# Fluctuation-dissipation theory of input-output interindustrial relations

Hiroshi Iyetomi<sup>\*</sup> and Yasuhiro Nakayama Department of Physics, Niigata University, Ikarashi, Niigata 950-2181, Japan

Hideaki Aoyama Department of Physics, Kyoto University, Kyoto 606-8501, Japan

> Yoshi Fujiwara ATR Laboratories, Kyoto 619-0288, Japan

Yuichi Ikeda Hitachi Research Laboratory, Hitachi Limited, Ibaraki 319-1221, Japan

Wataru Souma

College of Science and Technology, Nihon University, Chiba 274-8501, Japan (Received 23 February 2010; revised manuscript received 26 October 2010; published 11 January 2011)

In this study, the fluctuation-dissipation theory is invoked to shed light on input-output interindustrial relations at a macroscopic level by its application to indices of industrial production (IIP) data for Japan. Statistical noise arising from finiteness of the time series data is carefully removed by making use of the random matrix theory in an eigenvalue analysis of the correlation matrix; as a result, two dominant eigenmodes are detected. Our previous study successfully used these two modes to demonstrate the existence of intrinsic business cycles. Here a correlation matrix constructed from the two modes describes genuine interindustrial correlations in a statistically meaningful way. Furthermore, it enables us to quantitatively discuss the relationship between shipments of final demand goods and production of intermediate goods in a linear response framework. We also investigate distinctive external stimuli for the Japanese economy exerted by the current global economic crisis. These stimuli are derived from residuals of moving-average fluctuations of the IIP remaining after subtracting the long-period components arising from inherent business cycles. The observation reveals that the fluctuation-dissipation theory is applicable to an economic system that is supposed to be far from physical equilibrium.

DOI: 10.1103/PhysRevE.83.016103

PACS number(s): 89.65.Gh, 05.40.-a, 87.23.Ge, 89.75.Fb

## I. INTRODUCTION

Both the business cycle and the interindustrial relationship are long-standing basic issues in the field of macroeconomics, and they have been addressed by a number of economists. Recently, we analyzed [1] business cycles in Japan using indices of industrial production (IIP), an economic indicator that measures current conditions of production activities throughout the nation on a monthly basis. Careful noise elimination enabled us to extract business cycles with periods of 40 and 60 months that were hidden behind complicated stochastic behaviors of the indices.

In this accompanying paper, we focus our attention on the interindustrial relationship by analyzing IIP data in a framework of the linear response theory; the fluctuationdissipation (FD) theory plays a vital role in the analysis of IIP data. We also discuss the difference between moving-average fluctuations in the original data and long-period components arising from inherent business cycles. The residuals may be interpretable as a sign of external stimuli to the economic system. The recent worldwide recession offers us a good opportunity to conduct this study because it delivered an unprecedented shock to the economic system of Japan. tionally represented by a matrix in which each column lists the monetary value of an industry's inputs and each row lists the value of the industry's outputs, including final demand for consumption. Such a matrix, called the input-output table, was developed by Leontief [2,3]. This table thus measures how many goods of one industrial sector are used as inputs for production of goods by other industrial sectors and also the extent to which internal production activities are influenced by change in final demand. Leontief's input-output analysis can be regarded as a simplified model of Walras's general equilibrium theory [4] to implement real economic data for carrying out an empirical analysis of such economic interactions. Currently the basic input-output table is constructed every 5 years, according to the System of National Accounts, by the Ministry of Internal Affairs and Communications in Japan.

The interindustrial relations of an economy are conven-

It should be noted that the input-output table describes yearly averaged interindustrial relations. Although such a poor time resolution of the table may be tolerable for budgeting of the government on an annual basis, various day-to-day issues faced by practitioners require them to react promptly. We thus need a more elaborate methodology that enables investigation of the input-output interindustrial relationship with a much higher time resolution.

Econophysics [5–8] is a newly emerging discipline in which physical ideas and methodologies are applied for

<sup>\*</sup>hiyetomi@sc.niigata-u.ac.jp

understanding a wide variety of complex phenomena in economics. We adopt this approach to address the previously mentioned issues in macroeconomics; that is, we pay maximum attention to real data while drawing any conclusions. Of course, it is important to remember that real data are possibly contaminated with various kinds of noise. The random matrix theory (RMT), combined with principal component analysis, has been used successfully to extract genuine correlations between different stocks hidden behind complicated noisy market behavior [9–16]. Recently, dynamical correlations in time series data of stock prices have been analyzed by combining Fourier analysis with the RMT [17].

In this study, we further develop the noise elimination method initiated in previous studies. The null hypothesis that has been adopted thus far for extracting true mutual correlations corresponds to shuffling time series data in a completely random manner. Although we should distinguish between mutual correlations and autocorrelations, both these correlations are destroyed at the same time by the completely random shuffling. To solve this problem, rotational random shuffling of data in the time direction is introduced as an alternative null hypothesis; such randomization preserves autocorrelations involved in the original data. The new null hypothesis thus elucidates the concept of noise elimination for mutual correlations.

Furthermore, we borrow the concept of the FD theory [18] from physics to elucidate the interindustrial relationship and the response of an economic system to external stimuli. The theory establishes a direct relationship between the fluctuation properties of a system in equilibrium and its linear response properties. We assume that the validity of the FD theory in physical systems is also true for such an exotic system, as described by the IIP. Very recently, dynamics of the macroeconomy has been studied in the linear response theory by taking an explicit account of heterogeneity of microeconomic agents [19].

This paper is organized as follows. In Sec. II, we first provide a brief review of the noise elimination from the IIP using the RMT. A new null hypothesis based on rotational random shuffling is introduced in Sec. III. Section IV presents construction of a genuine correlation matrix for the IIP by consideration of only those dominant modes that are approved to be statistically meaningful by the RMT. We present development of a FD theory for input-output interindustrial relations in Sec. V. Then, in Sec. VI, we quantitatively discuss the relationship between shipments of final demand goods and production of intermediate goods. In Sec. VII, we elucidate the response of the industrial activities to external stimuli by subtracting long-period components arising from inherent business cycles from moving-average fluctuations in the original data. Section VIII concludes.

