# Cascading failures in ac electricity grids

Martin Rohden,<sup>1,\*</sup> Daniel Jung,<sup>1,†</sup> Samyak Tamrakar,<sup>1,2,‡</sup> and Stefan Kettemann<sup>1,3,§</sup>

<sup>1</sup>Jacobs University, Department of Physics and Earth Sciences, Campus Ring 1, 28759 Bremen, Germany

<sup>2</sup>Physics Department, Carl von Ossietzky Universität Oldenburg, Ammerländer Heerstraße 114-118, 26129 Oldenburg

<sup>3</sup>Pohang University of Science and Technology, Division of Advanced Materials Science, San 31, Hyoja-dong, Nam-gu,

Pohang 790-784, South Korea

(Received 30 March 2016; published 9 September 2016)

Sudden failure of a single transmission element in a power grid can induce a domino effect of cascading failures, which can lead to the isolation of a large number of consumers or even to the failure of the entire grid. Here we present results of the simulation of cascading failures in power grids, using an alternating current (AC) model. We first apply this model to a regular square grid topology. For a random placement of consumers and generators on the grid, the probability to find more than a certain number of unsupplied consumers decays as a power law and obeys a scaling law with respect to system size. Varying the transmitted power threshold above which a transmission line fails does not seem to change the power-law exponent  $q \approx 1.6$ . Furthermore, we study the influence of the placement of generators and consumers are especially vulnerable to cascading failures. As a real-world topology, we consider the German high-voltage transmission grid. Applying the dynamic AC model and considering a random placement of consumers, we find that the probability to disconnect more than a certain number of consumers depends strongly on the threshold. For large thresholds the decay is clearly exponential, while for small ones the decay is slow, indicating a power-law decay.

DOI: 10.1103/PhysRevE.94.032209

## I. INTRODUCTION

A reliable supply of electric power is of fundamental importance for the technical infrastructure of modern societies. In fact, the reliability of electric power grids has increased continuously in the past few decades [1]. However, large-scale power outages still occur and can affect millions of customers, which may result in catastrophic events. It is therefore important to understand which topological properties of power grids and which placements of generators and consumers on the grid are able to diminish the risk of large-scale outages.

Large-scale outages can often be traced back to the failure of a single transmission element [2–5]. The initial failure causes secondary failures, which can eventually lead to a whole cascade of failures. Cascading failures have been analyzed in various studies with different models and from different viewpoints [6–18]. Most of these previous studies analyze the influence of network topology on the cascade of failures using simplified topological flow models such as the *messenger model* introduced by Motter and Lai [6], which is used to study cascading failures in Refs. [19,20].

In this work we base the analysis of cascading failures on the *alternating current (AC) power flow equations* [21–24]. This allows us to study the influence of the physical properties of the grid, such as the placement of generators and consumers and the power capacitance of the transmission lines, on the probability and the extent of cascading failures. To this end, we consider a random placement of generators and consumers on a regular 2D grid graph (square grid) as well as on a

model topology for the German high-voltage transmission grid (380 kV) [25].

An evaluation of the statistics of power failures which occurred in real power grids over the past few decades shows that the probability of ending up with more than a certain number of unsupplied consumers often decays like a power law with the number of unsupplied consumers  $N_c$  [14,15]. Such a slow decay indicates a significant probability that a large number of consumers will become unsupplied. It is therefore of practical relevance to understand which properties of the grid are responsible for this behavior. In Ref. [14], a simple model has been suggested to simulate cascading failures that assumes that the load  $F_{ij}$  of a failing transmission line is redistributed in equal parts among the remaining transmission lines. This model can be solved analytically and yields a power-law distribution when the initial average power flow through the transmission lines prior to the failure  $F = \langle F_{ii} \rangle$ reaches a certain critical ratio of the threshold power  $F_{\rm th}$ . Below that value, an exponentially fast decay is found [14].

While most previous studies assume that the power-law dependence is only related to the network topology, in particular to scale-free topologies [26,27], the main goal of this work is to analyze which influence different placements of consumers and generators have on the probability to find  $N_c$  unsupplied consumers. We start by finding the stationary power flow in the fully connected grid. Then we initiate a power line failure by removing one transmission line by hand and find the new stationary power flow. The resulting redistribution of the power flow may trigger further line failures where the transmitted power exceeds a threshold  $F_{\text{th}}$ , which we set to be a certain ratio of the transmission line power capacitance. This chain of outages continues until we cannot find a stable solution for the stationary power flow anymore. We then record the number of consumers  $N_c$  that cannot be supplied anymore by the available generators. Note that

<sup>\*</sup>m.rohden@jacobs-university.de

<sup>&</sup>lt;sup>†</sup>d.jung@jacobs-university.de

<sup>&</sup>lt;sup>‡</sup>samyak.r.tamrakar@gmail.com

<sup>§</sup>s.kettemann@jacobs-university.de

consumers which are connected within clusters might also not be supplied if the number of generators does not at least equal the number of consumers in the specific cluster. This process is repeated by subsequently initiating the power outage with every transmission line of the grid, removing that line, and observing the resulting cascade. Thereby, we obtain the probability distribution of the *blackout size*  $N_c$  in dependence of the placement of consumers and generators and the threshold  $F_{th}$ .

