_{\alpha},k_{\beta})$ is

$$p_{\alpha}(k_{\alpha}) = \sum_{k_{\beta}} P(k_{\alpha}, k_{\beta}).$$
(4)

Related to interlayer joint degree distribution, the interconditional distribution is described as

$$P_{\alpha}(k_{\alpha}|k_{\beta}) = \frac{P_{\alpha}(k_{\alpha},k_{\beta})}{\sum_{k_{\beta}} P_{\alpha}(k_{\alpha},k_{\beta})},$$
(5)

which is defined as the probability of a node having degree k_{α} , given that the degree of its replica node is k_{β} .

Exchanging α with β in Eqs. (1)–(5), we define $r_{\beta}(k_{\beta})$, $r_{\beta}(k_{\beta}|l_{\beta})$, $p_{\beta}(k_{\beta})$, and $P_{\beta}(k_{\beta}|k_{\alpha})$, respectively. Using the definition of $P_{\beta}(k_{\beta}|k_{\alpha})$, the intralayer conditional distributions of node pairs are obtained as

$$r_{\alpha}(k_{\alpha},k_{\beta}|l_{\alpha},l_{\beta}) = P_{\beta}(k_{\beta}|k_{\alpha})r_{\alpha}(k_{\alpha}|l_{\alpha}), \tag{6}$$

the definition of which is the probability that a randomly chosen node pair having degree (l_{α}, l_{β}) is connected with another node pair having degree (k_{α}, k_{β}) , given an intralink in layer α .

III. THEORETICAL FRAMEWORK

The aim of this section is to develop a framework to analyze the robustness of antagonistic networks based on the cavity method. In preparation for evaluating the GC size from the macroscopic viewpoint, we examine the message flow for a single instance from the microscopic viewpoint in Sec. III A; Sec. III A 1 provides the message flow at stages t = 1 and t = 2and Secs. III A 2 and III A 3 describe how the flow behaves for $t \ge 3$ in cases Q and F, respectively. In Sec. III B we define macroscopic intralayer messages using microscopic ones. With the aid of the local tree approximation and self-averaging properties of a random network, we extend our formulation to the macroscopic level in Sec. III C; Sec. III C 1 provides the macroscopic message flow at stages t = 1 and 2 and Secs. III C 2 and III C 3 describe how the macroscopic flow behaves for $t \ge 3$ in cases Q and F, respectively.

A. Message flow from the microscopic viewpoint

1. First and second stages

We define a binary variable $\psi_{i_{\alpha}}$, which is set at 0 or 1 depending on whether or not the node suffers from the initial damage, respectively, and assign it to another variable, called the activity index $s_{i_{\alpha}}^{t=1}$ of i_{α} at stage t = 1,

$$s_{i_{\alpha}}^{t=1} \equiv \psi_{i_{\alpha}}.$$
 (7)

Note that the total fraction of the active nodes at the initial condition, namely, $\sum_{i_{\alpha}} \delta(\psi_{i_{\alpha}} = 1)/N$, is handled as a survival ratio in Sec. III C. To examine the first stage in layer α , we apply the cavity method presented in [47] for the given set of $\{s_{i_{\alpha}}^{1}\}$, which yields a set of self-consistent equations

$$m_{a_{\alpha} \to i_{\alpha}}^{t=1} = m_{j_{\alpha} \to a_{\alpha}}^{t=1} \quad (\partial a_{\alpha} = \{i_{\alpha}, j_{\alpha}\}), \tag{8}$$

$$m_{i_{\alpha} \to a_{\alpha}}^{t=1} = 1 - s_{i_{\alpha}}^{t=1} + s_{i_{\alpha}}^{t=1} \prod_{b_{\alpha} \in \partial i_{\alpha} \setminus a_{\alpha}} m_{b_{\alpha} \to i_{\alpha}}^{t=1}.$$
 (9)

Here $m_{i \rightarrow a}^{t} \in \{0, 1\}$ in general is the message from variable node *i* to function node *a* at the *t*th stage, which takes 0 when *i* belongs to a GC in the layer from which node *a* is removed and unity otherwise. The message $m_{a \rightarrow i}^{t} \in \{0, 1\}$, on the other hand, conveys 0 from function node *a* to variable node *i* at the *t*th stage when at least one $j \in \partial a \setminus i$ belongs to the GC, and unity otherwise. Using the solution of Eqs. (8) and (9), we derive the indices of the GC and the size of the GC in layer α at stage t = 1 as

$$\sigma_{\alpha}^{t=1} = \sum_{i_{\alpha}} \sigma_{i_{\alpha}}^{t=1} = \sum_{i_{\alpha}} s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right), \quad (10)$$

which also provides a message from i_{α} to the function node p on an interlink at stage t = 1 as

$$m_{i_{\alpha} \to p}^{t=1} = \sigma_{i_{\alpha}}^{t=1} = s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right).$$
(11)

Because of the antagonistic nature of the interlinks, the inverted value of Eq. (11) is propagated from the function node p to the replica node i_{β} of layer β after stage t = 1 as

$$1 - m_{i_{\alpha} \to p}^{t=1} = 1 - s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right)$$
$$= m_{p \to i_{\beta}}^{t=2}.$$
(12)

The second stage (t = 2) is considered the initial step for layer β . In contrast to the first stage at layer α , a particular set of $\psi_{i_{\beta}}$ is not involved (in other words $\psi_{i_{\beta}} = 1$), because the nodes in layer β are free from the initial damage and influenced only by the activity pattern of layer α at stage t = 1. The activity index of i_{β} at the start of stage t = 2 is

$$s_{i_{\beta}}^{t=2} = m_{p \to i_{\beta}}^{t=2} = 1 - s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right).$$
(13)

Note that $s_{i_{\beta}}^{t=2} = 1$ holds if either $\psi_{i_{\alpha}} = 0$ or $\psi_{i_{\alpha}} = 1$ and $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} = 0$ are satisfied. Given $\{s_{i_{\beta}}^{t=2}\}$, the cavity method provides the self-consistent equations

$$m_{a_{\beta} \to i_{\beta}}^{t=2} = m_{j_{\beta} \to a_{\beta}}^{t=2} \quad (\partial a_{\beta} = \{i_{\beta}, j_{\beta}\}), \tag{14}$$

$$m_{i_{\beta} \to a_{\beta}}^{t=2} = 1 - s_{i_{\beta}}^{t=2} + s_{i_{\beta}}^{t=2} \prod_{b_{\beta} \in \partial i_{\beta} \setminus a_{\beta}} m_{b_{\beta} \to i_{\beta}}^{t=2}.$$
 (15)

The solution of Eqs. (14) and (15) provides the GC size of layer β at stage t = 2 as

$$\sigma_{\beta}^{t=2} = \sum_{i_{\beta}} \sigma_{i_{\beta}}^{t=2} = \sum_{i_{\beta}} s_{i_{\beta}}^{t=2} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right).$$
(16)

The solution also provides the messages from i_{β} to p and the message from p to i_{α} as

$$m_{i_{\beta} \to p}^{t=2} = s_{i_{\beta}}^{t=2} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right), \tag{17}$$

$$m_{p \to i_{\alpha}}^{t=3} = 1 - m_{i_{\beta} \to p}^{t=2}.$$
 (18)

At the third and further stages, nodes in layer α receive the intermessages that are represented as Eq. (18).

