# Transport properties of Na<sup>+</sup> ions in Kr gas

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The transverse-diffusion coefficients for Na<sup>+</sup> ions drifting in Kr gas have been experimentally measured with estimated accuracies of better than 3% at E/N values ranging from 5 to 500 Td (1 Td=10<sup>-21</sup> V m<sup>2</sup>) at about 303 K. In addition, elaborate calculations of the transport properties for the Na<sup>+</sup>-Kr system have been made by using Monte Carlo simulations (MCS's) and the generalized Einstein-relation techniques. The experimental results are compared with values obtained from MCS's using the interaction potential of Koutselos, Mason, and Viehland [J. Chem. Phys. 93, 7125 (1990)] as input. Agreement within 3% for the entire range of E/N is observed in the comparisons. MCS-calculated values of longitudinal-diffusion coefficients are also compared with those determined experimentally by Thackston *et al.* [J. Chem. Phys. 73, 2012 (1980)].

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# INTRODUCTION

Accurate experimental determination of the longitudinal- and transverse-diffusion coefficients  $(D_L)$ and  $D_T$ , respectively) and mobility K of ions in a nonreacting neutral gas is of considerable importance to theoretical investigations of their mutual interaction forces [1-3]. Experimental data are often used as input in an inversion which leads to the calculation of the medium- to long-range interaction potential of the ion-neutral-atom pair. There now exists an extensive compilation of data on mobilities and longitudinaldiffusion coefficients [4-6], but little on transversediffusion coefficients [7-12]. The lack of data reported on  $D_T$  is basically due to some experimental limitations [2]. Because of its theoretical simplicity in analysis [13], there is a need for more experimental data on  $D_T$ .

Mobilities and values of  $D_L$  for Na<sup>+</sup> ions in Kr have been experimentally determined by Thackston *et al.* [14] and compiled by Ellis *et al.* [6]. This seems to be the only experimental data reported on Na<sup>+</sup>-Kr. To our knowledge, there has hitherto been no published data on the transverse-diffusion coefficient of Na<sup>+</sup>-Kr. Such studies of Na<sup>+</sup> ions in Ne and Ar have been reported [15].

In this paper, the experimental results of  $D_T/K$ , the ratio of the transverse-diffusion coefficient to mobility for Na<sup>+</sup> ions drifting in Kr gas at values of E/N (electricfield to neutral-gas density ratio) ranging from 5 to 500 Td (1Td=10<sup>-21</sup> V m<sup>2</sup>) at 303±1 K, are presented. Results of the calculations of  $D_T/K$  and  $D_L/K$  using Monte Carlo simulations (MCS's) and the generalized Einsteinrelation (GER's) techniques are also reported.

## **EXPERIMENTAL METHOD**

The measurements were made on a newly developed drift tube which was constructed for the study of transverse diffusion of ions in neutral gas at room temperature. The apparatus has been described in detail elsewhere [16] and the procedure for analyzing the results is the same as the one used in a previous investigation [15].

Systematic checks were made to ensure the nondependence of the data on gas pressure, drift length, conditions of ion injection into the drift tube, and ion-space-charge effects. In particular, repeated measurements made at five different drift lengths for E/N=70, 110, and 200 Td revealed no significant change in results, thus confirming that drift-tube end effects are negligible.

In the present measurements, sodium aluminosilicates were prepared and used as the thermionic emitters of Na<sup>+</sup> ions, according to the method described by Hogan et al. [17]. In all cases, the ions present were always more than 99% Na<sup>+</sup>. However, it should be noted that during the initial burn-in, the source produced a fairly large quantity (up to 25%) of the impurity  $K^+$  ions. As expected, the impurity  $K^+$  ions decayed rapidly during burn-in until only trace amounts were present. The drift chamber was always flushed and evacuated at least once with Kr gas before final introduction of the experimental gas. A residual gas analyzer, which was used to identify the neutral species present, found negligible traces of impurities, mostly residual water vapor, and Kr gas in its natural isotropic abundance. The Kr gas used in the experiments had a purity content of 99.999%.