We gain further insights by analyzing the share of links that can induce a cascade of failures depending on the existence of various fixed cluster sizes. We demonstrate that the existence of large clusters of generators and consumers makes the grid particularly vulnerable to cascading failures, since the likelihood for a whole cluster to break down at once appears to increase with increasing cluster size. We thereby show that the size of the power outage depends essentially on the placement of consumers and generators. Finally, we study a real-world topology, a model for the German high-voltage transmission grid [25]. In this irregular grid structure, we find a decay of  $\bar{p}(N_c)$  for large  $N_c$ , which depends strongly on the threshold  $F_{\text{th}}$ . For large  $F_{\text{th}}$  the decay is clearly exponential. For small  $F_{\rm th}$  the decay is slow and may indicate a power-law decay. Thus, there might be a critical value of the threshold  $F_{\text{th}}$  in the German grid below which the cumulative probability density becomes critical and decays with a power law.

## **II. POWER GRID MODEL**

We approximate the power grid as a network of *N* rotating synchronous machines, representing generators and motors. Each machine  $k \in \{1, ..., N\}$  is characterized by the net mechanical power  $P_k^{\text{mech}}$ , which is positive for a generator and negative for a consumer. The state of machine *k* is given by the angular frequency and the rotor angle (power angle)  $\theta_k(t)$ which is measured relatively to a reference machine rotating at the nominal grid angular frequency  $\omega_0 = 2\pi \times 50$  Hz. Correspondingly,  $\omega_k(t) = d\theta_k(t)/dt = \dot{\theta}_k$  gives the angular frequency deviation from the reference frequency  $\omega_0$ . The dynamics of the rotors are governed by the *swing equation* [21–23],

$$I_k \frac{\mathrm{d}^2 \theta_k}{\mathrm{d}t^2} + D_k \frac{\mathrm{d}\theta_k}{\mathrm{d}t} = P_k^{\mathrm{mech}} - P_k^{\mathrm{el}},\tag{1}$$

where  $P_k^{\text{el}}$  is the net electrical power transmitted from adjacent rotating machines through the transmission lines.  $I_k$  is the moment of inertia of the rotor times  $\omega_0$  and  $D_k$  measures the damping, which is mainly due to damper windings [28].

For simplicity, we neglect ohmic losses of transmission lines which can be considered small in high voltage levels [29]. Thus, the line admittance is purely inductive,  $Y_{k\ell} = 1/(i\omega L_{k\ell})$ , where  $L_{k\ell}$  is the inductance of the line  $(k, \ell)$ . Then, the magnitude of the voltage is constant throughout the grid,  $|U_k| = U_0 \forall k \in \{1, ..., N\}$ . For a common two-pole synchronous machine, the phase of the voltage equals the mechanical phase of the rotor. The expression for the active electric power then simplifies to

$$P_k^{\rm el} = \sum_{\ell=1}^N \frac{U_0^2}{\omega L_{k\ell}} \sin(\theta_k - \theta_\ell).$$
(2)

Substituting this result into the swing equation (1) yields the equations of motion,

$$I_k \frac{\mathrm{d}^2 \theta_k}{\mathrm{d}t^2} + D_k \frac{\mathrm{d}\theta_k}{\mathrm{d}t} = P_k^{\mathrm{mech}} - \sum_{\ell=1}^N \frac{U_0^2}{\omega L_{k\ell}} \sin(\theta_k - \theta_\ell). \quad (3)$$

Using the abbreviations

$$P_k = \frac{P_k^{\text{mech}} - D_k \omega_0}{I_k}, \quad \alpha_k = \frac{D_k}{I_k}$$
$$K_{k\ell} = \frac{U_0^2}{I_k \omega L_{k\ell}},$$

the oscillator model reads

$$\frac{\mathrm{d}^2\theta_k}{\mathrm{d}t^2} = P_k - \alpha_k \frac{\mathrm{d}\theta_k}{\mathrm{d}t} + \sum_{\ell=1}^N K_{k\ell} \sin(\theta_\ell - \theta_k). \tag{4}$$

These equations of motion are widely used to model power grids in power engineering, where it goes by the name *synchronous motor model* [30]. Notably, this model is similar to the *Kuramoto model*, which is studied extensively in statistical physics [31–33]. In the Kuramoto model, the inertia term is absent, so it can be seen as the overdamped limiting case.

