Geometry for a penguin-albatross rookery

Fabio Giavazzi[1,*] and Alberto Vailati[2,†]

[1]*Dipartimento di Biotecnologie Mediche e Medicina Traslazionale, Università degli Studi di Milano, Milano, Italy*
[2]*Dipartimento di Fisica, Università degli Studi di Milano, I-20133 Milano, Italy*

We introduce a simple ecological model describing the spatial organization of two interacting populations whose individuals are indifferent to conspecifics and avoid the proximity to heterospecifics. At small population densities Φ a nontrivial structure is observed where clusters of individuals arrange into a rhomboidal bipartite network with an average degree of 4. For $\Phi \to 0$ the length scale, order parameter, and susceptibility of the network exhibit power-law divergences compatible with hyperscaling, suggesting the existence of a zero-density nontrivial critical point. At larger densities a critical threshold Φ_c is identified above which the evolution toward a partially ordered configuration is prevented and the system becomes jammed in a fully mixed state.

I. INTRODUCTION

The interactions within and between species, either for survival or for the control of the territory, gives rise to a rich spatiotemporal dynamics [1,2], the most investigated case being probably the two species prey-predator system [3–6]. Models and experiments have outlined the emergence of collective behavior and pattern formation arising from simple local rules, a prototypical example being represented by the selfish herd model [7,8]. The study of the complex interplay between inter- and intraspecific interactions and the emergence, at different scales, of structures in the spatial distribution of the individuals, is an important topic in population ecology [9]. Within this framework, the investigation of systems incorporating species competing symmetrically, or quasisymmetrically, for space has attracted recent interest because it has been shown that it leads to the development of a robust diversification of ecosystems [10–12]. The symmetrical interaction can lead to the spatial segregation of species [13], and a competitive advantage can be determined by the emergent formation of spatial structures [14]. Even conspecific seabirds exhibit a spatial segregation into nonoverlapping colonies feeding from mutually exclusive areas [15,16]. Under some circumstances, different species organize into mixed colonies. A nice example is represented by the association between Rockhopper penguin (*Eudyptes chrysocome*) and Black-browed albatross (*Thalassarche melanophrys*) into common colonies, reported by several investigators involved in the census of birds in the Falkland Islands [17–22]. This peculiar system attracted the attention of several novelists in 19th century [23–25], who gave detailed—allegedly fictional—descriptions of the nontrivial topological distribution of the individuals of the two species. In particular, B. Morrell [23] described a penguin-albatross rookery where individuals were uniformly distributed on a nesting area in a staggered configuration where "each albatross is surrounded by four penguins; and each penguin has an albatross for its neighbor, in four directions."

In this work we argue that the description of the topology of a penguin-albatross rookery by Morrell and Poe could

be based on the empirical observation of a real colony rather than on imagination. Taking inspiration from them, we investigate numerically the evolution of the spatial distribution of two populations of "penguins" and "albatrosses" sharing the same area. We show that structures similar to those described in their narratives form by assuming that individuals behave according to two minimal rules: (i) individuals are indifferent to conspecifics (except for avoiding superposition) and (ii) avoid the proximity to heterospecifics. These behavioral rules determine a tendency of the system to phase separate into two clusters made of conspecific individuals. However, during the separation process the system becomes jammed into a structurally arrested configuration whose spatial structure strongly depends on the population density. At small population densities clusters of individuals arrange into a bipartite network spanning the entire system. The network is entirely made of rhomboidal cells, where two clusters of conspecific individuals sit at two opposite corners and two clusters of individuals of the other species occupy the other two corners. This partially ordered phase is characterized by a single dominant length scale represented by the distance between nearest heterospecific individuals and exhibits a nontrivial critical point at vanishing population density, characterized by a divergence of the order parameter in the presence of hyperscaling. At larger population densities the system becomes frozen into a structurally arrested configuration, where the motion of individuals is caged by those surrounding it.

II. MODEL ECOSYSTEM

Our model is of general relevance, but for simplicity we refer to the populations as "Penguins" (P) and "Albatrosses" (A). We assume that each population has an equal number N of individuals [26], initially distributed randomly with density $\Phi = N/M^2$ across the sites of a square lattice of side $M = 128$ with periodic boundary conditions.