## RESULTS

Experimental values of  $D_T/K$  of Na<sup>+</sup> ions in Kr were obtained from the transverse-ion-current-density spectrum for E/N values from 5-500 Td. The lower limit of E/N was set by the spectrum spread which became too broad and extended into the boundary of the drift tube. The geometry of the drift tube and electrical breakdown occurring when the drift tube voltage exceeded 300 V set the upper limit of E/N to 500 Td. Results were recorded over the pressure range of 13.3-26.7 Pa at temperatures between 302 and 304 K. To facilitate comparison, all data are adjusted to a fixed temperature of 303 K using the first-order correction formula [18], given by

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$$(D_T/K)[E/N, T + \Delta T] = (D_T/K)[E/N, T] + k_B \Delta T/e$$
,  
(1)

where the symbols are defined in Ref. [18]. The square brackets imply that the ratio  $D_T/K$  is a function of E/N and  $T + \Delta T$ .

Detailed experimental results are presented in Table I. Each value of  $D_T/K$  is the average of at least four independent measurements. Also given is the sample standard deviation associated with the averaged value. Judging from the magnitudes of the standard deviations in Table I, it is estimated that the random error in the values reported for  $D_T/K$  is less than  $\pm 2\%$ . The major contributions to the systematic error are in the pressure measurements ( $\pm 0.3\%$  or less), in the knowledge of the

TABLE I. Experimentally measured values of  $D_T/K$  for Na<sup>+</sup> in Kr with their associated standard deviations and reduced transverse-diffusion coefficient to mobility ratios. The  $D_T/K$  values have been adjusted to 303 K.

E/N	Average $D_T/K$	Standard deviation			calculated values using from Thackston at all			
(Td)	(mV)	(mV)	(%)	$D_T^{(r)}$	Irom I	nacksto	n et al.	
5	26.4	0.31	1.2	1.004				
10	27.2	0.25	0.9	1.005	TAE	BLE II.	MCS c	
15	28.3	0.20	0.7	1.005	and $D_I$	" using	the KN	
20	30.1	0.31	1.0	1.013	E/N	Tr		
25	32.3	0.26	0.8	1.017	(Td)	( <b>К</b> )	$(cm^2)$	
30	34.8	0.22	0.6	1.018				
35	38.9	0.38	1.0	1.048	5	306		
40	43.1	0.02	0.1	1.067	10	315		
45	48.3	0.11	0.2	1.092	15	330		
50	55.7	0.45	0.8	1.149	20	352		
55	64.1	0.07	0.1	1.205	25	379		
60	75.9	0.31	0.4	1.300	30	413		
65	86.5	0.58	0.7	1.353	35	460		
70	101	0.20	0.2	1.446	40	511		
80	128	1.19	0.9	1.538	45	571		
90	168	1.69	1.0	1.700	50	648		
100	208	2.20	1.1	1.795	55	738		
110	244	4.70	1.9	1.817	60	842		
120	285	3.93	1.4	1.835	65	969		
130	329	1.82	0.6	1.852	70	1 1 2 5		
140	378	4.91	1.3	1.873	80	1 460		
150	427	2.43	0.6	1.877	90	1 865		
160	478	6.03	1.3	1.872	100	2 371		
170	530	2.17	0.4	1.862	110	2 891		
180	582	5.83	1.0	1.839	120	3 488		
190	630	7.42	1.2	1.804	140	4 696		
200	678	5.48	0.8	1.765	160	6 0 0 3		
220	769	3.92	0.5	1.673	180	7 281		
240	870	11.79	1.4	1.605	200	8 745		
260	988	7.63	0.8	1.565	230	10 767		
280	1095	12.00	1.1	1.504	260	13 156		
300	1229	8.09	0.7	1.478	300	16215		
340	1500	8.28	0.6	1.413	340	19 401		
380	1741	21.15	1.2	1.320	380	22 7 1 9		
420	2054	11.37	0.6	1.279	420	26 297		
460	2277	20.41	0.9	1.185	460	30 108		
500	2551	32.84	1.3	1.126	500	33 786		



FIG. 1. Reduced transverse-diffusion coefficient to mobility ratios of Na<sup>+</sup> in Kr plotted against E/N. Squares (connected by straight lines) and diamonds represent the present experimental data and MCS-calculated values using the KMV potential, respectively. Triangles and crosses represent GERcalculated values using mobilities from MCS calculations and from Thackston *et al.* [14] respectively.