In order to find the stationary solutions of the oscillator model (4), it is equivalent to solve the static flow equations,

$$0 = P_k + \sum_{\ell=1}^{N} K_{k\ell} \sin(\theta_\ell - \theta_k).$$
 (5)

This can be done by a standard root-finding algorithm [34]. Note that regular square grids consisting of loops with four edges, which are studied in the following, may have multiple stationary (stable and unstable) solutions [38]. We ensure the stability of a stationary state by checking that the Jacobian matrix has no positive eigenvalues.

## **III. CASCADING FAILURE ALGORITHM**

In the following we describe the cascading failure algorithm used in this study. We determine the stable state with power flows  $F_{ij}$ . Note that we skip realizations for which no stable state exists or the initial maximal power flow  $\max(F_{ij})$  is already larger than the threshold power flow  $F_{\rm th}$ . Thereby it is ensured that the initial power flow through all lines is stable. Next, we remove one of the transmission lines of the network in order to induce a cascade of failures. It is again ensured that the network reaches a new stable state with a new power flow distribution  $F'_{ij}$ . All transmission lines for which the transferred power  $F'_{ij}$  exceeds  $F_{\text{th}}$  are removed from the grid, and the power flows are recalculated. This process is repeated until no transferred power exceeds  $F_{\rm th}$  or until the grid splits into different subgrids. In the latter case, we record the number of affected consumers, the blackout size  $N_c$ , which is the number of consumers that cannot be supplied by generators anymore. The whole process is repeated for each transmission line of the original grid being initially removed and for other placements of generators and consumers  $P_k$ . Figure 1 illustrates the cascading failure algorithm for the example of a  $6 \times 6$  square grid. Figure 1(a) illustrates the initial stable state before the initial line is removed, Fig. 1(b) the



FIG. 1. Example for a cascading failure in a  $6 \times 6$  square grid with open boundary conditions. Triangles denote generators, and disks denote consumers. Intact transmission lines are indicated by a thin dashed line. The stable state power flow is plotted in color scale ranging from 0 to 5 s<sup>-2</sup>. (a) Power flow in the fully connected grid. (b) Power flow after removing one line. (c) Power flow after removing all lines with a power flow larger than the power threshold  $F_{\rm th}$  after initial removal of a transmission line. (d) After several more steps (not shown) some nodes become isolated (indicated by crosses), so the simulation is stopped.

stable state after the removal of one transmission line (upper left side of the grid), Fig. 1(c) the second step of the cascade of failures, and Fig. 1(d) the final step with seven disconnected consumers.

#### A. Statistical analysis

We initiate a cascade of failures by manually removing each of the  $N_L$  transmission lines of the original grid separately. For a  $L \times L$  square grid graph with open boundary conditions, there exist  $N_L = 2L(L - 1)$  links, so we perform the cascading failure algorithm  $N_L$  times. We repeat this for R = 1000 realizations of random placements of generators and consumers  $P_k$ . For each realization r, we obtain the histogram  $E_r(N_c)$ , which counts the number of events that  $N_c$ consumers are unsupplied. From this, we obtain the normalized probability distribution function (pdf)

$$e_r(N_c) = \frac{E_r(N_c)}{N_L},\tag{6}$$

the share of initially removed transmission lines for which the cascade resulted in  $N_c$  isolated consumers.

Then we compute the *complementary cumulative distribution function* (ccdf)  $p_r(N_c)$  (in short *cumulative probability* in the following), yielding the probability that the number of unsupplied consumers (the *blackout size*) is larger than  $N_c$ ,

$$p_r(N_c) = \sum_{N'_c = N_c + 1}^{\infty} e_r(N'_c).$$
(7)

Finally, we obtain the ensemble average  $\bar{p}(N_c)$  over R = 1000 realizations,

$$\bar{p}(N_c) = \frac{1}{R} \sum_{r=1}^{R} p_r(N_c).$$
(8)

## B. Regular square grid topology with random placement of consumers

In the following, we present the results for the statistics of the number of unsupplied consumers for a  $L \times L$  square grid with open boundary conditions for different linear system sizes L. We consider a simple regular square grid topology with open boundary conditions this rather simple model to systematically study the influence of different consumer placements. Half of the nodes serve as consumers and the other half as generators. Each node k generates the net power  $P_k = \pm P$  (positive for generators, negative for consumers), with  $P = 1 \text{ s}^{-2}$ . The power capacity of all transmission lines is set to  $K_{ij} = K = 5 \text{ s}^{-2}$ . All machines have the same damping parameter  $\alpha_k = \alpha = 1 \text{ s}^{-1}$ .