## II. APPLICATION OF RANDOM MATRIX THEORY TO IIP

In Japan, the IIP are announced monthly by the Ministry of Economy, Trade, and Industry [20]. For this study, we will choose seasonally adjusted data instead of original data. Two classification schemes of the IIP are available: indices classified by industry and indices classified by use of goods. We adopt the latter classification scheme because we are



FIG. 1. Input-output relationship in industrial activities of economic system as measured by IIP in Japan. The numbers in the parentheses denote the classification index g in Table I.

interested in input-output interindustrial relations here, which are measured by correlations between shipments of final demand goods and production of intermediate goods in the IIP data. The concept is illustrated in Fig. 1. We emphasize that the inner loop of production existing in the economic system may give rise to a nonlinear feedback mechanism to complicate the dynamics of the system; outputs are reused by the system as inputs for its production activities. Table I lists the categories<sup>1</sup> of goods along with weights assigned to each of them for computing the average IIP. These weights are proportional to value added produced in the corresponding categories, and their total sum amounts to 10 000. Unfortunately, the resolution of the IIP data for the producer goods is quite poor, which are just categorized as mining and manufacturing and as others.

Figure 2 shows the temporal change of the averaged IIP data for production, shipments, and inventory during the period of January 1988 to June 2009. The ongoing global recession is traced back to the subprime mortgage crisis in the United States, which became apparent in 2007. The economic shock has affected Japan without exception, leading to a dramatic drop in the production activities of the country, as shown in the figure.

Since some of the entries, such as g = 16 and 17, are missing before January 1988, we use the data [20] for the 240 months from January 1988 to December 2007. Furthermore, this chosen period for the study excludes the abnormal behavior of the IIP data due to the Great Recession. We denote the IIP data for goods as  $S_{\alpha,g}(t_j)$ , where  $\alpha = 1,2$ , and 3 for production (value added), shipments, and inventory, respectively. Similarly, g = 1,2,...,21 denotes the 21 categories of goods, and  $t_j = j\Delta t$  with  $\Delta t = 1$  month and j =1,2,...,N(= 240); j = 1 and j = N correspond to 1/1988 and 12/2007, respectively. The logarithmic growth rate  $r_{\alpha,g}(t_j)$ is defined as

$$r_{\alpha,g}(t_j) := \log_{10} \left[ \frac{S_{\alpha,g}(t_{j+1})}{S_{\alpha,g}(t_j)} \right],\tag{1}$$

where *j* runs from 1 to N' := N - 1(=239). Then it is normalized as

$$w_{\alpha,g}(t_j) := \frac{r_{\alpha,g}(t_j) - \langle r_{\alpha,g} \rangle_t}{\sigma_{\alpha,g}},\tag{2}$$

<sup>&</sup>lt;sup>1</sup>See Ref. [1] for details on the classification.

TABLE I. Classification of goods according to IIP. First, the goods are classified into two categories, "final femand" and "producer," and then those two categories are divided into 19 and 2 subcategories, respectively. The central column specifies the index g for the subcategories, which has, in total, 21 values. The number in the parentheses associated with each species of goods shows its weight used to compute the averaged IIP; the weights are normalized so that their total sum is 10 000.

	No.	Final demand goods (4935.4)	
	Investment goods (2352.5)		
Capital goods (1662.1)	1	manufacturing equipment (530.7)	
	2	electricity (148.1)	
	3	communication and broadcasting (48.8	
	4	agriculture (31.0)	
	5	construction (129.6)	
	6	transport (381.3)	
	7	offices (175.4)	
	8	other capital goods (217.2)	
Construction goods (690.4)	9	construction (568.1)	
	10	engineering (122.3)	
	Consumer goods (2582.9)		
Durable consumer goods (1267.9)	11	housework (62.3)	
	12	heating and cooling equipment (62.5)	
	13	furniture and furnishings (43.4)	
	14	education and amusement (246.5)	
	15	motor vehicles (853.2)	
Nondurable	16	housework (649.7)	
Consumer goods (1315.0)	17	education and amusement (105.2)	
	18	clothing and footwear (92.2)	
	19	food and beverage (467.9)	
		Producer goods (5064.6)	
	20	mining and manufacturing (4601.7)	
	21	others (462.9)	

where  $\langle \cdot \rangle_t$  denotes average over time  $t_1, \ldots, t_{N'}$  and  $\sigma_{\alpha,g}$  is the standard deviation of  $r_{\alpha,g}$  over time. Definition (2) ensures that the set  $w_{\alpha,g} := \{w_{\alpha,g}(t_1), w_{\alpha,g}(t_2), \ldots, w_{\alpha,g}(t_{N'})\}$  has an average of zero and a standard deviation of 1.

Figure 3 shows an overview of how the volatility  $w_{\alpha,g}^2(t_j)$  of the standardized IIP data behaves on a time-goods plane.



FIG. 2. Averaged IIP data  $S_{\alpha}$  for production (thick solid line), shipment (thin solid line), and inventory (dotted line) as a function of time *t*. The correlation matrix is calculated using the data in the shaded area from January 1988 to December 2007.

Unfortunately, the visualization does not allow for detecting any correlations involved in the IIP data. One may even doubt whether useful information on interindustrial relations truly exists in the data.

To answer the obvious question that would arise here, we begin with calculating the equal-time correlation matrix C of  $\{w_{\alpha,g}\}$  according to

$$C_{\alpha,g;\beta,h} = \langle w_{\alpha,g}(t)w_{\beta,h}(t)\rangle_t,\tag{3}$$



FIG. 3. Bird's-eye view of volatility of standardized IIP data in a panel form, where the index  $\ell$  is defined as  $\ell := 21(\alpha - 1) + g$ .

whose diagonal elements are unity by definition of the normalized growth rate  $w_{\alpha,g}(t_j)$ . Since  $\alpha(\beta)$  runs from 1 to 3 and g(h) runs from 1 to 21, the matrix C has  $M \times M$ (M = 63) components. We denote the eigenvalues and the corresponding eigenvectors of the correlation matrix as  $\lambda^{(n)}$ and  $V^{(n)}$ , respectively:

$$\boldsymbol{C}\boldsymbol{V}^{(n)} = \boldsymbol{\lambda}^{(n)}\boldsymbol{V}^{(n)},\tag{4}$$

where the eigenvalues are sorted in descending order of their values and the norm of eigenvectors is set to unity.