2. Third and further stages in case Q

As already discussed in Sec. II A, we obtain the final robustness of layer α and layer β , which is the robustness

of layer α at the first stage (10) and that of layer β at the second stage (16), respectively,

$$\sigma_{i_{\alpha}}^{2t'+1} = \sigma_{i_{\alpha}}^{t=1}, \quad \sigma_{i_{\beta}}^{2t'} = \sigma_{i_{\beta}}^{t=2}(\mathbf{Q}).$$
(19)

3. Third and further stages in case F

In case F, nodes in layer α at stage t = 3 are influenced by only interlayer messages and are free from the initial activity indices, which provides the activity index of i_{α} as

$$s_{i_{\alpha}}^{t=3}(\mathbf{F}) = m_{p \to i_{\alpha}}^{t=3} = 1 - s_{i_{\beta}}^{t=2} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right).$$
(20)

Substituting $s_{i_{\alpha}}^{t=3}(F)$ for $s_{i_{\alpha}}^{t=1}$ in Eqs. (8) and (9),

$$m_{a_{\alpha} \to i_{\alpha}}^{t=3} = m_{j_{\alpha} \to a_{\alpha}}^{t=3} \quad (\partial a_{\alpha} = \{i_{\alpha}, j_{\alpha}\}), \tag{21}$$

$$m_{i_{\alpha} \to a_{\alpha}}^{t=3} = 1 - s_{i_{\alpha}}^{t=3}(\mathbf{F}) + s_{i_{\alpha}}^{t=3}(\mathbf{F}) \prod_{b_{\alpha} \in \partial i_{\alpha} \setminus a_{\alpha}} m_{b_{\alpha} \to i_{\alpha}}^{t=3}, \quad (22)$$

we obtain the solution $m_{a_{\alpha} \to i_{\alpha}}^{t=3}$, which derives the size of the GC in layer α as

$$\sigma_{\alpha}^{t=3} = \sum_{i_{\alpha}} \sigma_{i_{\alpha}}^{t=3} = \sum_{i_{\alpha}} s_{i_{\alpha}}^{t=3}(\mathbf{F}) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right).$$
(23)

We describe $s_{i_{\alpha}}^{t=4}(F)$, which denotes the message through the interlink, which each i_{β} receives at stage t = 4, as

$$s_{i_{\alpha}}^{t=4}(\mathbf{F}) = m_{p \to i_{\beta}}^{t=4} = m_{i_{\alpha} \to p}^{t=3} = 1 - \sigma_{i_{\alpha}}^{t=3}$$
$$= 1 - s_{i_{\beta}}^{t=3}(\mathbf{F}) \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=3} \right).$$
(24)

As discussed in Sec. II A, the percolation result at stage t = 4 becomes identical to that at stage t = 2 in case F. Therefore, we can determine the indices of each node in each network as

$$\sigma_{i_{\alpha}}^{2t'+1} = \sigma_{i_{\alpha}}^{t=3}, \quad \sigma_{i_{\beta}}^{2t'} = \sigma_{i_{\beta}}^{t=2}(F).$$
(25)

Consequently, we can terminate the repetition of the percolation at stage t = 2 by considering the networks converged [Fig. 1(a)].

The local message flows are categorized with the aid of the bipartite graph expression that is introduced in Sec. II B (see Fig. 3).

B. Cross-link from the microscopic viewpoint to the macroscopic one

We first focus on a pair of nodes i_{α} and i_{β} , the degrees of which are l_{α} and l_{β} , respectively. The node i_{α} is initially attached $s_{i_{\alpha}}^{t=1}$, which is described as Eq. (7). Using $s_{i_{\alpha}}^{t=1}$, we evaluate $q_{l_{\alpha}l_{\beta}}^{t,t=1}$, which denotes the fraction of the set of node pairs, the degrees of which are l_{α} and l_{β} , taking the value of unity at the first stage:

$$q_{l_{\alpha},l_{\beta}}^{\alpha,t=1} = \frac{\sum_{i_{\alpha}} \delta(|\partial i_{\alpha}| = l_{\alpha}) \delta(|\partial i_{\beta}| = l_{\beta}) s_{i_{\alpha}}^{\alpha,t=1}}{\sum_{i_{\alpha}} \delta(|\partial i_{\alpha}| = l_{\alpha}) \delta(|\partial i_{\beta}| = l_{\beta})}.$$
 (26)



FIG. 3. Diagram of message passing. (a) Message flow passing a function node in a layer, which corresponds to Eqs. (8), (14), and (21). (b) Message flow passing a variable node in a layer, which corresponds to Eqs. (9), (15), and (22). (c) Message flow for computing the size of the GC, which corresponds to Eqs. (10), (16), and (23). (d) Message flow from a variable node in a layer to a function node on an interlink, which corresponds to Eqs. (11) and (17). (e) Message flow from a function node on an interlink to a variable node in the layer, which corresponds to Eq. (12) or (18).