In order to precisely control the amount of randomness in the system, we use the following procedure to generate a random array  $P_k$  [39,40]: We start from a periodic arrangement of generators and consumers [41] and divide the graph into two subgraphs, one carrying all N/2 generators and the other all N/2 consumers. Then, p different nodes are chosen randomly from each subgraph, forming p generator-consumer pairs. Finally, each of these generator-consumer pairs is swapped. By generating a permutation of the periodic arrangement in this way, it is ensured that no node is swapped twice. The maximally disordered state is reached after  $p_{max} = N/4$ swaps, which is the case used throughout this study. There is a finite number of possible realizations, given by the *ensemble size*,

$$N_E = \binom{N/2}{p}.$$
(9)

In this study we always consider only a small subset of possible realizations, as  $N_E$  is a very large number already for the smallest considered systems.

The cumulative probability  $\bar{p}(N_c)$  is illustrated in Fig. 2 for various power flow thresholds  $f_{th} = F_{th}/K$  and linear system sizes L. For all considered threshold values  $f_{th}$ , the length of the tail of the cumulative probability  $\bar{p}(N_c)$  is increasing with increasing linear system size L. Larger systems possess more consumers, so the probability to obtain a large number of unsupplied consumers is increasing and the values of  $\bar{p}(N_c)$ are increasing with system size in the tail of the distribution.

The value  $\bar{p}(0)$  marks the probability that a cascade results in one or more unsupplied consumers. This probability is increasing for decreasing system size and decreases with the threshold value  $f_{\text{th}}$ . This is clearly seen for low critical values of  $f_{\text{th}} = 0.25$  and  $f_{\text{th}} = 0.3$  in Figs. 2(a) and 2(b), respectively. For  $f_{\text{th}} = 0.4$  and  $f_{\text{th}} = 0.5$ , illustrated in Fig. 2(c) and 2(d),



FIG. 2. Average cumulative probability  $\bar{p}(N_c)$  [cf. Eq. (8)] for square grids of different linear system size *L* and threshold  $f_{th} = F_{th}/K$ : (a)  $f_{th} = 0.25$ , (b)  $f_{th} = 0.3$ , (c)  $f_{th} = 0.4$ , (d)  $f_{th} = 0.5$ . Error bars show the standard error [cf. Eq. (8)].

these trends are less clearly visible since there are fewer observations available. Consequently, the relative standard deviation between different realizations is also increasing with increasing threshold values.

The results for the cumulative probabilities  $\bar{p}(N_c)$  in dependence of the threshold power flow  $f_{\rm th}$  for fixed system size are illustrated in Fig. 3. We observe an increase of  $\bar{p}(N_c)$  with decreasing  $f_{\rm th}$ . The length of the tails are almost independent of the threshold value  $f_{\rm th}$ , indicating that even for the largest threshold value  $f_{\rm th} = 0.5$  outages with a large number of unsupplied consumers occur.

For comparison, we also identify the initially existing consumer clusters that exist prior to the initialization of the cascade of failures. The size of an initially present consumer



FIG. 3. Average cumulative probability  $\bar{p}(N_c)$  (8) for square grids with varying thresholds  $f_{th}$  and fixed linear system size L (colored lines, red axis). The distribution of clusters of consumers  $\bar{p}(N_g)$  is shown as the black line (black axis). (a) L = 10, (b) L = 14. Error bars show the standard error [cf. Eq. (8)].

cluster is denoted by  $N_g$ . The corresponding complementary cumulative probability distribution  $\bar{p}(N_g)$  [cf. Eqs. (6)–(8)] represents the probability that the initial consumer clusters are larger than  $N_g$ , as illustrated in Fig. 3 (black curves). Interestingly, we do not find a direct relation between the consumer cluster size of the original grid  $N_g$  and the number of unsupplied consumers  $N_c$  after the cascade. The largest initial consumer clusters  $N_g$  are considerably larger than the largest number of unsupplied consumers after the cascade. For example, for the considered realizations with L = 10, the largest initial consumer cluster consists of 42 consumers, whereas the largest number of unsupplied consumers after the cascade is 26. We conjecture that in case of the square grid topology, the distribution of originally existing consumer clusters  $N_g$  provides an upper limit for the distribution of unsupplied consumers  $N_c$ , but we find no direct relation between  $N_g$  and  $N_c$  if the consumers are randomly distributed.

#### C. Scaling analysis

It has been noted in Ref. [14] that cascading failures may show evidence for *self-organized criticality* (SOC) [42]. SOC has been first found experimentally in rice piles, where the distribution of the size of avalanches has been found to follow a power law [43]. Numerical studies of avalanches in sand pile models have shown evidence for power-law behavior as well, but the accuracy of the numerical determination of the avalanche exponents q is limited by finite-size effects, resulting in values of the order of q = 1.6 [42,44,45]. The most effective way to evaluate the numerical data is by conducting a finitesize scaling analysis. Here we apply this strategy to evaluate our numerical data on cascading failures, using the scaling ansatz

$$\bar{p}(N_c, L) = N_c^{-q} f(N_c/L^d),$$
(10)

with some unknown scaling function  $f(N_c/L^d)$  and the effective scaling dimension d.