The system evolves trough discrete time steps. At each time step, an individual is randomly chosen from the two population with uniform probability. Let d be the Euclidean distance between the selected individual and the nearest individual of the opposite species. The individual attempts a move towards a new site, randomly chosen between one its four neighboring

*fabio.giavazzi@unimi.it
†alberto.vailati@unimi.it

lattice sites. If the new site is occupied by another individual, or if the distance d' between the new site and the nearest individual of the opposite species is smaller than d, the move is rejected. Otherwise it is accepted, and the position of the selected individual is updated accordingly.

In all the simulations discussed in the present work the total number of time steps N_s is chosen in order to ensure that a steady state is reached. A reliable indicator of the convergence to a statistically stable state is represented by the cross-correlation length ξ_c (see Sec. III B). During the evolution of the system a monotonic increase of the cross-correlation length is observed, until it reaches a plateau. For each population density a few exploratory simulations were run and a characteristic equilibration time n_0 was estimated by fitting $\xi_c(n)$ as a function of the simulation step n with a stretched exponential function: $f(n) = A\{1 - \exp[-(n/n_0)^\delta]\} + B$. The duration N_s was then set to a value at least 10 times larger than n_0.

III. RESULTS

The spatiotemporal dynamics of the system is strongly affected by the population density. For small Φ, individuals of one species form small clusters surrounded on average by four clusters of the other species (Fig. 1(a), Supplemental Material movie [27]), a configuration partially reminiscent of that depicted by the narratives of the 19th century. We define a cluster as a connected set of individuals of the same species, say, penguins: Two penguins (sitting on sites x and x', respectively) belong to the same cluster if a continuous path exists connecting x and x' made up of steps connecting neighboring sites all occupied by penguins. At intermediate Φ individuals arrange into large connected fractal clusters staggered across the plane [Fig. 1(b), Supplemental Material movie [27]]. These clusters are separated by convoluted aisles free of individuals. At a critical density $\Phi_c \approx 0.311$ the system undergoes a transition to a disordered phase [Fig. 1(c), Supplemental Material movie [27]]. The free aisles disappear and the distribution of species exhibits large concentration fluctuations determined by the strong spatial heterogeneity of the configurations (see the late stages of Supplemental Material movie [27]). At $\Phi > \Phi_c$ the system becomes jammed into a disordered configuration [Fig. 1(d), Supplemental Material movie [27]] where the limited number of free sites determines a structural arrest that prevents the development of long-range correlations, similarly to what happens in a glassy state. Thus the system exhibits two critical points at $\Phi_0 = 0$ and $\Phi_c = 0.311$. In the following we will describe the properties of these two critical points separately.

A. Structural arrest at high Φ

Due to the interspecific repulsion between individuals, a structural arrest transition takes place at packing fractions well below the close packing of individuals needed for the occurrence of the huddling transition reported in a single species [28–30]. The transition is characterized by the appearance of clusters of individuals with fractal dimension $d_f \approx 1.4$. The cluster size distribution close to the critical point is compatible with a power law with exponent -1.4, with an exponential cutoff. By choosing the average cluster size Ψ_a as a suitable order parameter one can appreciate that the transition is characterized by a continuous decrease of the order parameter [Fig. 2(a)], accompanied by a divergence of the susceptibility of the system χ_a [Fig. 2(b)], which represents the variance of the order parameter across different realizations. The average correlation length ξ_a of the clusters increases continuously at the transition [Fig. 2(c)], without showing any sign of divergence. This feature suggests that the critical behavior is not determined by the development of intraspecies long-range order.

The critical transition can be best understood as due to a percolation [31] of the empty phase. The availability of aisles of empty sites spanning the whole lattice represents an essential ingredient to achieve the long-range order exhibited by the system at $\Phi < \Phi_c$. Individuals cannot move when they are stuck inside a cluster of individuals of the same species. The only region accessible for movement is determined by
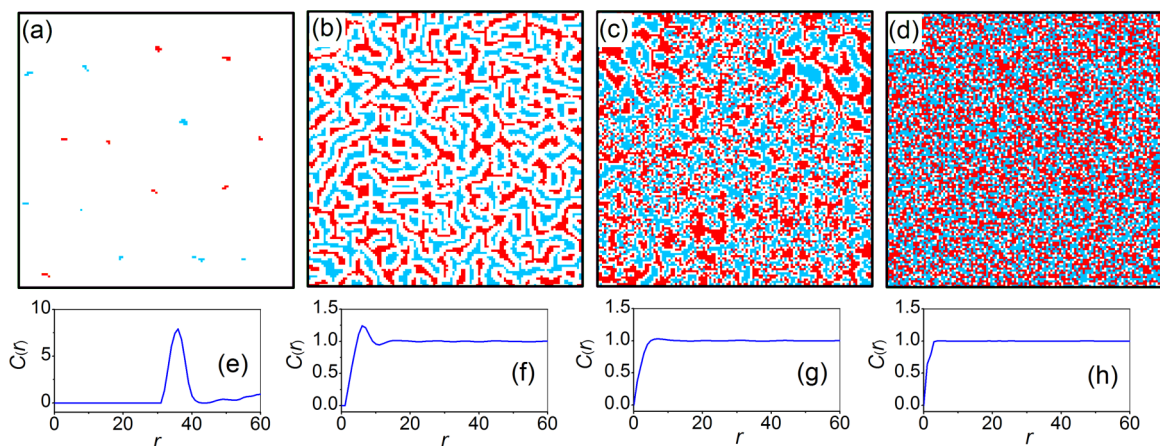