In the case of layer β , we define the relevant active probability as

$$q_{l_{\alpha},l_{\beta}}^{\beta,t=2} = \frac{\sum_{i_{\beta}} \delta(|\partial i_{\alpha}| = l_{\alpha})\delta(|\partial i_{\beta}| = l_{\beta})s_{i_{\beta}}^{\beta,t=2}}{\sum_{i_{\beta}} \delta(|\partial i_{\alpha}| = l_{\alpha})\delta(|\partial i_{\beta}| = l_{\beta})}.$$
 (27)

As in Eq. (26), we implicitly define $q_{l_{\alpha},l_{\beta}}^{\alpha,t=3}$ using $s_{i_{\alpha}}^{\alpha,t=3}$, which leads to Eq. (39). We compute the fraction of message $m_{j_{\alpha}\rightarrow a_{\alpha}}^{t=1} = m_{a_{\alpha}\rightarrow i_{\alpha}}^{t=1}$ taking unity, which is characterized by the degree of i_{α} and its replica node i_{β} , solving the relevant cavity equations [e.g., Eq. (33)] iteratively. We denote the macroscopic message by $I_{l_{\alpha},l_{\beta}}^{\alpha,t=1}$,

$$I_{l_{\alpha},l_{\beta}}^{\alpha,t=1} = \frac{\sum_{i_{\alpha}} \delta(|\partial i_{\alpha}| = l_{\alpha}) \sum_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{\alpha,t=1}}{l_{\alpha} \sum_{i_{\alpha}} \delta(|\partial i_{\alpha}| = l_{\alpha}) \delta(|\partial i_{\beta}| = l_{\beta})}.$$
 (28)

Similarly, in layer β ,

$$I_{l_{\alpha},l_{\beta}}^{\beta,t=2} = \frac{\sum_{i_{\beta}} \delta(|\partial i_{\beta}| = l_{\beta}) \sum_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{\beta,t=2}}{l_{\beta} \sum_{i_{\beta}} \delta(|\partial i_{\alpha}| = l_{\alpha}) \delta(|\partial i_{\beta}| = l_{\beta})}.$$
 (29)

Note that these macroscopic variables are relevant for network ensembles and thus may not be appropriate for each individual network; in particular, each active label of all the nodes is strongly correlated with the network connectivity.

C. Message flow from the macroscopic viewpoint

1. First and second stages

Let us consider the initial stage (t = 1) in layer α . First, we set only one parameter q as the initial survival probability, that is, the fraction of nonfailed nodes in layer α . In the case of RFs, we set

$$q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} = q_{k_{\alpha}}^{\alpha,t=1} = q.$$
 (30)

On the other hand, in the case of TAs, nodes having larger degrees in layer α are preferentially damaged. We therefore set

$$q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} = q_{k_{\alpha}}^{\alpha,t=1} = \begin{cases} 0 & (k_{\alpha} > \Theta) \\ \Delta & (k_{\alpha} = \Theta) \\ 1 & (k_{\alpha} < \Theta), \end{cases}$$
(31)

where Θ and Δ are uniquely determined so that

$$q = \sum_{l_{\beta}} \left(\Delta P(\Theta, l_{\beta}) + \sum_{l_{\alpha} < \Theta} P(l_{\alpha}, l_{\beta}) \right)$$
(32)

holds. Substituting $q_{k_{\alpha},k_{\beta}}^{\alpha,t=1}$ in the self-consistent equation,

$$I_{l_{\alpha},l_{\beta}}^{\alpha,t=1} = \sum_{k_{\alpha},k_{\beta}} r_{\alpha}(k_{\alpha},k_{\beta}|l_{\alpha},l_{\beta})$$
$$\times \left[1 - q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} + q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \left(I_{k_{\alpha},k_{\beta}}^{\alpha,t=1}\right)^{k_{\alpha}-1}\right], \quad (33)$$

which corresponds to Eqs. (8) and (9). Solving Eq. (33) iteratively, the solution $I_{l_{\alpha},l_{\beta}}^{\alpha,t=1}$ is determined, which offers the fraction of GC in layer α at stage t = 1,

$$\mu_{\alpha}^{t=1} = \sum_{k_{\alpha},k_{\beta}} P(k_{\alpha},k_{\beta}) q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \left[1 - \left(I_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \right)^{k_{\alpha}} \right].$$
(34)

Let us consider the second stage in layer β , in which $q_{k_{\alpha},k_{\beta}}^{\beta,t=2}$ denotes the probability that nodes do not fail at stage t = 2, the degree of which is k_{β} ; their replica node's degree is k_{α} . Considering the message flow [Eqs. (11)–(13)], $q_{k_{\alpha},k_{\beta}}^{\beta,t=2}$ is directly calculated from the solution $I_{l_{\alpha},l_{\beta}}^{\alpha,t=1}$ in Eq. (33) as

$$q_{k_{\alpha},k_{\beta}}^{\beta,t=2} = 1 - q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} + q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \left(I_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \right)^{k_{\alpha}}.$$
 (35)

Substituting $q_{k_{\alpha},k_{\beta}}^{\beta,t=2}$ in the self-consistent equation

$$I_{l_{\alpha},l_{\beta}}^{\beta,t=2} = \sum_{k_{\alpha},k_{\beta}} r_{\beta}(k_{\alpha},k_{\beta}|l_{\alpha},l_{\beta}) \\ \times \left[1 - q_{k_{\alpha},k_{\beta}}^{\beta,t=2} + q_{k_{\alpha},k_{\beta}}^{\beta,t=2} (I_{k_{\alpha},k_{\beta}}^{\beta,t=2})^{k_{\beta}-1}\right], \quad (36)$$

based on Eqs. (14) and (15), we compute the set of messages $I_{l_{\alpha},l_{\beta}}^{\beta,t=2}$, which offers the fraction of the GC in layer β at stage t = 2 as

$$\mu_{\beta}^{t=2} = \sum_{k_{\alpha},k_{\beta}} P(k_{\alpha},k_{\beta}) q_{k_{\alpha},k_{\beta}}^{\beta,t=2} \left[1 - \left(I_{k_{\alpha},k_{\beta}}^{\beta,t=2} \right)^{k_{\beta}} \right].$$
(37)

2. Third and further stages in case Q

In Sec. II A we already discussed that the final GC in layer α is identical to that at stage t = 1, while the final GC in layer β is identical to that at stage t = 2, which provides

$$\mu_{\alpha}(\mathbf{Q}) = \mu_{\alpha}^{t=1}, \quad \mu_{\beta}(\mathbf{Q}) = \mu_{\beta}^{t=2}.$$
 (38)

3. Third and further stages in case F

Here we naively compute $q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3}$, the fraction of nodes that have not failed at stage t = 3, based on Eq. (20),

$$q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3} = 1 - q_{k_{\alpha},k_{\beta}}^{\beta,t=2} + q_{k_{\alpha},k_{\beta}}^{\beta,t=2} \left(I_{k_{\alpha},k_{\beta}}^{\beta,t=2} \right)^{k_{\beta}}.$$
 (39)