We rescale our data according to the scaling ansatz (10)with parameters d = 0.8 and q = 1.6 illustrated in Fig. 4. For the scaling ansatz to be valid, the curves of all system sizes Lshould coincide. The agreement of the data with the scaling ansatz is best for small threshold values  $f_{\text{th}}$ . We cannot exclude that the worse agreement for larger  $f_{\rm th}$  is only due to the decreasing statistics, since fewer cascades are initiated. Note that the effective scaling dimension d = 0.8 is smaller than the spatial dimension. We find good agreement with the scaling ansatz for values in the range  $q = 1.6 \pm 0.2$ . For large  $f_{\text{th}}$  we still find the best agreement with the scaling ansatz for the same parameter values as for  $f_{\rm th} = 0.25$ . We thus find evidence for scaling of the average cumulative probability  $\bar{p}(N_c,L)$  with an exponent  $q \approx 1.6 \pm 0.2$ , which possibly does not depend on the threshold value. A power-law dependence corresponds to a horizontal line of  $N_c^q \bar{p}(N_c, L)$  in Fig. 4. We only find a power-law dependence for a small interval.

In contrast to the simple model of Ref. [14] where the load  $F_{ij}$  of a failing transmission line is redistributed in equal parts among the remaining transmission lines, we do not find evidence for a critical ratio of the threshold power  $F_{th}$  within the studied range of  $F_{th}$  values. Interestingly, by including the physical flow model (5) in the simulation, although we only



FIG. 4. Scaling law for d = 0.8 and q = 1.6. Average cumulative probability  $\bar{p}(N_c, L)$  (8) for square grids of different linear system size *L* and with different threshold power flow  $f_{\text{th}}$ : (a)  $f_{\text{th}} = 0.25$ , (b)  $f_{\text{th}} = 0.3$ , (c)  $f_{\text{th}} = 0.35$ , (d)  $f_{\text{th}} = 0.4$ . Error bars show the standard error [cf. Eq. (8)].

find a small interval with a pure power-law decay, the exponent  $q \approx 1.6 \pm 0.2$  is comparable to the one found in Ref. [14],  $q \approx 1.4$ .

## IV. IMPACT OF CLUSTERING IN THE REGULAR SQUARE GRID TOPOLOGY

To better understand the impact of the initial clustering of generators and consumers on the cascading failures, we analyze regular square grids (L = 12) with open boundary conditions and a periodic arrangement of generator and consumer clusters of fixed size  $(1 \times 2, 2 \times 2, 2 \times 4, \text{ etc.})$ . These may occur as particular cases of the random realizations, analyzed before. For these we are able to demonstrate a direct relation between the initial consumer cluster size  $N_g$  and the number of unsupplied consumers  $N_c$  after the cascade. Obviously, the largest possible clusters exist if all N/2 consumers are located on one side of the grid and all N/2 generators on the other side (6  $\times$  12 clusters). The smallest possible clusters are formed by pairs of connected consumers or generators  $(1 \times 2 \text{ clusters})$ . We also consider clusters with 4 (2  $\times$  2), 8 (2  $\times$  4), 12 (3  $\times$  4), 24 (4  $\times$  6), and 36  $(6 \times 6)$  consumers or generators per cluster. In order to obtain comparable results, we determine the maximal power flow in the initial grid  $F_{\text{max}}$  and set the threshold power flow to  $F_{\rm th} = F_{\rm max} + 0.1 \ {\rm s}^{-2}$ .

We measure two quantities and study their dependence on the cluster size  $N_g$ : The maximal number of unsupplied consumers that can occur, and the most likely number of unsupplied consumers  $N_c$  (cf. Fig. 5). The results demonstrate that the most likely value of  $N_c$  increases linearly with the cluster size  $N_g$  [cf. Fig. 5(a), red curve]. It is therefore the most probable event that exactly one consumer cluster disconnects from the grid. For small cluster sizes, the maximum outage



FIG. 5. (a) Blue: Maximal number of unsupplied consumers  $N_c$  that occurs for all possible initial line failures as function of cluster size  $N_g$ . Red: Most likely number of unsupplied consumers. (b) Blue: Number of supercritical links whose initial failure causes the maximal possible number of unsupplied consumers. Red: Number of critical links whose initial failure cause a power outage.

that can occur [cf. Fig. 5(a), blue curve] is slightly above the most probable outage, indicating that sometimes more than one cluster contains unsupplied consumers at the end of the cascade. For large cluster sizes (24 to 72 nodes), an outage of one cluster is at the same time the maximally possible event and the most probable one.