On the basis of the eigenvectors  $V^{(n)}$  thus obtained, the normalized growth rate  $w_{\alpha,g}(t_i)$  can be decomposed into

$$w_{\alpha,g}(t_j) = \sum_{n=1}^{M} a_n(t_j) \, V_{\alpha,g}^{(n)}.$$
 (5)

The correlation matrix C is also decomposable in terms of the eigenvalues and eigenvectors as

$$C = \sum_{n=1}^{M} \lambda^{(n)} V^{(n)} V^{(n)\mathrm{T}}.$$
 (6)

The eigenvalues satisfy the following trace constraint:

$$\sum_{n=1}^{M} \lambda^{(n)} = M. \tag{7}$$

By substituting Eq. (5) into Eq. (3) and comparing it with Eq. (6), we find that

$$\langle a_n(t)a_{n'}(t)\rangle_t = \delta_{nn'}\lambda^{(n)}.$$
(8)

The eigenvalue of each eigenmode thus represents the strength of fluctuations associated with the mode. Figure 4 shows the temporal variation of  $a_n^2(t)$ , which is in sharp contrast to the results shown in Fig. 3. The transformation of the base for describing the IIP data reveals that very few degrees of freedom actually are responsible for the complicated behavior of the IIP.

Thanks to the RMT, we are able to quantify how many eigenmodes should be considered. Probability distribution



PHYSICAL REVIEW E 83, 016103 (2011)



FIG. 5. Probability distribution function  $\rho(\lambda)$  for eigenvalues  $(\lambda)$ of correlation matrix derived from IIP data, in comparison with corresponding result of RMT represented by the solid curve.

function  $\rho(\lambda)$  for the eigenvalues  $(\lambda)$  of the correlation matrix C is shown in Fig. 5. It is compared with the corresponding result [21] of the RMT in the limit of infinite dimensions:

$$\rho(\lambda) = \begin{cases} \frac{Q}{2\pi} \frac{\sqrt{(\lambda_{+} - \lambda)(\lambda - \lambda_{-})}}{\lambda} & \text{for } \lambda_{-} \leq \lambda \leq \lambda_{+}, \\ 0 & \text{otherwise,} \end{cases} \tag{9}$$

where  $Q := N'/M \simeq 3.79$  (> 1) and the upper and lower bounds  $\lambda_\pm$  for  $\lambda$  are given as

$$\lambda_{\pm} = \frac{(1 \pm \sqrt{Q})^2}{Q} \simeq \begin{cases} 2.29, \\ 0.237. \end{cases}$$
(10)

We see that the largest and the second largest eigenvalues, designated as  $\lambda^{(1)}(\simeq 9.95)$  and  $\lambda^{(2)}(\simeq 3.83)$ , are well separated from the eigenvalue distribution predicted by the RMT, whereas the third largest eigenvalue,  $\lambda^{(3)}(\simeq 2.77)$ , is adjacent to the continuum. Therefore only 2 eigenmodes out of a total of 63 are of statistical significance according to the RMT.

Readers may be curious about the present construction of a correlation matrix by mixing up data for production, shipments, and inventory because these are very different species of data at first glance. Thus far, physicists have applied the RMT mainly to analyses of stock data having similar characteristics. In this sense, our approach is quite radical. However, production, shipments, and inventory form a trinity in the economic theory for business cycles so that those variables should be treated on an equal footing. Using the two dominant eigenmodes, in fact, we were successful in proving the existence of intrinsic business cycles [1].

In passing, we note that one may favor the growth rate itself defined by

$$r_{\alpha,g}(t_j) := \frac{S_{\alpha,g}(t_{j+1}) - S_{\alpha,g}(t_j)}{S_{\alpha,g}(t_j)} \tag{11}$$

for the present analysis over the logarithmic growth rate (1). If the relative change in  $S_{\alpha,g}(t_j)$  is small, we need not FIG. 4. Temporal variation of strength of fluctuations associated distinguish between Eqs. (1) and (11) numerically. To confirm that the results obtained here are insensitive to the choice of stochastic variables, we repeated the same calculation by using Eq. (11) and found no appreciable difference between the two calculations for the dominant eigenvalues and their associated eigenvectors. For instance, the first three largest eigenvalues 9.95, 3.83, and 2.77, as shown in Fig. 5, are replaced with 9.96, 3.73, and 2.78, respectively.

### **III. ROTATIONAL RANDOM SHUFFLING**

There are two major sources of noise in the IIP data. One of them, corresponding to thermal noise in physical systems, arises from elimination of a large number of degrees of freedom from our scope as hidden variables. This highlights the stochastic nature of the IIP and has a strong influence on autocorrelation of all goods. The other source of noise originates from the finite length of time series data. Such statistical noise hinders the detection of correlations among different goods in the IIP data. If one could have data of infinite length, statistical noise would disappear in the mutual correlations, and only thermal noise would remain. These two types of noise should be distinguished conceptually. The RMT is an effective tool for eliminating statistical noise from raw data to extract genuine mutual correlations.

However, the noise reduction method based on the RMT heavily depends on the following assumption: Stochastic variables would be totally independent if correlations between different variables were switched off. Such a null hypothesis simultaneously excludes both autocorrelations and mutual correlations. In the case of daily changes in Japanese stock prices that were available [17] to us, we found no detectable autocorrelations in the corresponding variables; therefore the RMT functions ideally. In contrast, the IIP data have significant autocorrelations, as shown in Fig. 6, where the autocorrelation function  $R_{\alpha,g}(t)$  of the normalized growth rate  $w_{\alpha,g}$  is defined as

$$R_{\alpha,g}(t_m) := \frac{1}{N' - m} \sum_{j=1}^{N' - m} w_{\alpha,g}(t_j) w_{\alpha,g}(t_{j+m}).$$
(12)

By definition,  $R_{\alpha,g}(0) = 1$ , and if there are no autocorrelations,  $R_{\alpha,g}(t_m) = 0$  for  $m \ge 1$ . We observe that both production  $(\alpha = 1)$  and shipments  $(\alpha = 2)$  have nontrivial values of autocorrelations at t = 1 month, whereas there is no clear evidence for autocorrelations for inventory  $(\alpha = 3)$  in the same time interval; the values averaged over 21 goods are  $R_1(1) \simeq -0.31$ ,  $R_2(1) \simeq -0.39$ , and  $R_3(1) \simeq 0.007$ . Beyond the 1-month time lag, however, we find no appreciable autocorrelations for any of these three categories.

To formulate the null hypothesis of the RMT using actual data, one may shuffle the IIP data completely in the time direction. In fact, the eigenvalue distribution of the resulting correlation matrix reduces to that of the RMT, as demonstrated in Fig. 7(a), where a total of  $10^5$  samples were generated. Departure from the RMT owing to finiteness of the data size is almost negligible even for such small-scale data as the IIP. This randomization process inevitably destroys both autocorrelations and mutual correlations. From a methodological point of view, it is favorable to deal with these two types of correlations separately.