FIG. 1. (Color online) Patterns formed at steady state at different population densities. Red and light blue (dark gray and light gray) pixels mark the position of individuals of the two species, while empty sites are white. (a) $\Phi = 0.002$, (b) $\Phi = 0.2$, (c) $\Phi = 0.31$, and (d) $\Phi = 0.4$. Panels (e)–(h) show the corresponding cross-correlation functions. The evolution of the patterns is shown in the Supplemental Material movie [27]. The size of the simulation box is 128×128 sites.
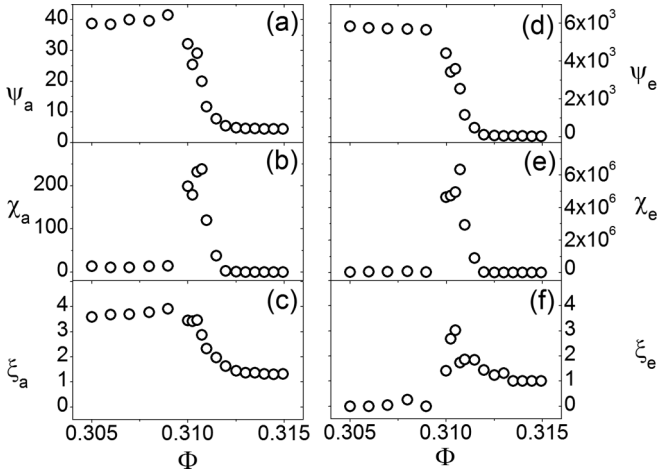
FIG. 2. [(a)–(d)] Order parameter, [(b)–(e)] susceptibility, and [(c)–(f)] correlation length of the system. The left column corresponds to a single species, while the right one to the empty region.
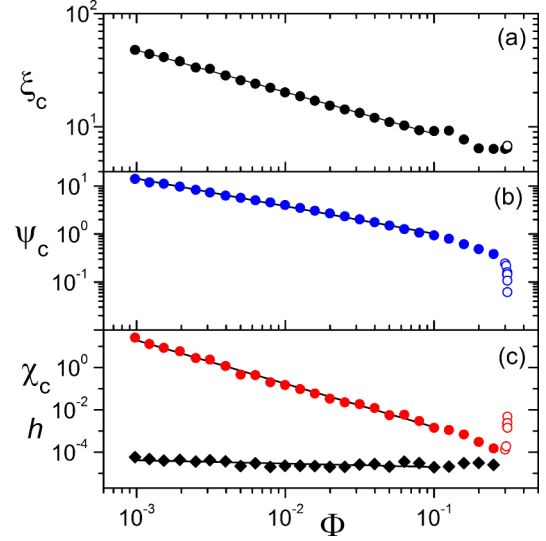

FIG. 3. (Color online) (a) Correlation length, (b) order parameter, and (c) susceptibility (circles) and hyperscaling parameter (diamonds) characterizing the mutual order of the two species defined from the cross-correlation function $C_c(r)$. Open and solid symbols represent averages on 20 and 100 independent realizations of the system, respectively. Lines represent the best fit of the data with a power law. The hyperscaling parameter h (diamonds) is reasonably constant, indicating that the power-law exponents are compatible with hyperscaling.