The number of critical links, i.e., those links that cause a power outage, is found to increase with the initial cluster size  $N_g$  up to a cluster size of eight [cf. Fig. 5(b), red curve]. For larger cluster sizes, this number of critical links is decreasing. Links that connect different clusters are found to be particularly vulnerable. The number of such links is decreasing with increasing cluster size, explaining the decrease in the number of critical links.

Next, we consider those links that cause the maximal observed outage, which we call *supercritical links*. The number of such supercritical links is first decreasing with cluster size and than increasing again [cf. Fig. 5(b), blue curve], until the number of supercritical links eventually matches the number of critical links in the limit of large cluster size. Only a few links cause a maximum outage for intermediate cluster sizes, whereas for large cluster sizes almost every link causes a maximal outage. We conclude that large clusters are favorable in the sense that only few transmission lines can cause an outage. On the other hand, if such an outage occurs, the impact is much more severe, i.e., more consumers are affected.

#### A. German transmission grid

So far, all results have been obtained for the regular square grid topology. To test our findings on a realistic grid topology, we consider a model for the German high-voltage transmission grid. The model grid is based on data from the SciGRID project [25], where only the 380-kV level is considered. The objective is to test our algorithm on a more realistic,



FIG. 6. Example for a cascade of line failures in the German grid model, with transmission capacity  $K = 10 \text{ s}^{-2}$  and threshold  $f_{\text{th}} = 0.35$ . Triangles denote generators, and disks denote consumers. (a) Initial power flow. (b) The cascade is initiated by removing the link marked with a red circle. (c) Another link fails. (d) Another four links fail, and the simulation is stopped because the grid is not fully connected anymore. We mark the disconnected component containing the unsupplied nodes and other failed lines in red.

nonregular topology, for which it is sufficient to consider the topology of the German grid, isolated from the rest of the pan-European electricity network. As before, we consider a binary distribution of generators and consumers and a constant line power transmission capacity. We also apply the cascading failure algorithm described above and consider R = 1000 realizations. Figure 6 demonstrates a cascade of single-line failures in the German grid model for a binary distribution of generators and consumers. Here, we use the parameters  $P = 1 \text{ s}^{-2}$ ,  $K = 10 \text{ s}^{-2}$ ,  $\alpha = 1 \text{ s}^{-1}$ , and  $f_{\text{th}} = 0.35$ .

We analyze cascading failures for various threshold values  $f_{\text{th}}$ . The results for the cumulative probabilities  $\bar{p}(N_c)$  are shown in Fig. 7. With decreasing  $f_{\text{th}}$ , the values of the average cumulative probability  $\bar{p}(N_c)$  are increasing. This behavior is



FIG. 7. Average cumulative probability  $\bar{p}(N_c)$  (8) as function of the number of unsupplied consumers  $N_c$  for the German grid model for various threshold values  $f_{\text{th}}$  (colored lines, red axis). The distribution of initial consumer clusters  $\bar{p}(N_g)$  is shown as the black line (black axis). Error bars show the standard error [cf. Eq. (8)].

similar to that of the square grids (cf. Fig. 2). Note that the German grid model contains 254 nodes and 317 links, so its size is comparable to a  $12 \times 12$  square grid, which has 264 links, so we can compare the results with the corresponding probability density  $\bar{p}(N_c)$  shown in Fig. 2. We observe clear differences: First, for the German grid  $\bar{p}(N_c)$  does not decay continuously but shows abrupt steps. Second, the decay of  $\bar{p}(N_c)$  for large  $N_c$  depends strongly on the threshold  $f_{\text{th}}$ . For large  $f_{\rm th}$  the decay is clearly exponential. For  $f_{\rm th} = 0.4$  the decay is slow and may indicate a power-law decay, although the statistics does not allow to extract a definite power exponent. There might be a critical value of the threshold  $f_{\rm th}$  below which the cumulative probability density becomes critical and decays with a power law. Thus, we may conclude that the difference in the topological structure of the German high-voltage grid as compared to the square grids leads to marked differences in the probability of cascading failures, whose origin needs to be studied in more detail. We also show in Fig. 2 the distribution of consumer clusters  $\bar{p}(N_g)$  (black line) existing in the grid before the cascade. We observe that for the smallest threshold value,  $f_{\rm th} = 0.4$ , the tail of the cumulative probability distribution of unsupplied consumers  $N_c$  is similar to the tail of the distribution of clusters, indicating that the probability of large outage sizes is closely related to the occurrence of large clusters in the German grid topology.