TABLE II. MCS calculated values of  $K_0$ ,  $D_T/K$ ,  $D_T^{(r)}$ ,  $D_L/K$ , and  $D_L^{(r)}$  using the KMV potential.

E/N	$T_{\rm eff}$	$K_0$	$D_T/K$		$D_L/K$	6.5
(Td)	( <b>K</b> )	$(cm^2 V^{-1} s^{-1})$	(mV)	$D_T^{(r)}$	(mV)	$D_L^{(r)}$
5	306	2.22	26.9	1.021	26.2	0.995
10	315	2.23	27.0	0.999	27.4	1.013
15	330	2.24	27.8	0.988	28.6	1.016
20	352	2.24	29.5	0.993	30.2	1.017
25	379	2.24	32.0	1.008	35.6	1.121
30	413	2.25	35.1	1.026	39.3	1.149
35	460	2.30	38.4	1.036	46.0	1.240
40	511	2.32	42.5	1.051	54.6	1.349
45	571	2.34	48.1	1.086	65.6	1.483
50	648	2.39	54.9	1.131	81.0	1.669
55	738	2.44	64.2	1.206	101	1.895
60	842	2.49	75.2	1.288	119	2.044
65	969	2.55	86.5	1.353	144	2.245
70	1 1 2 5	2.63	98.1	1.400	170	2.431
80	1 460	2.73	127	1.516	223	2.675
90	1 865	2.82	164	1.658	271	2.740
100	2 371	2.92	209	1.806	305	2.632
110	2 891	2.97	251	1.862	350	2.603
120	3 488	3.02	287	1.852	381	2.455
140	4 696	3.04	377	1.868	445	2.205
160	6 0 0 3	3.03	469	1.833	513	2.009
180	7 2 8 1	2.98	584	1.846	590	1.863
200	8 745	2.95	673	1.749	645	1.676
230	10767	2.86	818	1.635	779	1.557
260	13 156	2.80	997	1.577	902	1.427
300	16215	2.70	1222	1.467	1068	1.282
340	19 401	2.61	1485	1.398	1313	1.236
380	22 719	2.53	1741	1.319	1490	1.128
420	26 297	2.47	2017	1.255	1728	1.075
460	30 108	2.41	2270	1.180	1931	1.004
500	33 786	2.35	2540	1.120	2220	0.979

drift length ( $\pm 0.3\%$  or less), and in the approximations ( $\pm 0.3\%$  or less) made in deriving the equation for the ion-current-density spectrum. The total error for the experimental data is estimated to be  $\pm 3\%$ .

Ratios of the transverse- or longitudinal-diffusion coefficient to mobility  $(D_T/K \text{ or } D_L/K, \text{ respectively})$  can be presented effectively [12] in a reduced form given by

$$D_{T,L}^{(r)} = (e/k_B T_{\rm pol})(D_{T,L}/K) , \qquad (2)$$

where

$$T_{\rm pol} = T + M (N_0 K_{\rm pol} E / N)^2 / 3k_B .$$
 (3)

This reduced form removes both the  $(E/N)^2$  and the first-order temperature dependence of  $D_{T,L}/K$  and amplifies the details of the variation of the diffusion coefficients with E/N.

Figure 1 shows a plot of  $D_T^{(r)}$  with E/N for Na<sup>+</sup> ions in Kr.