As discussed in Sec. IV, it is necessary to evaluate Eq. (20) in detail. Substituting $q_{k_{\alpha},k_{\beta}}^{\alpha(F),t=3}$ in the self-consistent equation based on Eqs. (21) and (22),

$$I_{l_{\alpha},l_{\beta}}^{\alpha,t=3} = \sum_{k_{\alpha},k_{\beta}} r_{\alpha}(k_{\alpha},k_{\beta}|l_{\alpha},l_{\beta})$$
$$\times \left[1 - q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3} + q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3} (I_{k_{\alpha},k_{\beta}}^{\alpha,t=3})^{k_{\alpha}}\right], \quad (40)$$

we obtain a solution of $I_{k_{\alpha},k_{\beta}}^{\alpha,t=3}$, which yields the fraction of the GC in layer α

$$\mu_{\alpha}^{t=3} = \sum_{k_{\alpha},k_{\beta}} P(k_{\alpha},k_{\beta})q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3} \left[1 - \left(I_{k_{\alpha},k_{\beta}}^{\alpha,t=3}\right)^{k_{\alpha}}\right].$$
(41)

The GC in layer β at stage t = 4 is identical to that at stage t = 2, which means that the percolation process is only the repetition of the stage at stage t = 3 and the stage at stage t = 4 alternately. Therefore, the robustness of layer α and layer β is evaluated as

$$\mu_{\alpha}(\mathbf{F}) = \mu_{\alpha}^{t=3}, \quad \mu_{\beta}(\mathbf{F}) = \mu_{\beta}^{t=2},$$
 (42)

respectively.

IV. NUMERICAL TEST

A. Procedure

We conducted numerical experiments to confirm the validity of the method developed for analyzing the robustness of antagonistic networks. Here the procedure of the numerical experiments is briefly described.

(i) We constructed two random networks (layers) α and β , each of size $N = 10\,000$. The degree distribution of each layer was represented by $P_{\alpha}(4) = P_{\beta}(4) = 0.5$ and $P_{\alpha}(6) = P_{\beta}(6) = 0.5$.

(ii) To introduce intralayer degree-degree correlations, we set a Pearson coefficient in each layer C_{α} and C_{β} , respectively. For each layer, we randomly selected two pairs of connected nodes and rewired the intralinks, employing the algorithm in [40].

(iii) To introduce interlayer degree-degree correlations, we set a Pearson coefficient between layers C_I . We rewired the interlinks, reordering the indices of one layer. Note that it is necessary to suppose that $P(k_{\alpha} = x, k_{\beta} = y) = P(k_{\alpha} = y, k_{\beta} = x)$, because C_I does not determine $P(k_{\alpha}, k_{\beta})$ uniquely.

(iv) For the degree-correlated networks, we applied the Monte Carlo simulation described below. Setting an initial survival probability q, we chose initially failed nodes randomly depending on the type of failures (RFs or TAs). Failed nodes cause networks to decompose into connected components, each of which is detected using the algorithm in [52,53]. Note that we modified the open MATLAB code in [53] very slightly in the part of "case 4c ii," tracing the nest of "NodeLP(N)" sufficiently to reach its source such that "NodeLPmin" should be labeled the smallest cluster. We terminated the single



FIG. 4. Analytical results of the robustness of layer α versus q for the case where antagonistic networks suffer from RFs and the setting is case Q. Examples of the robustness of layer α at stage t = 1 are also plotted, which are the results of failure processes that are completed in the single layer α . Each dot is averaged 50 times, produced by numerical experiments.

instance if each active label of all nodes in layer α accorded with that at the last stage in a one-to-one manner. Similar procedures were tested 50 times at each initial survival probability q.

B. Methodological accuracy

1. Case Q

Figure 4 shows a comparison of the theoretical prediction obtained for the analysis and the experimental results, which exhibits excellent consistency. In particular, antagonistic interlinks do not affect the robustness in layer α , which is the GC at stage t = 1.

2. Case F

Figure 5 shows a comparison of the theoretical predictions of the robustness of layers α and β and the numerical results in the condition of case F. Experimental data for layer β exhibit excellent agreement with the theoretical predictions. However, with regard to the robustness in layer α , there exist significant discrepancies between the theoretical predictions and the numerical results.

V. ACCURACY IMPROVEMENT

The results in Sec. IV indicate that there are significant discrepancies in the GC size evaluation between theory and experiment for case F. The purpose of this section is to resolve this inconsistency. In Sec. V A we examine the cause of the discrepancies. To improve the accuracy of the theoretical evaluation, we derive alternative expressions for the microscopic equations (20)-(23) and corresponding



FIG. 5. Analytical results of the robustness of layer α versus q for the case where the antagonistic networks suffer from RFs and the setting is case F. Each dot is averaged 50 times, produced by numerical experiments.

macroscopic equations (39)-(41). Consequently, we introduce a heuristic treatment in the microscopic level in Sec. VB and extend it to derive macroscopic expressions in Sec. VC, the utility of which is verified by numerical experiments.

A. Clarifying the cause of discrepancy

The causes of the above discrepancies between theory and experiment lie in the transformation from microscopic variables to macroscopic ones at stage t = 3. To reach this resolution, we dissect the GC at stage t = 3, which is constitutively heterogeneous and divided into three subsets depending on the history of nodes.

(I) Nodes that belonged to the GC at stage t = 1. These nodes necessarily belonged to the GC at stage t = 3, which implies that strong correlations existed between $s_{i_{\alpha}}^{t=3}$ and $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ and thus $q_{k_{\alpha},k_{\beta}}^{\alpha(\mathrm{F}),t=3}$ and $I_{k_{\alpha},k_{\beta}}^{\alpha,t=3}$. Note that the last statement holds if and only if node i_{α} is classified as this class.

(II) Nodes that failed at stage t = 1. These nodes generally belonged to the GC by themselves.

(III) Nodes that belonged to one of the small components at stage t = 1. These nodes belonged to the GC at t = 3 with the aid of the node(s) of (II).

Because of the heterogeneity due to the hysteresis, the assumption that nodes of the same degree are statistically equivalent does not hold from the macroscopic viewpoint at stage t = 3. Therefore, Eq. (41), in which $q_{k_{\alpha},k_{\beta}}^{\alpha(F),t=3}$ and $I_{k_{\alpha},k_{\beta}}^{\alpha,t=3}$ are independent of each other, underestimates the GC size [see Figs. 6(a) and 6(b)].