the empty sites that separate clusters of different species. In the range $0.1 < \Phi < 0.311$ the empty phase has a labyrinthine structure spanning the entire system. The relevant parameters in this case are the largest cluster size Ψ_e [Fig. 2(d)], its variance χ_e [Fig. 2(e)], and the correlation length ξ_e of the empty region [Fig. 2(f)], determined from the size of the second-largest cluster of empty sites [31]. The correlation length ξ_e and the susceptibility χ_e diverge close to the critical point, while the order parameter decreases continuously from a finite value to zero. This finite value is close to the maximum cluster size that one would attain by using all the available empty sites. Ideally, we would like to determine the critical exponents at the phase transition by a systematic investigation of the finite-size scaling [31] in order to determine the universality class of the transition. Unfortunately, the slowing down of the system dynamics close to the transition (see the Supplemental Material movie [27]) induces divergence of the time needed to reach a stable configuration, which prevented us from achieving a reliable determination of the exponents.

B. Nontrivial critical point at $\Phi = 0$

At $\Phi < 0.31$ the individuals are able to arrange into clusters and an order related to the mutual spatial distribution of clusters of different species develops. The position of penguins and albatrosses is specified by two spatial distributions $P(x)$ and $A(x)$ respectively, where $P(x) = 1$ [$A(x) = 1$] for each site occupied by a penguin (albatross) and 0 otherwise. The statistical relation between the mutual positions of individuals of the two species is expressed by the radial cross-correlation function $C_c(r) = \langle P(x)A(x+r)\rangle/\Phi^2$. At small Φ the cross-correlation function exhibits a narrow peak [Fig. 1(e)], while the autocorrelation functions of $P(x)$ and $A(x)$ are almost featureless (not shown). By increasing Φ the cross-correlation peak becomes less defined [Figs. 1(f) and 1(g)] and for $\Phi > \Phi_c$ it disappears completely [Fig. 1(h)], thus marking the transition to a featureless phase where the mutual position between individuals of different species is uncorrelated. The position

of the peak of $C(r)$ represents the cross-correlation length ξ_c of the system [Fig. 3(a)], which provides a typical length scale for the distance between clusters of different species. The contrast of the first peak of the cross-correlation function $\Psi_c = C_{\max} - 1$ represents an order parameter suitable to characterize the formation of interspecies ordered structures [Fig. 3(b)]. The variance of the order parameter [Fig. 3(c)] provides the susceptibility χ_c to fluctuations. At small Φ these parameters exhibit a power-law behavior, $\Psi_c \propto \Phi^\beta$, $\chi_c \propto \Phi^{-\gamma}$, and $\xi_c \propto \Phi^{-\nu}$, with exponents $\beta = -0.575 \pm 0.015$, $\gamma = 2.05 \pm 0.06$, and $\nu = 0.37 \pm 0.01$. The exponents are compatible with the hyperscaling relation $2\beta + \gamma = d\nu$, where $d = 2$ is the dimensionality of the space. We test the goodness of the hyperscaling of our data by plotting the hyperscaling parameter $h = \chi_c/(\Psi_c\xi_c)^2$ [diamonds on Fig. 3(c)]. In the presence of hyperscaling this parameter must be constant. Hyperscaling suggests that $\Phi_c = 0$ exhibits the features of a critical point. At variance with the usual behavior of second-order phase transitions, in this case the order parameter diverges close to this critical point and the critical exponent β is negative, a feature exhibited by other systems where hyperscaling holds, such as the thermal denaturation of DNA [32]. The dynamic evolution of the patterns at small Φ is characterized by the transition from a disordered to an ordered configuration exhibiting on average fourfold symmetry (Supplemental Material movie [27]).

The initial configuration of the system corresponds to a random distribution of individuals of the two species. At this stage, the domain of influence of each individual is a Voronoi cell with an average of six sides. During the evolution, individuals belonging to the same species form clusters. At steady
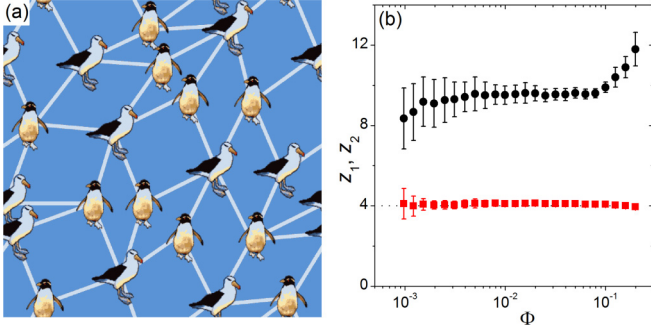
FIG. 4. (Color online) (a) Bipartite rhomboidal penguin-albatross network formed at $\Phi = 0.002$. (b) Average number of first neighbors z_1 (squares) and of second neighbors (circles) z_2 of the network. z_1 remains stable at 4 in a wide range of Φ, in agreement with the narratives by Morrell and Poe.