FIG. 6. Autocorrelation functions  $R_{\alpha,g}(t)$  of production ( $\alpha = 1$ ), shipments ( $\alpha = 2$ ), and inventory ( $\alpha = 3$ ) for each of the goods (g = 1, 2, ..., 21) at (a) t = 1, (b) t = 2, and (c) t = 3, where the index  $\ell$ on the horizontal axis is defined in the same way as in Fig. 3. The 95% confidence level for no autocorrelations is represented by the shaded band in each panel.

We instead propose to shuffle the data rotationally in the time direction, imposing the following periodic boundary condition on each of the time series:

$$w_{\alpha,g}(t_j) \to w_{\alpha,g}[t_{\mathrm{Mod}(j-\tau,N')}],$$
 (13)

where  $\tau \in [0, N' - 1]$  is a (pseudo-) random integer and is different for each  $\alpha$  and g. This randomization destroys only the mutual correlations involved in the data, with the autocorrelations left as they are; therefore it provides us with a null hypothesis more appropriate than that of the RMT.



FIG. 7. Same as Fig. 5, but for eigenvalues ( $\lambda$ ) of the correlation matrix obtained by shuffling IIP data (a) completely or (b) rotationally in the time direction. It should be noted that the autocorrelations involved in the IIP data are preserved in rotational shuffling.

Figure 7(b) shows the result in rotational shuffling with the same number of samples as that in the complete shuffling. We find that the existence of autocorrelations alone leads to departure from the RMT. The third largest eigenvalue  $\lambda^{(3)} \simeq 2.77$  becomes even closer to the upper limit,  $\lambda'_{+} = 2.47 \pm 0.20$ , of the eigenvalues obtained on the basis of the alternative null hypothesis, where the error is estimated at 95% confidence level. This result reinforces neglect of the third eigenmode by the RMT.

Thus this new method for data shuffling conceptually clarifies noise elimination for the correlation matrix, although the difference in the eigenvalue distribution from that of the RMT is practically not very dramatic. In addition, we note that the rotational shuffling of the stock price data in Japan reproduces the RMT result quite well, as is expected from the fact that no appreciable autocorrelations are observed there.

## IV. GENUINE CORRELATION MATRIX

In the current system of IIP data, the preceding careful arguments permit us to adopt

$$\boldsymbol{C}^{(G)} := \sum_{n=1}^{2} \lambda^{(n)} \boldsymbol{V}^{(n)} \boldsymbol{V}^{(n)T} + [\text{diagonal terms}] \qquad (14)$$

as a genuine correlation matrix, which consists of just the first and second eigenvector components in the spectral



FIG. 8. Eigenvector components corresponding to (a) largest and (b) second largest eigenvalues for correlation matrix of IIP time series data.

representation (6) of C plus the diagonal terms, thereby ensuring that all the diagonal components are 1. We note that self-correlations of stochastic variables always exist even if they are merely noise. The components of C are explicitly written as

$$C_{\ell m}^{(G)} = \begin{cases} 1 & \text{for } \ell = m, \\ \sum_{n=1}^{2} \lambda^{(n)} V_{\ell}^{(n)} V_{m}^{(n)} & \text{otherwise.} \end{cases}$$
(15)

The eigenvectors  $V^{(1)}$  and  $V^{(2)}$ , associated with  $\lambda^{(1)}$  and  $\lambda^{(2)}$ , are shown in Figs. 8(a) and 8(b), respectively. These two eigenvectors have characteristic features that distinguish them from each other. The eigenvector  $V^{(1)}$  represents an economic mode in which production and shipments of all goods expand (shrink) synchronously with decreasing (increasing) inventory of producer goods. This corresponds to the market mode obtained for the largest eigenvalue in the stock market analyses [9,10] and may be referred to as the "aggregate demand" mode, according to Keynes's principle of effective demand: Both shipments and production in all the sectors are moved jointly by aggregate demand [22]. On the other hand, the eigenvector  $V^{(2)}$  is a mode that apparently represents dynamics of inventory, that is, accumulation or clearance of inventory, for most goods, including producer goods. We further find positive correlation between production enhancement and inventory accumulation for most goods. This finding indicates that production has a kind of inertia in its response to change of demands.

Accordingly, we project raw fluctuations of  $w_{\ell}(t_j)$  onto the first and second eigenmodes; that is, only the first two terms are retained, and the remaining terms are regarded as just noise in expansion (5):

$$w_{\ell}(t_j) = \sum_{n=1}^{2} a_n(t_j) \, V_{\ell}^{(n)} + [\text{noise}].$$
(16)

This process extracts statistically meaningful information on mutual correlations among  $w_{\ell}(t_j)$ , as has been already discussed in Secs. II and III. Collaboration of these two modes results in inherent business cycles with periods of 40 and 60 months throughout the economy. The cycles are accounted for by time lags in information flow between demand of goods by consumers and decision making of firms on production [1]; inventory fills this information gap. If we were to single out the most dominant mode alone in Eq. (5), all  $w_{\ell}(t_j)$  would oscillate without phase difference. As will be shown later, each good possesses its own characteristics in the phase relations among production, shipments, and inventory.

Temporal change of the two principal factors  $a_1(t)$  and  $a_2(t)$  is plotted in Fig. 9(a). Since the functional behavior of these variables is very noisy, we take their simple moving average, defined as

$$\overline{a_n(t_j)} := \frac{1}{2\xi + 1} \sum_{k=-\xi}^{\xi} a_n(t_{j+k}), \tag{17}$$

where  $\xi$  is a characteristic time scale for smoothing. This process eliminates thermal noise present in the original data. Actually, the moving average was taken with  $\xi = 6$ ; the results for  $\overline{a_1(t)}$  and  $\overline{a_2(t)}$  are shown in Fig. 9(b). We see that the moving-average operation significantly reduces the level of noise present in  $a_n(t)$ . It is noteworthy that Fig. 9 indicates the existence of some mechanical relationship between  $a_1(t)$ and  $a_2(t)$ . This finding is ascertained more quantitatively from Fig. 10, in which the correlation coefficient between  $\overline{a_1(t)}$  and  $\overline{a_2(t-\tau)}$ ,

$$C_{\overline{a_1} \ \overline{a_2}}(\tau) = \frac{\langle \overline{a_1(t)} \ \overline{a_2(t-\tau)} \rangle_t}{\sqrt{\langle \overline{a_1(t)}^2 \rangle_t \ \langle \overline{a_2(t)}^2 \rangle_t}},\tag{18}$$

is plotted as a function of time lag  $\tau$ . A correlation as large as 0.7 is detected between the two dominant modes around  $\tau = 10$  months. Detailed study of the underlying dynamics in the economic system is in progress and will be reported elsewhere.