In order not to treat the heterogeneity argued above, we drop the term of the past intralayer messages to obscure (or encapsulate) the connectivity in the relevant layer in Eq. (20). Note that, unlike interdependent networks, this treatment has a





FIG. 6. Possible transitions for case F from (a) a microscopic, (b) a naive macroscopic, and (c) a modified macroscopic viewpoint. Note that, for simplicity, we considered the situations where the number of active nodes that were isolated from the GC in layer β was almost negligible. The meaning of each symbol is similar to that in Fig. 1. The red triangle in (b) represents the nodes that are not judged constituents of the GC, which causes the discrepancies in Fig. 5. The red circle in (c) represents the nodes that are actually inactive according to the rule of antagonistic interlinks [see also (a)], but regarded active by dropping the self-feedback term from Eq. (43) as Eq. (44).

possibility of considering that some failed nodes belong to the GC in layer α at stage t = 3 in antagonistic networks, which depends on the percolation result of layer β at stage t = 2 (see Fig. 7). To compensate for this inconsistency, we deduct them

(a)

Х

(b)

Х

(c)



FIG. 7. Possible transitions from a modified macroscopic viewpoint for case F in the situations where the number of active nodes that are isolated from the GC in layer β is not negligible. Because some initially destroyed nodes in layer α are revived at stage t = 3 because of antagonistic interlinks, some nodes that are treated as active by encapsulation involuntarily belong to the GC at stage t = 3, which allows the impossible state transitions that are highlighted by stars. To compensate for this inconvenience, we modified the evaluation of the GC size as Eq. (47), highlighted by a small diagonal line. The meaning of each symbol is the same as in Fig. 1.

using the original interlayer message to derive the GC size at the microscopic level.

B. Heuristic

To keep our model analytically tractable, we define the provisional active variables $s_{i_{\alpha}}^{*t=3}$ instead of $s_{i_{\alpha}}^{t=3}(F)$ such that they do not include messages in layer α . We first describe Eq. (20) in detail and show that $s_{i_{\alpha}}^{t=3}(F)$ includes the past intralayer messages in layer α ,

$$s_{i_{\alpha}}^{t=3}(\mathbf{F}) = s_{i_{\alpha}}^{t=1} + \left(1 - s_{i_{\alpha}}^{t=1}\right) \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} - s_{i_{\alpha}}^{t=1} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}\right).$$
(43)

It is clear that $s_{i_{\alpha}}^{t=3}$ F is already influenced by the connectivity of layer α , because it includes the term $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1}$, which is indeed $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ itself in the case that node i_{α} has the history (I) (see Table I). Supposing the term $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} = 0$

TABLE I. All possible cases of $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1}$, $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$, or $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$, denoted by $M_{i_{\alpha}}^{t=1}$, $M_{i_{\alpha}}^{t=3}$, and $M_{i_{\alpha}}^{*t=3}$, respectively. Here ϕ denotes the unrealizable case in which $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1}$ vanishes and $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ is equal to unity at the same time.

$s_{i_{\alpha}}^{t=1}$	$M_{i_{\alpha}}^{t=1}$	$M_{i_{\alpha}}^{t=3}$	$M^{*t=3}_{i_{\alpha}}$	$s_{i_{\alpha}}^{t=1}M_{i_{\alpha}}^{t=1}M_{i_{\alpha}}^{t=3}$	$s_{i_{\alpha}}^{t=1}M_{i_{\alpha}}^{t=1}M_{i_{\alpha}}^{*t=3}$
1	1	1	1	1	1
1	1	0	0	0	0
1	0	0	0	0	0
1	0	1	1	ϕ	ϕ

in Eq. (43), we define the provisional active variables

$$s_{i_{\alpha}}^{*t=3}(\mathbf{F}) \equiv s_{i_{\alpha}}^{t=1} + \left(1 - s_{i_{\alpha}}^{t=1}\right) \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}.$$
 (44)

Substituting $s_{i_{\alpha}}^{*t=3}(F)$ for $s_{i_{\alpha}}^{t=3}(F)$ in Eqs. (21) and (22), we compute the provisional message $m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$. Employing $s_{i_{\alpha}}^{*t=3}$ and $m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$ in Eq. (23), we can compute the provisional GC,

$$\sigma_{i_{\alpha}}^{*t=3} \equiv s_{i_{\alpha}}^{*t=3}(\mathbf{F}) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3} \right), \tag{45}$$

the sum of which is larger than the actual GC size in particular when the number of isolated nodes in layer β is not negligible (see Fig. 7).

Our idea is to replace only intralayer messages in Eq. (23): Employing $m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$ instead of $m_{a_{\alpha} \to i_{\alpha}}^{t=3}$, we approximately describe the GC label of each node at stage t = 3 [Eq. (23)] in detail:

$$\sigma_{i_{\alpha}}^{t=3}(\mathbf{F}) \approx s_{i_{\alpha}}^{*t=3}(\mathbf{F}) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3} \right) - s_{i_{\alpha}}^{t=1} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \\ \times \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3} \right),$$

$$(46)$$

where it is possible to replace $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ with $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$, which is due to the correlations between the product of intralayer messages at stage t = 1 and that at stage t = 3 (see Table I). Therefore, we renew Eq. (46) as

$$\sigma_{i_{\alpha}}^{t=3}(\mathbf{F}) \approx \sigma_{i_{\alpha}}^{*t=3} - s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right) \times \left(\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3} \right), \quad (47)$$

where the second term of Eq. (47) corresponds to deducting the nodes that incorrectly belong to the GC.

C. Macroscopic evaluation and numerical validation

For the purposes of macroscopic analysis, we denote the ratio of active nodes at t = 3 based on Eq. (44),

$$q_{k_{\alpha},k_{\beta}}^{*\alpha(\mathrm{F}),t=3} = q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} + (1 - q_{k_{\alpha},k_{\beta}}^{\alpha,t=1}) (I_{k_{\alpha},k_{\beta}}^{\beta,t=2})^{k_{\beta}}.$$
 (48)

Substituting $q_{k_{\alpha},k_{\beta}}^{*\alpha(\mathrm{F}),t=3}$ in the self-consistent equation

$$I_{l_{\alpha},l_{\beta}}^{*\alpha,t=3} = \sum_{k_{\alpha},k_{\beta}} r_{\alpha}(k_{\alpha},k_{\beta}|l_{\alpha},l_{\beta})$$
$$\times \left[1 - q_{k_{\alpha},k_{\beta}}^{*\alpha(\mathrm{F}),t=3} + q_{k_{\alpha},k_{\beta}}^{*\alpha(\mathrm{F}),t=3} (I_{k_{\alpha},k_{\beta}}^{*\alpha,t=3})^{k_{\alpha}}\right], \quad (49)$$