## V. CONCLUSIONS AND OUTLOOK

Single line failures can induce a cascade of failures leading to power outages for a potentially large number of consumers. Although cascading failures have been studied before, most of the studies were based on simple flow models. In this article we contribute to a more realistic model for cascading failures based on an AC load flow model. We developed an algorithm for cascading failures which is widely applicable to different network topologies and parameter ranges.

First, we analyzed regular square grids with a random placement of generators and consumers. We considered different threshold values for the transmitted power  $F_{\text{th}}$ . We identified a scaling law with a power-law exponent for the probability to obtain a minimum number of unsupplied consumers (blackout size) in dependence of the system size. In contrast to the simplified model of Ref. [14], we do not find evidence for a critical ratio of the threshold power  $F_{\rm th}$  but rather find the same dependence in a wide range of threshold values  $f_{\rm th}$ . Moreover, we find a power-law decay with power  $q = 1.6 \pm 0.2$  which is comparable to the one found in Ref. [14],  $q \approx 1.4$ . We do not find a direct relation between the distribution of initial clusters and the resulting blackout sizes after the cascading failures. Often, clusters of consumers are found to split into multiple parts during a cascade of line failures, which might explain why the observed blackout size is often smaller than the initial cluster size.

Second, we studied regular square grids with a periodic arrangement of consumer clusters of fixed size. We demonstrated that large consumer clusters often lead to large outages, whereas small clusters typically lead to only small outages for the initial failure of a single line. In contrast, the number of critical links, i.e., links which cause a cascade if they fail, is decreasing with increasing cluster size. Here we can also identify a direct relation between consumer cluster sizes and the number of unsupplied consumers after a cascade has been initiated.

Finally, we studied a real-world topology, a model for the German high-voltage transmission grid. In this irregular grid structure, we find a qualitatively similar behavior for the complementary cumulative probability  $\bar{p}(N_c)$  that a blackout of a certain minimum size occurs, but we also observe clear differences: First, for the German grid,  $\bar{p}(N_c)$  does not decay continuously but shows abrupt steps, which could be only due to the small number of observations. Second, the decay of  $\bar{p}(N_c)$  for large  $N_c$  depends strongly on the threshold  $f_{\text{th}}$ . For large  $f_{\text{th}}$  the decay is clearly exponential. For  $f_{\text{th}} = 0.4$  the decay is slow and may indicate a power-law decay, although the statistics does not allow us to extract a definite power exponent. There might be a critical value of the threshold  $f_{\text{th}}$  below which the cumulative probability density becomes critical and decays with a power law.

Thus, we may conclude that the difference in the topological structure of the German high-voltage grid as compared to the square grids leads to marked differences in the probability of cascading failures. Its origin and dependence on various topological measures will have to be studied in more detail in future research. It remains to find criteria for the question which arrangements of generators and consumers are beneficial for a particular grid topology in order to minimize the chance and extent of blackouts. In future research we will also study the influence of heterogeneous transmission line capacities and more realistic distributions for the consumed and generated power at each node.

### ACKNOWLEDGMENTS

We gratefully acknowledge the support of BMWi and BMBF, CoNDyNet, FKZ: 03SF0472D.

- Versorgungsqualität SAIDI-Werte 2006-2014, Bundesnetzagentur (2015).
- [2] P. Pourbeik, P. Kundur, and C. Taylor, IEEE Power Energy Mag. 4, 22 (2006).
- [3] M. Vaiman, K. Bell, Y. Chen, B. Chowdhury, I. Dobson, P. Hines, M. Papic, S. Miller, and P. Zhang, IEEE Trans. Power Syst. 27, 631 (2012).
- [4] G. Zhang, Z. Li, B. Zhang, and W. A. Halang, Physica A 392, 3273 (2013).
- [5] P. Hines, K. Balasubramaniam, and E. C. Sanchez, IEEE Potent. 28, 24 (2009).
- [6] A. E. Motter and Y.-C. Lai, Phys. Rev. E 66, 065102(R) (2002).
- [7] A. Bernstein, D. Bienstock, D. Hay, M. Uzunoglu, and G. Zussman, ACM SIGMETRICS Perform. Eval. Rev. 40, 33 (2012).
- [8] A. Bernstein, D. Bienstock, D. Hay, M. Uzunoglu, and G. Zussman, *Proceedings of IEEE INFOCOM 2014* (IEEE, 2014), pp. 2634–2642.
- [9] S. Soltan, D. Mazauric, and G. Zussman, Proceedings of the 5th International Conference on Future Energy Systems, June 11– 13, 2014, Cambridge, United Kingdom (ACM e-Energy, New York, 2014).
- [10] W. Wang, Q. Cai, Y. Sun, and H. He, Proceedings of IEEE Global Telecommunications Conference (GLOBECOM 2011) (IEEE, 2011).
- [11] R. Leelaruji and V. Knazkins, Australasia Universities Power Engineering Conference (AUPEC 2008) (IEEE, 2008).