#### V. FD THEORY

The FD theory plays a central role in nonequilibrium statistical mechanics because this theory establishes a general relation between fluctuation properties of a physical system in equilibrium and response properties of the system to small external perturbations. We assume that the theory still is applicable to the economic system under study here. This assumption provides us with a framework to derive inputoutput interindustrial relations in the system. Its validity in view of how the system responded to the recent economic crisis will be discussed later.



FIG. 9. Two principal factors  $a_1$  and  $a_2$  as a function of time t. (a) Originally obtained results; (b) results smoothed by the movingaverage operation (17) with  $\xi = 6$ . The eigenmodes for business fluctuations were determined using the data in the shaded areas.

Let us denote our variable  $w_{\alpha,g}$  as  $w_{\ell}$ , whose index  $\ell := 21(\alpha - 1) + g$  runs from 1 to 63. We assume that  $w_{\ell}$  obeys dynamics governed by a Hamiltonian  $H(\{w\}, \{x\})$ , where  $\{x\}$  is a set of hidden variables in the system, which encompass all variables in the current economics. These numerous variables interact with each other in a nonlinear chaotic way, and hence the temporal change of  $w_{\ell}$  appears stochastic in the same way as that of a Brownian particle. The existence of underlying dynamics in the IIP [1], as demonstrated in Figs. 9 and 10, strongly supports this idea borrowed from mechanics of motion.

Actually, however, the economy of a nation is quite open now; therefore it could potentially be subjected to perturbations such as disasters, political issues, and trade issues. We thus add external forces  $\epsilon_{\ell}(t)$  to the system; then the total Hamiltonian  $\mathcal{H}$  becomes

$$\mathcal{H} = H(\lbrace w \rbrace, \lbrace x \rbrace) - \sum_{\ell=1}^{M} \epsilon_{\ell}(t) w_{\ell}.$$
 (19)



FIG. 10. Correlation coefficient  $C_{\overline{a_1} \overline{a_2}}(\tau)$  between  $\overline{a_1(t)}$  and  $\overline{a_2(t-\tau)}$  as a function of time lag  $\tau$ , calculated during the normal period (from January 1988 to December 2007) with  $\xi = 6$  for the moving-average operation. A comparison of this result with that obtained using unsmoothed data with  $\xi = 0$  is also shown in the figure.

This extra term represents external perturbations to the equation of motion for  $w_{\ell}$ :

$$\frac{dp_{\ell}}{dt} = -\frac{\partial H}{\partial w_{\ell}} + \epsilon_{\ell}(t), \qquad (20)$$

$$\frac{dw_{\ell}}{dt} = \frac{\partial H}{\partial p_{\ell}},\tag{21}$$

where  $p_{\ell}$  is the momentum conjugate to  $w_{\ell}$ . Therefore  $\epsilon_{\ell}(t)$  directly affects  $w_{\ell}$  at time t, the effect of which then extends to other ws through direct and indirect interactions among them.

For simplicity, let us assume that  $\{\epsilon\}$  is constant in time. Thus the perturbation set  $\{\epsilon\}$  induces a static shift of the equilibrium positions of the variables  $\{w\}$ , which otherwise move stochastically around the origin  $w_{\ell} = 0$ . If the perturbation is weak, the shift  $\langle w_{\ell} \rangle$  thus induced can be expressed by the following linear response relation:

$$\langle w_{\ell} \rangle = \sum_{m=1}^{M} \chi_{\ell m} \epsilon_m, \qquad (22)$$

where the ensemble average denoted by  $\langle \cdot \rangle$  has replaced the time average. The coefficients { $\chi$ } are the result of the interactions, and they are called magnetic susceptibility while describing the physics of magnetic materials.

Once such a set of susceptibilities is available, we can quantify the response of the economic system to external perturbations. For instance, suppose that the government adopts an economic policy to increase the shipment of one of the final demand goods by  $\langle w_m \rangle$  with a stimulus  $\epsilon_m$ . The resulting changes in production, shipments, and inventory of goods are given as

Since  $\epsilon_m$  is not an observable quantity, it should be appropriate to eliminate  $\epsilon_m$  appearing in Eq. (23) and express ripple effects on the economy in terms of  $\langle w_m \rangle$  as

$$\langle w_1 \rangle = \frac{\chi_{1m}}{\chi_{mm}} \langle w_m \rangle,$$
  
$$\vdots$$
  
$$\langle w_M \rangle = \frac{\chi_{Mm}}{\chi_{mm}} \langle w_m \rangle.$$
  
(24)

In Sec. VI, we demonstrate that Eq. (24) can provide quantitative information on input-output interindustrial relations in Japan. Furthermore, one may make reverse use of the linear response relation (22) to distinguish and detect external perturbation from observed economic changes in  $\{w\}$ ; this is discussed in Sec. VII.

Now, the remaining problem is how to calculate  $\{\chi\}$ . To this end, we invoke the concept of the FD theorem in statistical physics. If we assume that the stochastic process of  $\{w\}$  is characterized by Gibbs's ensemble, then the probability density function (pdf)  $P(\{w\},\{\epsilon\})$  for  $\{w\}$  is given as

$$P(\{w\},\{\epsilon\}) \propto \exp[-\beta \mathcal{H}(\{w\},\{\epsilon\})], \tag{25}$$

where the hidden variables  $\{x\}$  have been integrated out and  $\beta$  is the inverse temperature of the economic system. For a weak perturbation, Eq. (25) is expanded to the first order of  $\{\epsilon\}$  as

$$P(\{w\},\{\epsilon\}) \simeq P(\{w\}) \left(1 + \beta \sum_{\ell=1}^{M} \epsilon_{\ell} w_{\ell}\right), \qquad (26)$$

where

$$P(\{w\}) \propto \exp[-\beta H(\{w\})] \tag{27}$$

is the pdf in the absence of  $\{\epsilon\}$ . Equation (26) enables us to calculate the induced change in  $w_{\ell}$  by  $\{\epsilon\}$  as

$$\langle w_{\ell} \rangle = \int P(\{w\}, \{\epsilon\}) w_{\ell} d\{w\}$$

$$\simeq \beta \sum_{m=1}^{M} \epsilon_{m} \int P(\{w\}) w_{\ell} w_{m} d\{w\}$$

$$= \beta \sum_{m=1}^{M} \epsilon_{m} \langle w_{\ell} w_{m} \rangle_{0},$$

$$(28)$$

where  $\langle \cdot \rangle_0$  denotes the ensemble average without perturbation. Comparison of Eqs. (22) and (28) gives one of the outcomes of the FD theorem:

$$\boldsymbol{\chi} = \boldsymbol{\beta} \boldsymbol{C}^{(0)}, \tag{29}$$

where  $C^{(0)}$  denotes a correlation matrix in the absence of external perturbations.