The Koutselos-Mason-Viehland (KMV) interaction potential [19] of the Na<sup>+</sup>-Kr system, which is the only reliable one available, is used as input in our high-speed MCS calculations. Details of the method have been previously reported [8,12,20] and results are shown in Table II. Confidence in the broad applicability of the KMV potential for the alkali-metal-ion-rare-gas interaction has been built up from several previous investigations involving different interaction pairs [8–12]. Values of  $T_{\rm eff}$ ,  $K_0$ ,  $D_T/K$ ,  $D_T^{(r)}$ ,  $D_L/K$ , and  $D_L^{(r)}$  are presented. Repeated calculations covering 2.5 million collisions show that the reproducibility of the results is better than 1% for mobility and 2.5% for  $D_{T,L}$ . These results are shown graphically in Figs. 1 and 2, as plots of  $D_T^{(r)}$  and  $D_L^{(r)}$ , respectively.

To our knowledge, the only experimentally determined values of  $K_0$  and  $ND_L$  for Na<sup>+</sup> in Kr are those due to Thackston *et al.* [14]. These results were later compiled by Ellis *et al.* [6] and derived quantities of  $T_{\rm eff}$  were determined. Unfortunately, we have found that both the derived values of  $T_{\rm eff}$  and the numerical values of  $ND_L$ 



FIG. 2. Reduced longitudinal-diffusion coefficient to mobility ratios of Na<sup>+</sup> in Kr plotted against E/N. Squares and diamonds represent the experimental data of Thackston *et al.* [14] and MCS-calculated values using the KMV potential, respectively. Triangles and crosses represent GER-calculated values using mobilities from MCS calculations and from Thackston *et al.* [14], respectively.

given by Ellis *et al.* were inconsistent with the original data of Thackston *et al.* Consequently, we decided to calculate the quantities of  $T_{\text{eff}}$  and  $D_L/K$  based on the original data of Thackston *et al.* Values of  $D_L/K$  and  $D_L^{(r)}$  are given in Table III and plotted in Fig. 2.

We have also used Waldman and Mason's parametrized version [21] of the generalized Einstein relations to compute  $D_T/K$  and  $D_L/K$  values from both the mobility data of the MCS calculations and those of Thackston *et al.* [14]. Results are given in Table IV and plots of  $D_T^{(r)}$  and  $D_L^{(r)}$  are included in Figs. 1 and 2, respectively.

TABLE III. Derived values of  $v_d$ ,  $T_{\text{eff}}$ ,  $D_L/K$ , and  $D_L^{(r)}$  from the experimental mobilities  $K_0$  and longitudinal-diffusion coefficients  $ND_L$  of Thackston *et al.* [14].

<i>E / N</i> (Td)	$\frac{ND_L}{(10^{-18} \text{ cm}^{-1} \text{ s}^{-1})}$	$(cm^2 V^{-1} s^{-1})$	$(\mathrm{cm} \mathrm{s}^{-1})$	${T_{ m eff}} ({f K})$	$D_L/K$ (mV)	$D_L^{(r)}$
7	1.53	2.20	4 1 3 8	309	25.9	0.975
8	1.53	2.20	4 729	311	25.9	0.970
10	1.58	2.20	5911	315	26.7	0.990
15	1.68	2.20	8 867	329	28.4	1.010
20	1.83	2.21	11 877	350	30.8	1.038
30	2.54	2.26	18218	414	41.8	1.224
40	3.65	2.31	24 828	510	58.8	1.454
60	9.20	2.52	40 627	857	136	2.327
80	18.2	2.78	59 7 59	1 503	244	2.919
100	27.0	2.98	80 073	2 457	337	2.913
150	45.0	3.12	125 752	5615	537	2.357
200	60.0	3.04	163 370	9 268	735	1.910
250	77.0	2.90	194 808	13 050	988	1.685
300	89.9	2.78	224 096	17 172	1 204	1.445
400	145	2.59	278 373	26 3 32	2 084	1.427
500	193	2.45	329 158	36 696	2932	1.293

