

## Erratum: Critical wetting transitions in two-dimensional systems subject to long-ranged boundary fields [Phys. Rev. E **79**, 041144 (2009)]

A. Drzewiński, A. Maciołek, A. Barasiński, and S. Dietrich  
(Received 5 June 2009; published 7 July 2009)

 DOI: [10.1103/PhysRevE.80.019901](https://doi.org/10.1103/PhysRevE.80.019901)

PACS number(s): 05.50.+q, 68.35.Rh, 68.08.Bc, 99.10.Cd

In our paper the study of interface localization-delocalization (ILD) transitions allows one to infer wetting transitions in the corresponding semi-infinite systems governed by the Hamiltonian [see Eq. (3) in our paper]

$$\mathcal{H} = -J \left( \sum_{\langle kj, k'j' \rangle} \sigma_{k,j} \sigma_{k',j'} + \sum_{j=1}^{\infty} V_j^{\text{ext}} \sum_k \sigma_{k,j} + H \sum_{k,j} \sigma_{k,j} \right), \quad (1)$$

with  $J > 0$  and the external potential  $V_j^{\text{ext}} = \frac{h_1}{j^p}$  with  $p > 0$ . Our numerical data correspond to  $h_1 > 0$  so that at the surface there is a preference for the spin-up phase with magnetization  $= +|m_b|$ , where  $m_b < 0$  is the bulk magnetization corresponding to the bulk field  $H = 0^-$ .

For this Hamiltonian the effective interface potential  $\omega$  at  $T=0$  has the form (see Fig. 1)

$$\frac{\omega(\ell, T=0, H=0^-)}{2J} = \begin{cases} h_1 \sum_{j=\ell+1}^{\infty} \frac{1}{j^p}, & \ell \geq 1 \\ h_1 \zeta(p) - 1, & \ell = 0, \end{cases} \quad (2)$$

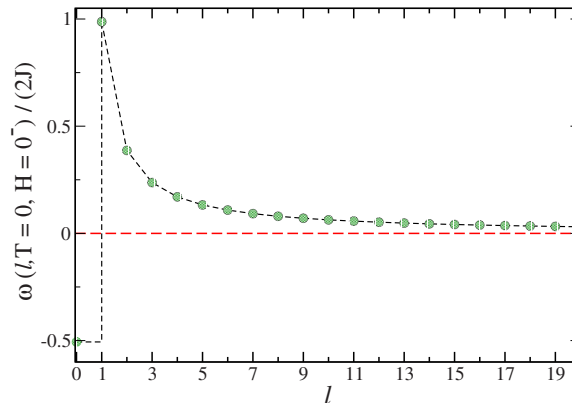


FIG. 1. (Color online) Effective interface potential  $\omega(\ell, T=0, H=0^-)/2J$  [see Eq. (2)] for  $p=2$  and  $h_1=0.3$ .

where  $\zeta(p)$  is the Riemann zeta function; here the thickness  $\ell$  of the wetting film is taken to be the number of surface layers with magnetization  $|m_b|=+1$ . The equilibrium wetting film thickness minimizes  $\omega$ .

At  $T=0$  the wetting transition occurs if  $\omega(\ell=0)=\omega(\ell=\infty)=0$ , i.e., at  $h_1=1/\zeta(p)$ . For  $h_1 < 1/\zeta(p)$  the system is partially wet. This agrees with the condition for the ILD transition in the strip [see Eq. (9) of our paper]. [In the first sum in the last line of Eq. (8) and in Eq. (9) the sum over  $n$  should be from 1 to  $\infty$ .]

Figure 1 shows that our numerical data cannot be compared with the study in Ref. [1] because the latter corresponds to effective interface potentials which are attractive for large  $\ell$ . Accordingly our findings do not challenge the reliability of effective interface models.

We are grateful to A. Parry and N. Bernardino for having alerted us to this issue.

---

[1] D. M. Kroll and R. Lipowsky, Phys. Rev. B **28**, 5273 (1983).