Limitations on concentration measurements and gradient discerning times in cellular systems

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This work reports on two results. At first we revisit the Berg and Purcell calculation that provides a lower bound to the error in concentration measurement by cells by considering the realistic case when the cell starts measuring the moment it comes in contact with the chemoattractants, instead of measuring after equilibrating with the chemotactic concentration as done in the classic Berg and Purcell paper. We find that the error in concentration measurement is still the same as evaluated by Berg and Purcell. We next derive a lower bound on measurement time below which it is not possible for the cell to discern extracellular chemotactic gradients through spatial sensing mechanisms. This bound is independent of diffusion rate and concentration of the chemoattracts and is instead set by detachment rate of ligands from the cell receptors. The result could help explain experimental observations.

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I. INTRODUCTION

Cells have to make sense out of their surroundings. They have to go toward sources of nutrition and away from sources of danger. Berg and Purcell [1], found that concentration measurement by cells such as Escherichia coli approach one of optimum design. The fluctuations in ligand receptor binding were later incorporated in Reference [2] who found that cells perform concentration measurements within limits set by these fluctuations. Since then, there have been many theoretical works in understanding the limitations imposed on concentration measurements by the cell. Some of these include [3] that considered effect of ligand diffusion on fluctuations in occupancy of receptors. Understanding how increasing number of receptors could affect reduction in measurement noise due to receptor noise were considered in Ref. [4]. Corrections to contributions to diffusive arrival of ligands obtained in Ref. [2] were considered in Ref. [5]. Constraints placed by energy consumption in concentration measurement of cells was considered in Refs. [6-8]. Limits to concentration sensing by a cell in an environment of interfering ligands was considered in Ref. [9]. Maximum likelihood estimation of concentration of ligands by looking at history of attachment detachment of ligands to the receptor was considered in Ref. [10]. Most calculations that involve evaluation of limits to concentration detection assume that the cell receptors have reached equilibrium with the surrounding ligands after which the attachment or detachment dynamics is considered. However, realistic calculations to error measurements should consider measurement of concentration beginning the moment exposure to chemoattractants occurs. We label this scenario as the nonequilibrium case in

In addition to concentration measurement, cells also have to measure extracellular gradients of chemoattractants. Cells have evolved to measure gradients with great accuracy. Cells can detect gradients of 1 - 2% difference across the cell [11-13]. An optimal response has been seen in the cells where the difference in receptor occupancy between front and back of cells is only 10 occupied receptors [14]. Understanding the limits to chemotactic gradient measurement by the cells were studied in Ref. [15] which idealized the cell as a perfectly absorbing sphere and a perfectly monitoring sphere. Fluctuation dissipation theorem was used to consider limitations imposed on gradient sensing due to ligand receptor kinetics in Ref. [16]. Modelling the surface of the cell as an Ising spin chain showed improved ability to detect gradients if receptor cooperativity was introduced [17]. A signal transduction modeling to understand gradient readouts was considered in Ref. [18]. If one considers the works [15,16], it was shown that the limits to gradient measurements by the cell went like $\frac{1}{\sqrt{DacT}}$, up to multiplicative constants. This implies that one could decrease the time of measurement T to an arbitrary degree by increasing the concentration of chemoattractants

comparison to the equilibrium case which was considered in the Berg-Purcell calculation where the receptor is assumed to have been in equilibrium with the ligand concentration throughout the measurement history. In the first part of the paper we show that the bound on error in concentration measurement in the nonequilibrium scenario is still given by what is expected from Berg-Purcell like studies, in the limit of large measurement times. Berg and Purcell considered the quantity getting measured to be the fraction of time the receptor is occupied by the ligand. On considering a generic quantity defining a measurement as the linear combination of powers of the fraction of time a receptor is occupied, we show that the error in concentration measurement is similar in the equilibrium and nonequilibrium cases for large measurement times.

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and still be able to discern the gradient. In the present work we represent the cell as sphere covered with receptors with a point source of chemoattractants a particular distance away. We show that there is a bound on the measurement time, below which it is not possible to discern the concentration gradient set up by the source. This time being independent of the concentration and diffusion rates of the chemoattractant, implies one cannot simply increase these parameters to decrease the measurement times to being as small as possible. For *Dictyostelium* we evaluate this time to be around ~4 s. The time duration of the pulses not being sufficient to discern gradients, could be a reason as to why *Dictyostelium* cells subjected to 5-s pulses [20] do not respond to the chemotactic gradients of varied concentrations, while being placed at varied distances from the chemoattractant source.

In the next section we present the calculation which evaluates the error in concentration measurement by cells, when the concentration detection starts as soon as the cells are bought in contact with chemoattractants. The calculation of the minimum time below which it is not possible to discern concentration gradients is presented in the following section. We end with conclusions.

II. CONCENTRATION MEASUREMENTS COMMENCING AFTER IMMEDIATE EXPOSURE

Cells would in general start measuring chemotactic concentration the moment they come in touch with the chemoattractants. Berg and Purcell [1] assumed the receptor has already equilibrated with the surrounding chemotactic concentration before concentration measurement are considered. Let us instead consider the more realistic case when the cell has started measuring the moment it got in touch with the chemotactic concentration. To understand this, assume the cell represented by a sphere is dropped into a chemoattractant concentration c as shown in Fig. 1. The occupation probability of each cell receptor obeys the relation

$$\frac{dp(t)}{dt} = k_{+}c[1-p(t)] - k_{-}p(t).$$
(1)

With the initial condition p(0) = 0, we get the solution

$$p(t) = \frac{k_+ c}{k_+ c + k_-} [1 - e^{-(k_+ c + k_-)t}].$$
 (2)