FIG. 8. Modified analytical results of the robustness of layer α versus q for the case where antagonistic networks suffer from RFs and the setting is case F. Examples of the robustness of layer α at stage t = 1 are also plotted, which are the results of failure processes that are completed in the single layer α . Each dot is averaged 50 times, produced by numerical experiments.

we compute the set of messages $I_{l_{\alpha},l_{\beta}}^{*\alpha,t=3}$, which yields the fraction of the GC

$$\mu_{\alpha(\mathrm{F})}^{\star t=3} = \sum_{k_{\alpha},k_{\beta}} P(k_{\alpha},k_{\beta}) q_{k_{\alpha},k_{\beta}}^{*\alpha(\mathrm{F}),t=3} \left[1 - \left(I_{k_{\alpha},k_{\beta}}^{*\alpha,t=3} \right)^{k_{\alpha}} \right] - q_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \left[1 - \left(I_{k_{\alpha},k_{\beta}}^{\beta,t=2} \right)^{k_{\beta}} \right] \left[\left(I_{k_{\alpha},k_{\beta}}^{\alpha,t=1} \right)^{k_{\alpha}} - \left(I_{k_{\alpha},k_{\beta}}^{*\alpha,t=3} \right)^{k_{\alpha}} \right]$$
(50)

based on Eq. (47).

Figure 8 shows the size of the GCs in layer α predicted by Eq. (50) in the case where no degree-degree correlations exist, which is in excellent agreement with the experimental data. Similar accuracy was also achieved for the other parameter sets.

VI. RESULT

A. Influence of degree-degree correlations

Here we argue the influence of interlayer and intralayer degree-degree correlations on the robustness of each layer, the specific topologies of which are defined in Sec. IV A. We narrow the argument to the case F and TAs because the percolation processes in each layer in case Q can be reduced to those of a single network and the influence of degree-degree-correlations on the robustness in RFs is considerably smaller than that in TAs. In Figs. 9 and 10 we show examples of the robustness of layer α and the effects of various interlayer and intralayer degree-degree correlations.

We focus on the effect of degree-degree correlations on critical (minimum) robustness ($\mu_{\alpha} \approx 0$) and maximum robustness ($\mu_{\alpha} \approx 1$) of layer α , respectively. The thresholds at which μ_{α} vanishes depend on only its intralayer



FIG. 9. Plot of GC size in layer α versus the initial parameter q for the case where antagonistic networks suffer from TAs, the remaining effect of which is case F. The GC size in layer α at stage t = 1 is also shown, which is the result of the percolation process that is completed in the single layer α . The intralayer correlations in layer α are fixed to be positive, e.g., $C_{\alpha} = 0.6$, to focus on the effect of intralayer degree-degree correlations in layer β and interlayer degree-degree correlations on the robustness of layer α .





FIG. 10. Plot of GC size in layer α versus the initial parameter q for the case where antagonistic networks suffer from TAs, the remaining effect of which is case F. The GC size in layer α at stage t = 1 is also shown, which is the result of the percolation process that is completed in a single network. The intralayer correlations in layer α are fixed to be negative, e.g., $C_{\alpha} = -0.4$, to focus on the effect of intralayer degree-degree correlations in layer α .

degree-degree correlations, which are characterized by the Pearson coefficient C_{α} . This is because the GC at t = 1 plays the role of the core of the final GC and thus if there exists no GC at t = 1, no node can belong to the GC thereafter. Meanwhile, the maximum robustness of layer α is affected by both intralayer degree-degree correlations and interlayer degree-degree correlations, the transition point of which accords with that for critical robustness of layer β .

Let us consider the advantageous conditions for the maximum robustness of layer α , the parameter example of which is $C_{\alpha} = 0.6$, $C_{\beta} = -0.4$, and $C_I = -1$ in Fig. 9. Layer α is more robust against TAs if its degree-degree correlations are positive. In addition, layer α is more robust if layer β is more fragile due to antagonistic properties. Layer β is the most fragile if its intralayer degree-degree correlations are negative and it suffers from TAs, which are realized in the cases where interlayer degree-degree correlations are highly negative. In this case, nodes of higher degree in layer β become inactive because they connect with nodes of lower degree in layer α that tend to be active due to TAs in layer α . Considering the above, it is natural that the maximum robustness of layer α is the most fragile if the parameter example is $C_{\alpha} = -0.4, C_{\beta} = 0.6$, and $C_I = 1$ in Fig. 10, because each sign of the parameter is opposite that in the conditions that are advantageous for layer α .

B. Possible relevance to real-world systems

As for significance to real-world systems, the antagonistic networks may serve as a model of complex ecological interactions between endangered species and invaders. References [54,55] report such ecological relationship in the Amami islands in Japan: Populations of endangered species, such as the Amami rabbit (*Pentalagus furnessi*) and the Amami Ishikawa's frog (*Odorrana splendida*) were restored to almost the original level at the areas where invaders such as the small Indian mongoose (*Herpestes auropunctatus*) were exterminated, whereas few were observed in places where the invaders were established.

Our analysis shows that if initially damaged nodes can be reactivated (case F), the GC size is restored to the original level owing to the antagonistic inhibition to the replica nodes as long as the fraction of the initial damage is sufficiently small. This is consistent with the above reports. The endangered species in the reports have relatively high reproduction rates and short life cycles, which may fit the conditions of case F. On the other hand, species such as large mammals and primates have low reproduction rates and long life cycles and may correspond to case Q, for which populations of the species cannot be restored only by the antagonistic interaction. However, our analysis, in conjunction with Refs. [54,55], implies that, even in such cases, the combination of increasing the reactivation rate of endangered species by human-induced methods such as relocation and extermination of invasive species is an effective scheme for conserving ecological systems.

VII. SUMMARY

In this paper we developed an analytical methodology based on the cavity method to study the robustness of duplex networks coupled with antagonistic interlinks, considering intralayer and interlayer degree-degree correlations. We investigated two scenarios according to whether initially failed nodes are able to revive (case F) or not (case Q) with the aid of their replica nodes. In both cases Q and F we showed that the failure process periodically repeated because of the peculiarity of the antagonistic property of interlinks and the percolation transition exhibited continuity.

The oscillation was due to hysteresis of each layer, which led to the inconsistency between theory and experiment particularly for case F. Therefore, we introduced a heuristic treatment for improving the theoretical prediction accuracy, the utility of which was verified by numerical experiments.