- [12] Y. Koç, M. Warnier, P. Van Mieghem, R. E. Kooij, and F. M. T. Brazier, Physica A 402, 169 (2014).
- [13] B. A. Carreras, I. Dobson, V. E. Lynch, and D. E. Newman, Chaos 12, 985 (2002).
- [14] I. Dobson, B. A. Carreras, V. E. Lynch, and D. E. Newman, Chaos 17, 026103 (2007).
- [15] B. Paralchev, Statistical Analysis and Modelling of Power Failures, Bachelor's thesis, Jacobs University Bremen, Germany 2014.
- [16] S. Pahwa, C. Scoglio, and A. Scala, Nat. Sci. Rep. 4, 3694 (2014).
- [17] I. Simonsen, L. Buzna, K. Peters, S. Bornholdt, and D. Helbing, Phys. Rev. Lett. **100**, 218701 (2008).
- [18] A. Plietzsch, P. Schultz, J. Heitzig, and J. Kurths, Eur. Phys. J. Spec. Top. 225, 551 (2016).
- [19] D. Witthaut and M. Timme, Eur. Phys. J. B 86, 377 (2013).
- [20] D. Witthaut and M. Timme, Phys. Rev. E 92, 032809 (2015).
- [21] A. R. Bergen and D. J. Hill, IEEE Trans. Power App. Syst. 100, 1 (1981).
- [22] M. Rohden, A. Sorge, M. Timme, and D. Witthaut, Phys. Rev. Lett. 109, 064101 (2012).
- [23] G. Filatrella, A. H. Nielsen, and N. F. Pedersen, Eur. Phys. J. B 61, 485 (2008).
- [24] M. Rohden, A. Sorge, D. Witthaut, and M. Timme, Chaos 24, 013123 (2014).
- [25] W. Medjroubi, C. Matke, and D. Kleinhans, SciGRID v0.1, NEXTENERGY-EWE Research Centre for Energy Technology, http://www.scigrid.de/.

- [26] L. Cuadra, S. Salcedo-Sanz, J. Ser, S. Jimenez-Fernandez, and Z. Geem, Energies 8, 9211 (2015).
- [27] K. I. Goh, D. S. Lee, B. Kahng, and D. Kim, Phys. Rev. Lett. 91, 148701 (2003).
- [28] J. Machowski, J. W. Bialek, and J. R. Bumby, *Power System Dynamics: Stability and Control* (Wiley, West Sussex, 2008).
- [29] K. Heuck, K.-D. Dettmann, and D. Schulz, *Elekrtische Energieversorgung* (Springer, Wiesbaden, 2013).
- [30] T. Nishikawa and A. E. Motter, New J. Phys. 17, 015012 (2015).
- [31] Y. Kuramoto, Lect. Notes Phys. 39, 420 (1975).
- [32] S. H. Strogatz, Physica D 143, 1 (2000).
- [33] J. A. Acebron, L. L. Bonilla, C. J. Perez Vicente, F. Ritort, and R. Spigler, Rev. Mod. Phys. 77, 137 (2005).
- [34] We use the PYTHON function scipy.optimize.fsolve [35,36], which is a wrapper around the functions hybrd and hybrj of the FORTRAN package MINPACK [37].
- [35] E. Jones, T. Oliphant, P. Peterson, and others, *SciPy: Open Source Scientific Tools for Python*, http://www.scipy.org (2001).

- [36] T. E. Oliphant, IEEE Comput. Sci. Eng. 9, 10 (2007).
- [37] J. J. Moré, B. S. Garbow, and K. E. Hillstrom, User Guide for MINPACK-1, Tech. Rep. ANL-80-74 (Argonne National Laboratory, Argonne, IL, USA, 1980).
- [38] R. Delabays, T. Coletta, and P. Jacquod, J. Math. Phys. 57, 032701 (2016).
- [39] D. Labavić, R. Suciu, H. Meyer-Ortmanns, and S. Kettemann, Eur. Phys. J. Spec. Top. 223, 2517 (2014).
- [40] D. Jung and S. Kettemann, Phys. Rev. E 94, 012307 (2016).
- [41] Similarly to the "antiferromagnetic ground state" known from solid-state theory.
- [42] P. Bak, C. Tang, and K. Wiesenfeld, Phys. Rev. Lett. 59, 381 (1987).
- [43] V. Frette, K. Christensen, A. Malthe-Sorenssen, J. Feder, T. Jossang, and P. Meakin, Nature 379, 49 (1996).
- [44] M. Paczuski and S. Boettcher, Phys. Rev. Lett. 77, 111 (1996).
- [45] S. Lübeck and K. D. Usadel, Phys. Rev. E 55, 4095 (1997).