Making use of Eq. (29), we can rewrite the relation (24) of economic ripple effects caused by the increase in shipments of final demand goods as

$$\langle w_1 \rangle = C_{1m}^{(0)} \langle w_m \rangle,$$

$$\vdots \qquad (30)$$

$$\langle w_M \rangle = C_{Mm}^{(0)} \langle w_m \rangle.$$

This is just one example of possible interindustrial relations derived from the present formulation. What we should emphasize here is that Eq. (30) has a rather general form in the framework of linear response.

For instance, we do not need to determine the temperature of the economic system for  $\beta$ . Even the assumption of Gibbs's ensemble, for example, as given in Eqs. (25) and (27), may be too restrictive because the assumption (26) about the pdf is sufficient to derive Eq. (30). We also recall Onsager's regression hypothesis [23,24], on which the FD theorem relies. Once one accepts the hypothesis, one can readily derive Eq. (30). According to Onsager, the response of a system in equilibrium to an external field shares an identical law with its response to a spontaneous fluctuation. In other words, the regression of spontaneous fluctuations at equilibrium takes place in the same way as the relaxation of nonequilibrium disturbances does. Let us suppose that the nonequilibrium disturbances  $\langle w_i \rangle$  and  $\langle w_m \rangle$  are linearly related through

$$\langle w_i \rangle = \kappa \langle w_m \rangle. \tag{31}$$

Accordingly, the spontaneous fluctuations  $w_i$  and  $w_m$  satisfy the same relation as Eq. (31):

$$w_i = \kappa w_m. \tag{32}$$

The ensemble average of Eq. (32) multiplied by  $w_m$  on both sides determines the proportionality coefficient  $\kappa$  as

$$\kappa = C_{im}^{(0)}.\tag{33}$$

We thus see that Eq. (30) is directly derivable from Onsager's hypothesis.

We note that the correlation matrix appearing in Eqs. (29) and (30) should be measured for a system not subject to any perturbations. However, the genuine correlation matrix  $C_{\ell m}^{(G)}$  determined by Eq. (15) is possibly contaminated with various kinds of external economic shocks. While such forces may easily affect the stochastic motion of each w, it is legitimate to assume that their influence on the correlations among ws is much weaker; otherwise, the external factors would have to work coherently to change springs connecting pairs of ws. This consideration justifies the replacement of  $C_{\ell m}^{(0)}$  in Eq. (30) with  $C_{\ell m}^{(G)}$ .

#### VI. INTERINDUSTRIAL RELATIONS

We are now in a position to quantitatively estimate the strength of the interindustrial relations by making use of the genuine correlation matrix through Eq. (30). In particular, we focus on ripple effects on production of intermediate goods that are triggered by applying an external stimulus to consumption of final demand goods:

$$\langle w_{1,20} \rangle = C_{1,20;2,g}^{(G)} \langle w_{2,g} \rangle,$$
 (34)

$$\langle w_{1,21} \rangle = C_{1,21;2,g}^{(G)} \langle w_{2,g} \rangle,$$
 (35)

with g = 1, 2, ..., 19.

The results for production of intermediate goods for mining and manufacturing (g = 20) are shown in Fig. 11 (top), in which those obtained with the original correlation matrix are also added for comparison. This figure shows



FIG. 11. Input-output interindustrial relations based on the genuine and original correlation matrices. (a) Vertical axis indicates the extent to which the logarithmic growth rate of production of intermediate goods for mining and manufacturing (g = 20) is accordingly elevated when the logarithmic growth rate of shipments of each of the final demand goods specified on the horizontal axis is increased by one unit. (b) Same as (a), except showing the relationship between production of intermediate goods for others (g = 21) and shipments of each of the final demand goods.

an increase in the logarithmic growth rate of production of intermediate goods that is predicted from unit increments of the logarithmic growth rate of shipments of each of the final demand goods. As expected, increase in shipments of final demand goods with large weights, as represented by manufacturing equipment (g = 1), construction (g = 9), motor vehicles (g = 15), housework (g = 16), and food and beverage (g = 19), certainly causes large ripple effects in the production of intermediate goods. If the original correlation matrix is replaced with the genuine one, then the relative importance of species of final goods is interchanged between construction and motor vehicles. This is understandable because sales of cars are sometimes promoted just for inventory adjustment, having no effect on the growth of production of intermediate goods. We also note that the original correlation matrix significantly underestimates the effects of furniture and furnishing (g = 13) and nondurable consumer goods (g = 17, 18, 19). It is noteworthy that furniture and furnishing and clothing and footwear, having much smaller weights than the major final demand goods, have comparable contributions; some feedback mechanism must be working through the inner loop in the economic system.

Figure 11 (bottom) shows the corresponding results for production of intermediate goods for others (g = 21), whose

TABLE II. Phases of periodic oscillations with T = 60 and 40 of production (P), shipments (S), and inventory (I) for all goods, in the unit of degrees ranging from -180 to 180. They are measured relative to the production of g = 20.

	T = 60			T = 40		
Goods	Р	S	Ι	Р	S	Ι
1	-8.5	-8.2	-60.3	-9.8	-9.5	-117.7
2	-22.2	-45.3	-75.8	-35.4	-95.4	-130.4
3	-43.8	-40.9	-81.3	-92.2	-85.7	-133.6
4	-16.0	26.8	-85.6	-22.1	17.6	-135.9
5	4.4	61.6	-60.5	4.0	31.0	-117.8
6	2.5	15.8	-42.0	2.4	11.9	-88.2
7	-24.2	-7.0	-33.1	-40.3	-7.8	-65.1
8	-15.6	-4.3	-58.9	-21.3	-4.6	-116.1
9	13.8	37.0	-80.0	10.7	22.0	-132.9
10	56.2	28.4	-93.5	29.1	18.3	-139.6
11	9.2	83.8	-56.6	7.6	39.3	-113.4
12	-10.2	105.8	-54.3	-12.3	50.7	-110.5
13	7.8	48.1	-63.2	6.6	26.2	-120.6
14	-3.8	-4.2	-396.6	-4.0	-4.5	-74.6
15	1.2	-0.7	-24.2	1.1	-0.7	-40.3
16	9.7	27.2	-78.5	8.0	17.8	-132.1
17	11.6	35.4	-80.9	9.3	21.3	-133.4
18	-15.6	6.0	-66.1	-21.2	5.3	-123.2
19	44.6	55.2	-73.0	24.9	28.7	-128.6
20	0	15.8	-89.5	0	11.9	-137.8
21	7.4	40.7	-89.9	6.3	23.5	-138.0
Average	-0.6	25.2	-66.6	-0.6	16.9	-123.7

weight is 1 order of magnitude smaller than that of mining and manufacturing. The important species of final demand goods are common in both categories of intermediate goods. In contrast, the original correlation matrix significantly underestimates the effects of different final demand goods such as those given by g = 3 to g = 7.