We also argued the most advantageous situations for layer α in terms of degree-degree correlations employ bimodal networks. While the minimum robustness of layer α ($\mu_{\alpha} \approx 0$) was affected from only intralayer degree-degree correlations

in layer α , the maximum robustness of layer α ($\mu_{\alpha} \approx 1$) was influenced from various degree-degree correlations; the most robust situations are positive intralayer degree-degree correlations in layer α , negative intralayer degree-degree correlations in layer β , and negative interlayer degree-degree correlations. As for significance to real-world systems, a possible relevance to ecological systems that are composed of endangered and invasive species was mentioned.

Future works should include the construction of a model and an analytical framework that works with more realistic settings in the network approach.

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APPENDIX: RECONFIRMING THE PERIODICITY BY THE HEURISTIC

Although we already made sure that the percolation processes oscillate in both cases Q and F in Sec. II A, here we reconfirm these by the heuristic in Sec. V B.

1. Case Q

The active variable of each node at stage t = 3 is provided as

$$s_{i_{\alpha}}^{t=3}(\mathbf{Q}) = s_{i_{\alpha}}^{t=1} m_{p \to i_{\alpha}}^{t=3} = \left(s_{i_{\alpha}}^{t=1}\right)^{2} + s_{i_{\alpha}}^{t=1} \left(1 - s_{i_{\alpha}}^{t=1}\right) \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} - \left(s_{i_{\alpha}}^{t=1}\right)^{2} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}\right) = s_{i_{\alpha}}^{t=1} \left[1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}\right)\right].$$
(A1)

Substituting $m_{a_{\alpha} \to i_{\alpha}}^{t=1} = 0$ in Eq. (A1), we obtain $s_{i_{\alpha}}^{*t=3}(Q) = s_{i_{\alpha}}^{t=1}$, which also shows that $m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$ is equivalent to $m_{a_{\alpha} \to i_{\alpha}}^{t=1}$. Using $s_{i_{\alpha}}^{*t=3}(Q)$ and $m_{a_{\alpha} \to i_{\alpha}}^{*t=3}$, we obtain the result on the GC size, namely,

$$\sigma_{i_{\alpha}}^{t=3} \approx s_{i_{\alpha}}^{t=3}(\mathbb{Q}) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{*t=3} \right) = s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right)^{2} + \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right) \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}$$

$$= s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \right) = \sigma_{i_{\alpha}}^{t=1}.$$
(A2)

2. Case F

Substituting Eq. (43) into Eq. (24), we describe $s_{i_{\beta}}^{t=4}(F)$ in detail,

$$s_{i_{\beta}}^{t=4}(\mathbf{F}) = 1 - s_{i_{\alpha}}^{*t=3}(\mathbf{F}) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right) + s_{i_{\alpha}}^{t=1} \left(\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right) \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right).$$
(A3)

Note that $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ is replaced with $\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3}$ because of Table I. Supposing that $\prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2}$ vanishes in Eq. (A3), which also replaces $s_{i_{\alpha}}^{*t=3}$ with $s_{i_{\alpha}}^{t=1}$ because of Eq. (44), we define

$$s_{i_{\beta}}^{*t=4}(\mathbf{F}) \equiv 1 - s_{i_{\alpha}}^{t=1} \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right) + s_{i_{\alpha}}^{t=1} \left(\prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right) = 1 - s_{i_{\alpha}}^{t=1} + s_{i_{\alpha}}^{t=1} \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=1} = s_{i_{\beta}}^{t=2}, \quad (A4)$$

where the last equal sign is because of Eq. (13). Employing $s_{i_{\beta}}^{*t=4}(F)$ instead of $s_{i_{\beta}}^{t=2}$ in Eqs. (14) and (15), we derive the provisional message at t = 4 and $m_{a_{\beta} \to i_{\beta}}^{*t=4}$, which completely accords with $m_{a_{\beta} \to i_{\beta}}^{t=2}$ because of Eq. (A4).

The label of the GC at t = 4 is derived using the original interlayer message $s_{i_{\beta}}^{t=4}(F)$ and the provisional message $m_{a_{\beta} \to i_{\beta}}^{*t=4}$, namely,

$$\sigma_{i_{\beta}}^{t=4} \approx s_{i_{\beta}}^{t=4}(\mathbf{F}) \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{*t=4} \right) = s_{i_{\beta}}^{t=4}(\mathbf{F}) \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right) = s_{i_{\alpha}}^{t=2} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right) - \left(1 - s_{i_{\alpha}}^{t=1} \right) \left(1 - \prod_{a_{\alpha} \in \partial i_{\alpha}} m_{a_{\alpha} \to i_{\alpha}}^{t=3} \right) \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \left(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} \right) = \sigma_{i_{\beta}}^{t=2}, \tag{A5}$$

where the last equal sign is derived because $\prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2} (1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{t=2})$ always vanishes. We also note that the label of the provisional GC, which is defined as $\sigma_{i_{\beta}}^{*t=4} \equiv s_{i_{\beta}}^{*t=4}(F)(1 - \prod_{a_{\beta} \in \partial i_{\beta}} m_{a_{\beta} \to i_{\beta}}^{*t=4}) = \sigma_{i_{\beta}}^{t=2}$, also correctly evaluates the GC at t = 4, because there exist no nodes that incorrectly belong to the GC at t = 4 even if we set $s_{i_{\beta}}^{*t=4}(F)$ instead of $s_{i_{\beta}}^{t=4}(F)$.