state the spatial distribution of the clusters is characterized by a bipartite random network [33] obtained by joining clusters of penguins with the nearest clusters of albatrosses. The network is made of rhomboidal loops where two clusters of penguins sit at opposite vertices of a rhombus and two cluster of albatrosses at the other two vertices [Fig. 4(a)]. Similar configurations have been reported for the bioconvection of bacteria [34,35], for the spoke pattern convection of simple fluids occurring at high Rayleigh numbers [36,37], and for the solutal convection of nanoparticles at high solutal Rayleigh numbers [38,39].

The average degree z_1 of a cluster can be determined by averaging the number of nearest neighbors of each cluster in the network [Fig. 4(b)] [40]. At small Φ, $z_1 \approx 4$, and the average number of second neighbors is $z_2 \approx 8$, in agreement with the fact that on average z_1 second neighbors are shared by the nearest neighbors. At larger fractions, z_1 remains stable at 4, while z_2 increases to 12, a value compatible with a tree structure of the bipartite network.

The peculiar pattern exhibited at vanishing Φ has proven to be robust against perturbations breaking the perfect symmetry between the two species. In particular, we ran simulations in the presence (i) of a moderate unbalance in the number N_1 and N_2 of individuals within each population and (ii) of a difference in the mobility of the two species, implemented by introducing a systematic bias in the probability P of selecting an individual from a given species for attempting a move. Within the investigated range $1 \leqslant \frac{N_1}{N_2} \leqslant 2, 0.1 \leqslant P \leqslant 0.5$ no significant variation can be detected in the spatial distributions at steady state when compared to the unperturbed, fully symmetric, situation.

IV. DISCUSSION

A. Topology of the network

The average fourfold coordination of the vertices of the network can be understood from Euler's theorem. The theorem states that for an infinite network $R + V - E = 0$, where R is the number of rhomboidal loops, V the number of vertices, and E the number of edges connecting two clusters of different species. Each loop has four edges and, since each edge is shared by two cells, $E = 2R$. On the same argument, the number of

edges is related to the number of vertices by $E = z_1 V/2$. By combining these relations one gets that $z_1 = 4$. Therefore, when the population density is small, each cluster of one species is surrounded on average by four clusters of the other species, a configuration that mirrors the structure reported for the distribution of individuals in the penguin-albatross rookery described in the narratives by Morrell and Poe [41]. Under this condition the system is dominated by a single characteristic length scale corresponding to the average distance between nearest clusters of different species ξ_c [Fig. 3(a)]. This fourfold coordination of the network formed by the two species differs dramatically from the trivial average sixfold coordination, which is expected in colonies of individuals belonging to the same species. Indeed, the development of a fourfold coordination puts more stringent constraints on the mutual position of the individuals than the sixfold coordination, which is also present for a random distribution of individuals [39,42].

B. Topology of a penguin-albatross rookery in the narratives of the 19th century

At the moment, no systematic investigation of the topology of mixed colonies of penguins and albatrosses based on empirical data has been performed to our knowledge. However, the similarities between the networks generated by our model and those described by Morrell and Poe suggest that their description of a penguin-albatross rookery could be inspired by reality more than expected. The book *A Narrative of Four Voyages* by Benjamin Morrell was first published in 1832 by J&J Harper in New York, while *The Narrative of Arthur Gordon Pym of Nantucket* by Edgar Allan Poe was published by the same publisher in 1838. Literary criticism [43] has shown that Poe largely borrowed from Morrell, making it unlikely that Poe himself witnessed real penguin-albatross rookery. Benjamin Morrell was a sea captain and he actually performed several expeditions across the world. The reliability of his narratives has been questioned by historians and geographers, but the fact that his book contains a large amount of accurate information convinced other investigators that his accounts were partially based on reality [44]. As far as the observation of a penguin albatross-rookery is concerned, Morrell provides a precise description of the geographical location of the colony in New Island, one of the Falkland Islands. Indeed, associations between the Rockhopper penguin and the Black Browed albatross have regularly been reported from recent bird population censuses on these islands [17–22]. This fact, together with the description of the topology of the rookery akin to that described by our model, strongly suggests that the report on the mixed rookery by Morrell could be based on actual observation rather than on fiction or literary imagination.