Presence of a correlation between two stochastic variables A and B does not indicate the existence of a mechanical connection between them. Actually, correlating A and B might be driven by a third variable C; then there would be no causality relationship between A and B. To address this question, we provide detailed information on phase relations in the business cycles identified in the previous study [1]. Table II lists phases of the cyclic motion of production, shipments, and inventory at T = 60 and 40 for each of the goods. We can see that shipments of final demand goods are ahead of or almost in phase with production of intermediate goods; the resolution limit (1 month) is  $6^{\circ}$  and  $9^{\circ}$  for T = 60 and 40, respectively. Electricity (g = 2) and communication and broadcasting (g = 3)are exceptions to this observation. The wave of production arrives first and then that of shipments follows for electricity; the production activity for communication and broadcasting behaves significantly out of phase with that averaged over goods. Since we have adopted a static approximation for the interindustrial relations, it may be more appropriate to average the phase relations over frequency. The results are shown in Fig. 12 and Table III. The frequency-averaged phase relations in the cyclic behavior of the economic fluctuations thus support our postulate that production of intermediate goods is driven



FIG. 12. Frequency-averaged phases of production, shipments, and inventory for each type of goods, measured relatively to production of g = 20 in the unit of degrees.

by increasing shipments of final demand goods, with a few exceptions.

### VII. EXTERNAL STIMULI

Finally, we try to identify the presence of external stimuli hidden in real data by inversely using the linear response relationship (22). The recent global economic crisis certainly has delivered an extremely large shock to the economic system of Japan, as is clearly shown in Fig. 2. In our previous paper [1], however, we demonstrated that the crisis has simply increased the level of fluctuations associated with the dominant modes that were determined from the data during the normal time, instead of destroying the industrial structure itself; this is also manifested here, as shown in Fig. 9. And the information on collective movement of the IIP that we could extract from the dominant modes remains intact even in such an abnormal situation. This result thus conforms to the idea of Onsager's regression hypothesis, indicating the validity of the FD theory even in an economic system that is supposed to be far away from equilibrium.

Since approximation (15) has been adopted for the correlation matrix, we consider only two independent external fields  $\{\eta_1, \eta_2\}$  that are coupled to the normal coordinates  $\{a_1, a_2\}$ associated with the two dominant eigenmodes  $V^{(1)}$  and  $V^{(2)}$ , respectively. The total Hamiltonian [Eq. (19)] is therefore simplified to

$$\mathcal{H} = H(a_1, a_2, \{x\}) - \eta_1 a_1 - \eta_2 a_2.$$
(36)

TABLE III. Frequency-averaged phases of periodic motion of production (P), shipments (S), and inventory (I) for final demand and producer goods in the unit of degrees. The results for final demand goods were obtained by averaging over goods excluding g = 2 and 3; the numbers in parentheses are those obtained with all final demand goods.

		Produce	Producer goods		
	Final demand goods	g = 20	g = 21		
Р	-0.69 (-2.35)	0	0.99		
S	2.99 (0.20)	5.84	5.15		
I	-28.7 (-29.5)	-37.4	-37.5		

The reduced external fields  $\{\eta\}$  in Eq. (36) are derived from the original ones in Eq. (19) through

$$\eta_n = \sum_{\ell=1}^M \epsilon_\ell V_\ell^{(n)}.$$
(37)

Then Eq. (22) is projected onto the two-dimensional reduced state space as

$$\begin{pmatrix} \langle a_1 \rangle \\ \langle a_2 \rangle \end{pmatrix} = \begin{pmatrix} \hat{\chi}_{11} & \hat{\chi}_{12} \\ \hat{\chi}_{21} & \hat{\chi}_{22} \end{pmatrix} \begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix}, \tag{38}$$

where

$$\langle a_n \rangle = \sum_{\ell=1}^M \langle w_\ell \rangle \, V_\ell^{(n)},\tag{39}$$

and the reduced susceptibilities  $\{\hat{\chi}\}$  are defined as

$$\hat{\chi}_{mn} = \sum_{i=1}^{M} \sum_{j=1}^{M} V_i^{(m)} \chi_{ij} V_j^{(n)}.$$
(40)

The relative values of  $\{\hat{\chi}\}$  with reference to  $\hat{\chi}_{11}$  are calculated from  $C^{(G)}$  as

$$\begin{pmatrix} \hat{\chi}_{11} & \hat{\chi}_{12} \\ \hat{\chi}_{21} & \hat{\chi}_{22} \end{pmatrix} = \beta \begin{pmatrix} 1 & 1.30 \times 10^{-3} \\ 1.30 \times 10^{-3} & 0.433 \end{pmatrix}.$$
 (41)

These results shows that the two eigenmodes are almost decoupled from each other, which is understandable from the orthogonality [Eq. (8)] of the normal coordinates.

One can obtain  $\{\eta\}$  using the inverse of Eq. (38) along with Eqs. (39) and (40), although it is not so straightforward. We first recall that  $\langle w_{\ell} \rangle$  on the right-hand side of Eq. (39) is the deviation of  $w_{\ell}$  from the equilibrium value induced by external perturbation and not fluctuations of  $w_{\ell}$  directly observed in the real data. We then identify  $\langle w_{\ell} \rangle$  as residuals obtained by subtracting the long-period components arising from the inherent business cycles from moving-average fluctuations of the IIP.