- [1] P. Erdős and A. Réyni, Publ. Math. 6, 290 (1959).
- [2] B. Bollobás, S. Janson, and O. Riordan, Random Struct. Alg. 31, 3 (2007).
- [3] S. H. Strogatz, Nature (London) 410, 268 (2001).
- [4] D. J. Watts, Small Worlds: The Dynamics of Networks between Order and Randomness (Princeton University Press, Princeton, 1999).
- [5] R. Albert and A.-L. Barabási, Rev. Mod. Phys. 74, 47 (2002).
- [6] M. E. J. Newman, SIAM Rev. 45, 167 (2003).
- [7] S. N. Dorogovtsev and J. F. F. Mendes, Adv. Phys. 51, 1079 (2002).
- [8] S. Boccaletti, V. Latora, Y. Moreno, M. Chavez, and D.-U. Hwang, Phys. Rep. 424, 175 (2006).
- [9] M. Barthelémy, Phys. Rep. 499, 1 (2011).
- [10] D. S. Callaway, M. E. J. Newman, S. H. Strogatz, and D. J. Watts, Phys. Rev. Lett. 85, 5468 (2000).
- [11] S. Boccaletti, G. Bianconi, R. Criado, C. I. del Genio, J. Gómez-Gardenes, M. Romance, I. Sendina-Nadal, Z. Wang, and M. Zanin, Phys. Rep. 544, 1 (2014).
- [12] M. Kivelä, A. Arenas, M. Barthélemy, J. P. Gleeson, Y. Moreno, and M. A. Porter, J. Complex Netw. 2, 203 (2014).
- [13] G. D'Agostino and A. Scala, Networks of Networks: The Last Frontier of Complexity (Springer, Berlin, 2014).
- [14] J. Gao, S. V. Buldyrev, S. Havlin, and H. E. Stanley, Phys. Rev. Lett. 107, 195701 (2011).
- [15] R. G. Morris and M. Barthelemy, Phys. Lett. 109, 128703 (2012).
- [16] G. J. Baxter, S. N. Dorogovtsev, A. V. Goltsev, and J. F. F. Mendes, Phys. Lett. **109**, 248701 (2012).
- [17] S. V. Buldyrev, N. W. Shere, and G. A. Cwilich, Phys. Rev. E 83, 016112 (2011).
- [18] X. Huang, J. Gao, S. V. Buldyrev, S. Havlin, and H. E. Stanley, Phys. Rev. E 83, 065101 (2011).
- [19] Y. Hu, B. Ksherim, R. Cohen, and S. Havlin, Phys. Rev. E 84, 066116 (2011).
- [20] D. Zhou, J. Gao, H. E. Stanley, and S. Havlin, Phys. Rev. E 87, 052812 (2013).
- [21] J. Shao, S. V. Buldyrev, S. Havlin, and H. E. Stanley, Phys. Rev. E 83, 036116 (2011).

- [22] R. Parshani, C. Rozenblat, D. Ietri, C. Ducruet, and S. Havlin, Europhys. Lett. 92, 68002 (2010).
- [23] Z. Wang, A. Szolnoki, and M. Perc, Europhys. Lett. 97, 48001 (2012).
- [24] B. Podobnik, D. Horvatić, M. Dickison, and H. E. Stanley, Europhys. Lett. 100, 50004 (2012).
- [25] X. Huang, S. Shao, H. Wang, S. V. Buldyrev, H. E. Stanley, and S. Havlin, Europhys. Lett. 101, 18002 (2013).
- [26] G. Dong, L. Tian, R. Du, J. Gao, H. E. Stanley, and S. Havlin, Europhys. Lett. **102**, 68004 (2013).
- [27] M. M. Danziger, A. Bashan, Y. Berezin, and S. Havlin, J. Complex Netw. 2, 460 (2014).
- [28] D. Cellai, E. López, J. Zhou, J. P. Gleeson, and G. Bianconi, Phys. Rev. E 88, 052811 (2013).
- [29] S. V. Buldyrev, R. Parshani, G. Paul, H. E. Stanley, and S. Havlin, Nature (London) 464, 1025 (2010).
- [30] R. Parshani, S. V. Buldyrev, and S. Havlin, Phys. Rev. Lett. 105, 048701 (2010).
- [31] S. M. Rinaldi, J. P. Peerenboom, and T. K. Kelly, IEEE Control Syst. 21, 11 (2001).
- [32] K. Zhao and G. Bianconi, J. Stat. Mech. (2013) P05005.
- [33] K. Zhao and G. Bianconi, J. Stat. Phys. 152, 1069 (2013).
- [34] B. Kotnis and J. Kuri, Phys. Rev. E 91, 032805 (2015).
- [35] G. P. Boswell, N. F. Britton, and N. R. Franks, Proc. R. Soc. London B 265, 1921 (1998).
- [36] A. Fall, M.-J. Fortin, M. Manseau, and D. O'Brien, Ecosystems 10, 448 (2007).
- [37] D. L. Urban, E. S. Minor, E. A. Treml, and R. S. Schick, Ecol. Lett. 12, 260 (2009).
- [38] J. M. Tylianakis, E. Laliberté, A. Nielsen, and J. Bascompte, Biol. Cons. 143, 2270 (2010).
- [39] G. Huth, A. Lesne, F. Munoz, and E. Pitard, Physica A 416, 290 (2014).
- [40] M. E. J. Newman, Phys. Rev. Lett. 89, 208701 (2002).
- [41] S. Maslov and K. Sneppen, Science 296, 910 (2002).
- [42] M. E. J. Newman, Phys. Rev. E 67, 026126 (2003).
- [43] B. Min, S. D. Yi, K.-M. Lee, and K.-I. Goh, Phys. Rev. E 89, 042811 (2014).

- [44] D. Zhou, H. E. Stanley, G. D'Agostino, and A. Scala, Phys. Rev. E 86, 066103 (2012).
- [45] J. Pearl, Probabilistic Reasoning in Intelligent Systems: Networks of Plausible Inference (Kaufmann, San Francisco, 1988).
- [46] A. K. Hartmann and M. Weigt, *Phase Transitions in Combinatorial Optimization Problems: Basics, Algorithms and Statistical Mechanics* (Wiley-VCH, Berlin, 2005).
- [47] M. Mézard and Montanari, *Information, Physics, and Computation* (Oxford University Press, Oxford, 2009).
- [48] S. Watanabe and Y. Kabashima, Phys. Rev. E 89, 012808 (2014).
- [49] S. W. Son, G. Bizhani, C. Christensen, P. Grassberger, and M. Paczuski, Europhys. Lett. 97, 68002 (2012).

- [50] M. E. J. Newman, S. H. Strogatz, and D. J. Watts, Phys. Rev. E 64, 026118 (2001).
- [51] The fraction of nodes that are incorrectly evaluated as the constituents of the GC at t = 3 in case F, which are marked with a star in Fig. 7, is evaluated to be 0.02 at most in our study for duplex networks, each layer of which is a regular random network in which the degree of the node is 4.
- [52] J. Hoshen and R. Kopelman, Phys. Rev. B 14, 3438 (1976).
- [53] A. Al-Futaisi and T. W. Patzek, Physica A 321, 665 (2003).
- [54] Y. Watari, S. Nishijima, M. Fukasawa, F. Yamada, S. Abe, and T. Myashita, Ecol. Evol. 3, 4711 (2013).
- [55] K. Fukasawa, T. Miyashita, T. Hashimoto, M. Tatara, and S. Abe, Proc. R. Soc. London B 280, 20132075 (2013).