C. Preliminary comparison with a real mixed colony

The ideal fourfold coordination reported by Morrell and found in our model system can be difficult to observe in natural systems, due to the presence of asymmetries in the behavior or in the interaction between the two species. As an example, in the case of the penguin-albatross mixed colony, asymmetries can arise from slight differences in the

breeding periods and by the constraints in the distributions of individuals determined by peculiarities of the terrain [45]. In general, one could expect that one of the species settles on the terrain at first, thus providing a template for the final structure of the colony, followed by the second [45]. An analysis of real mixed colonies to develop a model that accurately describes all the peculiarities related to temporal and spatial asymmetries would require an extensive set of data showing the time evolution of the position of birds within the colonies during different breeding seasons and in different nesting areas. At the moment we are not aware of the availability of empirical data showing the dynamics of a mixed colony. Indeed, the acquisition of such data from an aerial position through a breeding season would require a significant technical and economical effort. However, we are in the position of attempting a preliminary analysis by using an aerial photograph of a mixed colony kindly made available to us by Ian Strange and Georgina Strange [Fig. 5(a)] [21,46]. The photograph was taken under uncontrolled conditions: We do not know whether the albatrosses or the penguins occupied the site first and we are not aware of the stage of breeding for the two species. Moreover, during other breeding seasons the conditions can differ significantly from those reported by the photograph. Notwithstanding the peculiar nature of the conditions of the photograph, we think that a preliminary analysis of the geometry of a penguin-albatross rookery is meaningful to start highlighting some general features of a real natural system.

The presence of an interspecies organization is evidenced by the fact that penguins tend to be located at the boundaries of the Voronoi cells generated by albatrosses [Fig. 5(b)]. Moreover, the cross-correlations function of the position of penguins and albatrosses exhibits a marked peak [Fig. 5(c)]. The radial distribution function of penguins is almost feature-less, as it occurs in our model, but that of albatrosses presents a marked peak witnessing the presence of intraspecies order [Fig. 5(c)], a feature at variance with our model. Intraspecies order is also witnessed by the area distribution of the Voronoi cells generated by albatrosses [Fig. 5(d)]. This distribution differs substantially from that associated to a set of randomly distributed generators (spatial Poisson process [42]) and is compatible with the distribution generated by a Hasegawa-Tanemura adjustment process [47]. In a Hasegawa-Tanemura adjustment process each individual changes its position in the attempt to be as close as possible to the center of its Voronoi cell, defined as the center of mass of the vertices of the cell. In the absence of a second species, such a process leads to the spread across the colony of conspecific individuals. The analysis of the photograph allows us to identify a tentative scenario of the temporal and spatial interactions between the birds that accounts for the observed structures. (1) Albatrosses have occupied the nesting site first and spread themselves across the available area. Then they have built their nest. The nests of Black-browed albatrosses are piles made of mud and grass whose position, once established, cannot be changed easily. (2) Once the nests have been built and the position of albatrosses became stable, penguins have joined and filled the free space. Again, we are not aware whether the penguins shown in Fig. 5(c) had already established the position of their nests. The temporal asymmetry suggested above implies
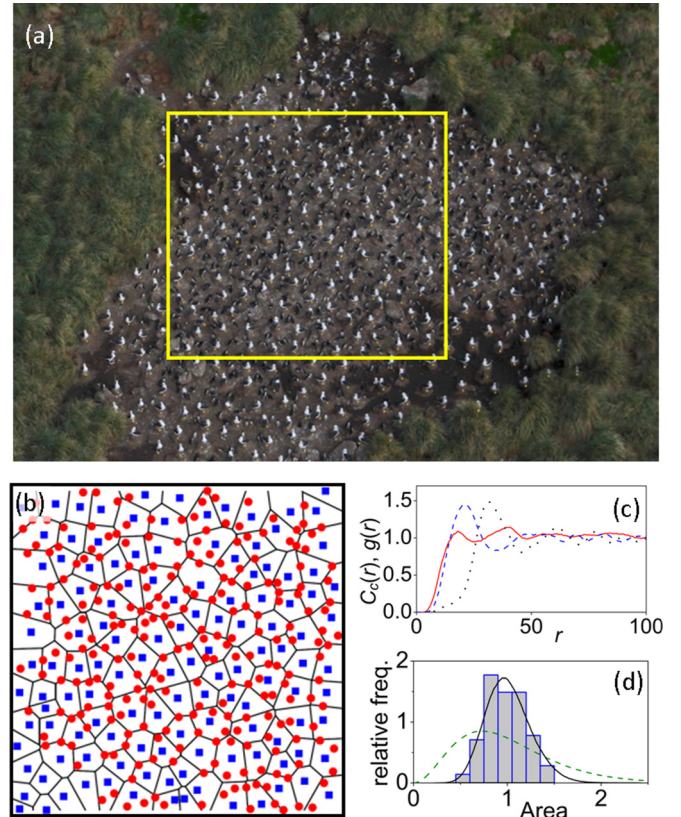