To extract  $\langle w_{\ell} \rangle$ , we first define the Fourier transform of the coefficients  $a_n(t_j)$  as

$$a_n(t_j) = \frac{1}{\sqrt{N'}} \sum_{k=1}^{N'-1} \widetilde{a}_n(\omega_k) e^{-i\omega_k t_j}, \qquad (42)$$

with the Fourier frequency  $\omega_k = 2\pi k/(N'\Delta t)$  and hence  $\omega_k t_j = 2\pi k j/N'$ . The relevant long-period component  $a_n^{(LP)}(t_j)$  is obtained by limiting the sum over k only to k = 1 (T = 240), k = 2 (T = 120), k = 4 (T = 60), and k = 6 (T = 40) or by summing all of the terms with periods larger than 2 years ( $k \leq 9$ ) in Eq. (42). The formula for  $\langle w_\ell \rangle$  is finally expanded as

$$\langle w_{\ell}(t) \rangle = \sum_{n=1}^{2} \overline{a_n(t)} \, V_{\alpha,g}^{(n)} - \sum_{n=1}^{2} a_n^{(\text{LP})}(t) \, V_{\alpha,g}^{(n)}.$$
 (43)

Figure 13, for which we arbitrarily set  $\beta = 1$  in Eq. (41), shows the external fields  $\eta_1$  and  $\eta_2$  thus derived from  $\langle w_\ell \rangle$ . Two computational schemes were adopted to evaluate the long-period components in the IIP data, and no appreciable difference was observed between the two results. Here the economic system was assumed to respond instantaneously to



FIG. 13. External stimuli  $\eta_1$  and  $\eta_2$ , derived from IIP data through the linear response relation (38), shown as a function of time in (a) and (b), respectively; the system is assumed to respond instantly to the applied external fields. The solid curves depict results obtained only with the terms of k = 1, 2, 4, and 6 in Eq. (42), and the dotted curves depict those calculated with the terms of  $k \leq 9$ . The shaded area is the same as depicted in Figs. 2 and 9.

the applied external fields without any time delay. Referring to Fig. 2, we clearly confirm that such a large external shock as manifested in  $\eta_1$  causes the drastic drop in industrial activities in Japan. We also see that another large shock in  $\eta_2$ , which leads to reduction in inventory, immediately accompanies the first shock. In contrast, the maximum fluctuation levels of  $\eta_1$  and  $\eta_2$  are 0.1 and 0.2, respectively, in the normal period (before the end of 2007).

### VIII. CONCLUSION

This paper has described our attempt to utilize the FD theory for elucidating the nature of input-output correlations in the Japanese industry on the basis of IIP data. We were able to quantitatively estimate the strength of correlations between goods by using the genuine correlation matrix obtained in this study. We were also successful in extracting external stimuli over the last two decades. The noise reduction along with the RMT enabled us to detect economic signals hidden behind the complicated dynamics of the IIP. The strong coincidence between the sudden change in IIP data and the external shocks described here may prove that the present method is capable of predicting the input-output interindustrial relationship with a much higher time resolution than the annual resolution. We thus expect the results of this study to provide a new HIROSHI IYETOMI et al.

methodology for gaining deeper understanding of complex economic phenomena at a macroscopic level.

#### ACKNOWLEDGMENTS

The present study was supported in part by the Program for Promoting Methodological Innovation in Humanities and

- H. Iyetomi, Y. Nakayama, H. Yoshikawa, H. Aoyama, Y. Fujiwara, Y. Ikeda, and W. Souma, e-print arXiv:0912.0857 (to be published).
- [2] W. Leontief, Rev. Econ. Statist. 18, 105 (1986).
- [3] W. Leontief, *Input-Output Economics*, 2nd ed. (Oxford University Press, New York, 1986).
- [4] L. Walras, *Elements of Pure Economics* (Allen and Unwin, London, 1954).
- [5] R. Mantegna and H. Stanley, An Introduction to Econophysics: Correlations and Complexity in Finance (Cambridge University Press, Cambridge, 2000).
- [6] J. Bouchaud and M. Potters, *Theory of Financial Risks: From Statistical Physics to Risk Management*, 2nd ed. (Cambridge University Press, Cambridge, 2003).
- [7] M. Aoki and H. Yoshikawa, Reconstructing Macroeconomics: A Perspective from Statistical Physics and Combinatorial Stochastic Processes (Cambridge University Press, New York, 2007).
- [8] H. Aoyama, Y. Fujiwara, Y. Ikeda, H. Iyetomi, and W. Souma, *Econophysics and Companies: Statistical Life and Death in Complex Business Networks* (Cambridge University Press, Cambridge, 2010).
- [9] L. Laloux, P. Cizeau, J. P. Bouchaud, and M. Potters, Phys. Rev. Lett. 83, 1467 (1999).

Social Sciences by Cross-Disciplinary Fusing of the Japan Society for the Promotion of Science and by the Ministry of Education, Science, Sports, and Culture, Grants-in-Aid for Scientific Research (B), Grant No. 20330060 (2008-10) and No. 22300080 (2010-12). We would also like to thank Hiroshi Yoshikawa for providing continual advice and encouragement.

- [10] V. Plerou, P. Gopikrishnan, B. Rosenow, L. A. N. Amaral, and H. E. Stanley, Phys. Rev. Lett. 83, 1471 (1999).
- [11] V. Plerou, P. Gopikrishnan, B. Rosenow, L. A. N. Amaral, T. Guhr, and H. E. Stanley, Phys. Rev. E 65, 066126 (2002).
- [12] A. Utsugi, K. Ino, and M. Oshikawa, Phys. Rev. E 70, 026110 (2004).
- [13] D. H. Kim and H. Jeong, Phys. Rev. E 72, 046133 (2005).
- [14] V. Kulkarni and N. Deo, Eur. Phys. J. B 60, 101 (2007).
- [15] R. K. Pan and S. Sinha, Phys. Rev. E 76, 046116 (2007).
- [16] J. Shen and B. Zheng, Europhys. Lett. 86, 48005 (2009).
- [17] Y. Nakayama and H. Iyetomi, Prog. Theor. Phys. Suppl. 179, 60 (2009).
- [18] L. D. Landau and E. M. Lifshitz, *Statistical Physics, Part 1*, 3rd ed. (Butterworth-Heinemann, Oxford, 1980).
- [19] R. J. Hawkins and M. Aoki, Econ. Open-Access Open-Assess. E-J. 3, 2009 (2009).
- [20] [http://www.meti.go.jp/english/statistics/tyo/iip/index.html].
- [21] A. M. Sengupta and P. P. Mitra, Phys. Rev. E **60**, 3389 (1999).
- [22] J. M. Keynes, *The General Theory of Employment, Interest and Money* (Macmillan, London, 1936).
- [23] L. Onsager, Phys. Rev. 37, 405 (1931).
- [24] L. Onsager, Phys. Rev. 38, 2265 (1931).