## DISCUSSION

Figure 1 shows that the present experimental data of  $D_T^{(r)}$  agree so well with those of the MCS calculations using the KMV interaction potential that the latter may also be represented by almost the same curve. The only slight difference is that the experimental values give a smoother curve around the peak region from E/N=110-170 Td, which thus yields a more accurate value of the peak position of the  $D_T^{(r)}$  plot. The values calculated from the GER using the MCS-derived mobility values as input agree very well with both experimental and the directly calculated MCS values at E/N above 200 Td and at very low E/N. However, these GER values are 4-11 % below them for intermediate E/N values between 40 and 120 Td. GER calculations are not

expected to be accurate when the ion-neutral-atom mass ratio is not negligible and in the E/N region where the mobility rises rapidly. This inaccuracy of the GER is similar to that found in the three-temperature theory [2] from which the GER formulas were derived. At moderate and high values of E/N, there is a significant difference between the GER-calculated values derived from the mobility data of Thackston *et al.* [14] and the rest of the  $D_T^{(r)}$  data. This discrepancy is probably due to a systematic error in the mobility measurements made with the Georgia Tech drift-tube mass spectrometer. Such trends were similarly observed for Na<sup>+</sup> ions in Ar and Ne [15] and others [7,22]. A brief discussion concerning the discrepancy is presented by Hogan and Ong [15].

A comparison of the various  $D_L^{(r)}$  data is shown in Fig.

TABLE IV. GER calculated values of  $D_T/K$ ,  $D_T^{(r)}$ ,  $D_L/K$ , and  $D_L^{(r)}$  using the MCS mobility values, and the experimental mobility data of Thackston *et al.* [14].

	Based on MCS mobility values				Based on experimental mobility values [14]			
E/N	$D_T/K$		$D_T/K$		$D_T/K$		$D_T/K$	
(Td)	(mV)	$D_T^{(r)}$	(mV)	$D_L^{(r)}$	(mV)	$D_T^{(r)}$	(mV)	$D_L^{(r)}$
5	26.4	1.002	27.5	1.044				
7					26.9	1.013	26.7	1.005
8					26.7	1.000	26.9	1.008
10	27.0	0.999	27.7	1.025	27.0	0.999	27.4	1.014
15	28.1	0.999	29.4	1.045	28.1	0.999	29.3	1.041
20	29.7	1.000	31.5	1.061	29.6	0.997	32.5	1.094
25	31.8	1.003	34.8	1.097				
30	34.4	1.006	41.8	1.223	34.5	1.009	41.5	1.214
35	37.9	1.022	49.1	1.324				
40	41.7	1.031	53.2	1.315	41.7	1.031	57.1	1.412
45	46.3	1.046	64.6	1.460				
50	52.2	1.076	80.8	1.666				
55	59.2	1.113	95.3	1.791				
60	67.2	1.152	116	1.988	68.4	1.172	119	2.032
65	77.2	1.207	148	2.314				
70	89.4	1.277	169	2.414				
80	116	1.390	206	2.469	119	1.429	223	2.678
90	148	1.500	269	2.727				
100	189	1.634	325	2.810	195	1.689	321	2.778
110	230	1.710	365	2.714				
120	277	1.786	407	2.625				
140	372	1.845	477	2.366				
150					445	1.954	585	2.572
160	473	1.852	538	2.107				
180	572	1.809	615	1.945				
200	683	1.778	702	1.827	723	1.881	761	1.979
230	836	1.673	808	1.617				
250					1009	1.722	945	1.613
260	1018	1.612	989	1.566				
300	1246	1.498	1098	1.320	1320	1.587	1235	1.484
340	1485	1.399	1289	1.215				
380	1735	1.315	1528	1.158				
400					2008	1.377	1835	1.258
420	2006	1.249	1830	1.140				
460	2290	1.192	2028	1.056				
500	2562	1.131	2085	0.921	2841	1.254	2425	1.071

well with those calculated from MCS mobility data using the GER relations. The experimental data of Thackston *et al.* [14] and the GER calculations using their mobility data were found to be higher than the MCS-derived data at moderate and high values of E/N. Relatively close agreement exists for all data below the peak region of about E/N = 100 Td.