FIG. 5. (Color online) (a) Aerial photograph of a Rockhopper penguin–Black-browed albatross colony at Elephant Jason Island, Falkland Islands [21]. Image copyright by Ian J. Strange and Georgina Strange, *Design in Nature*. The box delimits the region of interest. (b) Position of the albatrosses (blue squares) and penguins (red circles) in the region of interest. The network represents the Voronoi diagram generated by albatrosses. The size of the selected region is about 18 m × 18 m. (c) Spatial cross-correlation function $C_c(r)$ for the position of penguins and albatrosses (dotted line); radial distribution functions $g(r)$ of penguins (solid line) and albatrosses (dashed line). (d) Histogram of the frequency distribution of the area of the Voronoi cells generated by albatrosses. Areas are normalized by the average area. The solid line represents the Hasegawa-Tanemura distribution [47]. The dashed line corresponds to a random distribution of albatrosses.

that the species that settled first on the terrain had already reached a stationary condition when the other one joined, thus preventing a mutual redistribution of individuals. As a result of the asymmetric interaction between the species, the structures shown in Fig. 5 are not compatible with the bipartite network predicted by our model and reported by Morrell and Poe.

V. CONCLUSIONS

We have shown that a simple model where two species interact by mutually avoiding each other leads to the development of clusters of individuals arranged to form interspecies structures, where each cluster is surrounded on average by four clusters of the other species. The system exhibits a nontrivial critical

point at vanishing population densities. At larger population densities, the spatiotemporal dynamics is determined by a percolation of the empty region. A preliminary analysis of empirical data shows that real systems develop interspecies order, but the structure is affected by nonsymmetrical interactions between the species. A more detailed investigation of a real system with symmetrical interactions is made difficult by the scant empirical data currently available and would ideally require the analysis of time-lapse movies of mixed colonies across several breeding seasons and different breeding sites.

[1] I. D. Couzin and J. Krause, Adv. Study Behav. **32**, 1 (2003).

[2] T. Vicsek and A. Zafeiris, Phys. Rep. **517**, 71 (2012).

[3] D. A. Vasseur and W. Fox, Nature **460**, 1007 (2009).

[4] A. Kamimura and T. Ohira, New J. Phys. **12**, 053013 (2010).

[5] T. Vicsek, Nature **466**, 43 (2010).

[6] L. Angelani, Phys. Rev. Lett. **109**, 118104 (2012).

[7] W. D. Hamilton, J. Theor. Biol. **31**, 295 (1971).

[8] A. J. King, A. M. Wilson, S. D. Wilshin, J. Lowe, H. Haddadi, S. Hailes, and A. J. Morton, Curr. Biol. **22**, R561 (2012).

[9] P. Legendre and M. J. Fortin, Vegetatio **80**, 107 (1989).

[10] J. Mathiesen, N. Mitarai, K. Sneppen, and A. Trusina, Phys. Rev. Lett. **107**, 188101 (2011).

[11] N. Mitarai, J. Mathiesen, and K. Sneppen, Phys. Rev. E **86**, 011929 (2012).

[12] J. Juul, K. Sneppen, and J. Mathiesen, Phys. Rev. E **87**, 042702 (2013).

[13] N. J. Gotelli, G. R. Graves, and C. Rahbek, Proc. Natl. Acad. Sci. U.S.A. **107**, 5030 (2010).

[14] E. Heinsalu, E. Hernandez-Garcia, and C. Lopez, Phys. Rev. Lett. **110**, 258101 (2013).

[15] E. D. Wakefield *et al.*, Science **341**, 68 (2013).

[16] H. Weimerskirch, Science **341**, 35 (2013).

[17] I. J. Strange, Le Gerfaut **72**, 137 (1982).

[18] N. Huin, Falkland Islands Penguin Census 2005/06, Report to Falklands Conservation (2007).

[19] L. Demongin, M. Poisbleau, R. Matias, and J. P. Granadeiro, Ornitol. Neotrop. **20**, 629 (2009).

[20] S. Crofts, A. Wolfaardt, and A. Baylis, "Storm damage to colonial seabirds at Beauchene Island and Steeple Jason Island during 13-14 December 2010", Report to Falklands Conservation (2011).

[21] I. J. Strange and G. Strange, "Black-browed albatross Thalassarche melanophris. Aerial photographic survey methodology: advantages over ground surveying in the Falkland Islands", Results and comparisons with surveys in 2005 and 2010 (Design in Nature, Falkland Islands, 2011).

[22] A. Baylis, "2010 Archipelago-wide census of gentoo and rockhopper penguins breeding at the Falkland Islands", Report to Falklands Conservation (2012).

[23] B. Morrell, *A Narrative of Four Voyages* (J & J Harper, New York, 1832).

[24] E. Allan Poe, *The Narrative of Arthur Gordon Pym of Nantucket* (Harper & Brothers, New York, 1838).

[25] J. Verne, *Le Sphinx des glaces* (Hetzel, Paris, 1897).

[26] Representative values of the ratio of penguins to albatrosses in mixed colonies obtained from direct observation range from 5:1 [17] to 1:3 [20].

[27] See Supplemental Material at http://link.aps.org/supplemental/10.1103/PhysRevE.89.052706 for a movie showing the time evolution of the system at different population densities.

[28] D. P. Zitterbart, B. Wienecke, J. P. Butler, and B. Fabry, PLoS ONE **6**, e20260 (2011).

[29] M. Canals and F. Bozinovic, Complexity **17**, 35 (2011).

[30] R. C. Gerum, B. Fabry, C. Metzner, M. Beaulieu, A. Ancel, and D. P. Zitterbart, New J. Phys. **15**, 125022 (2013).

[31] K. Binder, Rep. Prog. Phys. **60**, 487 (1997).

[32] S. Buyukdagli and M. Joyeux, Phys. Rev. E **73**, 051910 (2006).

[33] M. E. J. Newman, S. H. Strogatz, and D. J. Watts, Phys Rev. E **64**, 026118 (2001).

[34] J. R. Platt, Science **133**, 1766 (1961).

[35] I. M. Janosi, J. O. Kessler, and V. K. Horvath, Phys. Rev. E **58**, 4793 (1998).

[36] F. H. Busse and J. A. Whitehead, J. Fluid Mech. **66**, 67 (1974).

[37] F. H. Busse, in *Fluid and Gasdynamics*, edited by G. H. Schnerr, R. Bohning, W. Frank, and K. Bühler (Springer, Wien, 1994), Acta Mech. (Suppl.) vol. 4, p. 11, doi:10.1007/978-3-7091-9310-5_2

[38] R. Cerbino, A. Vailati, and M. Giglio, Phys. Rev. E **66**, 055301(R) (2002).

[39] S. Mazzoni, F. Giavazzi, R. Cerbino, M. Giglio, and A. Vailati, Phys. Rev. Lett. **100**, 188104 (2008).

[40] One cluster A is the nearest neighbor of a cluster P of the other species if, and only if, an individual belonging to A is the nearest neighbor of an individual belonging to P.

[41] The average number of individuals per cluster decreases as the population density Φ is decreased. As an example, the average number of individuals per cluster in configurations like that shown in Fig. 1(a) ($\Phi = 0.002$) is of the order of 3, while in the limit of very small Φ ($\Phi = 4 \times 10^{-5}$) we find that the average number of individuals per cluster is of the order of 1.5.

[42] A. Okabe, B. Boots, K. Sugihara, and S. N. Chiu, *Spatial Tessellations* (Wiley, New York, 1999).

[43] L. Gitelman, Nineteenth-Cent. Lit. **47**, 349 (1992).

[44] R. T. Gould, *Enigmas* (Philip Allan & Co., London, 1929).

[45] Nicola Saino (private communication).

[46] The image was dilated in the vertical direction by 15% to compensate for the deformation due to a parallax angle of about 30° with respect to the perpendicular to the terrain. No attempt was made to introduce additional corrections for perspective.

[47] M. Hasegawa and M. Tanemura, Ann. Inst. Statist. Math. **28**, 509 (1976).