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Further evidence of the accuracy of the MCS calculations may be obtained from a plot of  $D_T/K$  and  $D_L/K$ with  $(E/N)^2$ . Theoretically as E/N tends to zero,  $D_T/K$ and  $D_L/K$  should merge and vary linearly. This behavior is shown in Fig. 3, and both  $D_T/K$  and  $D_L/K$ approach the theoretical limit of 26.1 mV for T=303 K as  $E/N \rightarrow 0$  as predicted by the Einstein equation,

$$D_{T,L}/K = k_B T/e . (4)$$

Our experimental values of  $D_T/K$  are included in the plots for comparison.

It can be seen from Figs. 1 and 2 that the transversediffusion peaks at around E/N = 150 Td, in contrast to the peak at about 90 Td for the longitudinal diffusion. This apparent shift in the peak can be explained in terms of an additional collision energy attributed to the forward ion drift velocity in the longitudinal-diffusion coefficient. This shift is similarly observed for the Li<sup>+</sup>-He system [20]. The MCS calculations are expected to be more jagged since they are strictly based on random numbers and no correlation exists between calculations of (even) neighborhood E/N points. However, the jaggedness is always maintained within the errors of the calculations, which are only a function of the statistical size of the samples used as input. On the other hand, the experimental data are inherently smoother, as expected, since errors in all experimental input parameters permeate



FIG. 3. Plots against  $(E/N)^2$  of the MCS-calculated  $D_T/K$  (diamonds) and  $D_L/K$  (triangles) using the KMV potential and present experimental  $D_T/K$  (squares connected by straight lines), showing their expected zero-field limit of 26.1 mV at 303 K.



FIG. 4. Plots against the ion-to-neutral-atom mass ratio (m/M), of E/N at which the maximum  $D_T^{(r)}$  occurs, for various Na<sup>+</sup>-noble-gas systems.

smoothly across different E/N settings.

The Na<sup>+</sup>-Ne and Na<sup>+</sup>-Ar [15], and the Na<sup>+</sup>-Kr systems appear to indicate a relationship between m/M(ion-to-neutral-atom mass ratio) and the value of E/Ncorresponding to the maximum value of  $D_T^{(r)}$  for the alkali-metal-ion-noble-gas system. For Na<sup>+</sup>-Ne, Na<sup>+</sup>-Ar, and Na<sup>+</sup>-Kr, the values of m/M are 1.15, 0.575, and 0.274, respectively. The maximum values of  $D_T^{(r)}$  occur at E/N=40, 90, and 150 Td, respectively. From a recent preliminary experimental study of Na<sup>+</sup> ions in Xe gas [23], a plot of  $D_T^{(r)}$  gives a maximum value at 180 Td. The m/M value of Na<sup>+</sup>-Xe is 0.175. A relatively smooth trend of m/M with E/N corresponding to the maximum value of  $D_T^{(r)}$  is shown in Fig. 4. Attempts to obtain a more general correlation between the peaks of transport coefficients and the values of E/N at which the peaks occur for different m/M ratios are in progress.

#### CONCLUSION

We have presented experimental data of  $D_T/K$  for the heretofore unstudied system of Na<sup>+</sup> ions in Kr with an overall accuracy of 3% in the range E/N=5-500 Td. They provide values which may be used as input data for the determination of ion-neutral-atom interaction potential particularly in the medium- and long-range regions, and with which theoretical calculations can be compared. Our MCS calculated values of  $D_T/K$  and  $D_L/K$  have reaffirmed the validity of the KMV potential curve for the Na<sup>+</sup>-Kr system. Experimental values of  $D_T/K$  were found to be in excellent agreement with the calculations of the MCS method based on the KMV potential. The present experimental data would hopefully contribute to the scarce pool of transverse diffusion coefficient data from which a general functional form of its peak value at a certain value of E/N for different ion-neutral-atom mass ratios can hopefully be established.

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