Let us say that the receptor occupancy at time t is given by O(t), which equals to zero if the receptor is not occupied and 1 if its occupied. The fraction of time a receptor is occupied in a time duration T is

$$m(T) = \frac{1}{T} \int_0^T O(t) dt.$$
(3)

This is something that the cell can measure. If we split the time T in to $\frac{T}{dt}$ intervals, then the probability of having n intervals in which receptor was occupied is simply

$$e^{-\int_0^T p(t)dt} \frac{\left[\int_0^T p(t)dt\right]^n}{n!}.$$
(4)



FIG. 1. In (a) we have a cell in a constant concentration of chemoattractants. The limitations to concentration measurement were studied in the classic work of Berg-Purcell, where the cell was idealized as a sphere in a constant concentration of chemoattractants. In (b) the cell moves toward a higher concentration of chemoattractants. We simulate this scenario with the sphere dropped into a bottle of chemoattractants. In this paper we show that the error in concentration measurement in (b) is similar to the error in (a) in limit of large measurement times. (Images created with BioRender.com.)

Hence an average estimate of m(T) would be

$$\langle m(T) \rangle = \frac{1}{T} \sum_{n=0,\infty} ne^{-\int_0^T p(t)dt} \frac{\left[\int_0^T p(t)dt\right]^n}{n!} = \frac{1}{T} \int_0^T p(t)dt$$
$$= \frac{k_+c}{k_+c+k_-} - \frac{(1-e^{-(k_+c+k_-)T})}{(k_+c+k_-)T} \frac{k_+c}{k_+c+k_-},$$
(5)

which in limit of large T becomes

$$\langle m(T) \rangle = \frac{k_+ c}{k_+ c + k_-} \left[1 - \frac{1}{(k_+ c + k_-)T} \right]$$
$$\sim \frac{k_+ c}{k_+ c + k_-}, \quad T \to \infty.$$
(6)

The average occupancy could be inverted by the cell to evaluate the chemoattractant concentration *c*. To evaluate the noise in this estimate of *c* we need to evaluate $\langle m(T)^2 \rangle$ for which we need to evaluate $G(t, t') = \langle O(t)O(t') \rangle$. Assume the receptor is occupied at time *t*. The receptor would be still found occupied at time *t'*, if the receptor remains attached for some time τ after time *t*, then has *n* pairs of time intervals t_d^i, t_a^i with $i \in [0, n]$, where the receptor is in a detached or attached state, such that $\tau + \sum_{i=1,n} t_d^i + t_a^i = t - t'$. For any particular timing combination, the probability of realizing the same is

$$e^{-\int_0^\tau k_- dx} [k_+ ck_-]^n \Pi_{i=1,n} dt_d^i dt_a^i e^{-k_+ c \int_0^{t_d^i} dx} e^{-k_- \int_0^{t_a^i} dx}.$$
 (7)

One then sums over all possible timing combinations and all values of *n* to get the probability of still finding the receptor occupied at time t'. Let us call this probability P(t, t'). Since the only information getting into the evaluation of P(t, t') is that the ligand is attached at time *t* and since the ligand attachment detachment is Markovian, the past history before attachment at time *t* is irrelevant, so whether the receptor has equilibrated with the ligands or not is irrelevant in evaluation of P(t, t'). It is obvious that G(t, t') = p(t)P(t, t') t < t'. Since, we already know from Ref. [1] that in the case where receptor has equilibrated with the chemoattractant concentration

$$G_{\rm eq}(t,t') = \bar{p}^2 + \bar{p}(1-\bar{p})e^{-\frac{k-|t-t'|}{1-\bar{p}}}$$
(8)

with

$$\bar{p} = \frac{k_+ c}{k_+ c + k_-}.$$
 (9)

We get

$$P(t,t') = \bar{p} + (1-\bar{p})e^{-\frac{k-|t-t'|}{1-\bar{p}}}.$$
(10)

So,

$$G(t, t') = p(t) \left[\frac{k_{+}c}{k_{+}c + k_{-}} + \frac{k_{-}}{k_{+}c + k_{-}} e^{-(k_{+}c + k_{-})|t'-t|} \right],$$

× t < t'. (11)

Hence,

$$\langle m(T)^2 \rangle = \frac{1}{T^2} \int_0^T \int_0^T G(t, t') dt dt'$$

= $\frac{1}{T^2} \int_0^T dt \int_0^t dt' p(t') P(t', t)$
+ $\frac{1}{T^2} \int_0^T dt' \int_0^{t'} dt p(t) P(t, t')$

$$= \frac{2}{T^2} \int_0^T dt' \int_0^{t'} dt \, p(t) P(t, t').$$
(12)

The second last equation arises from the one above, because P(t, t') = P(t', t) and on interchanging $t' \leftrightarrow t$, we find the two integrals in the sum are equal. Hence as shown in the Appendix, in the limit where T is very large $(T^{-1} < < k_+c, k_-)$,

$$\langle m(T)^2 \rangle - \langle m(T) \rangle^2 = \left(\frac{k_+c}{k_+c+k_-}\right)^2 \frac{2}{(k_+c+k_-)T} \frac{k_-}{k_+c}$$
(13)

and

$$\frac{\delta c}{c} = \sqrt{\frac{2(k_{+}c + k_{-})}{T}} \left[\frac{k_{-}}{k_{+}c}\right] \frac{1}{k_{-}}.$$
 (14)

This is the same result as obtained by Berg and Purcell [1]. What we hence see is that the error in the nonequilibrium case is the same as the equilibrium case when $(k_+c + k_-)T \gg 1$. Berg and Purcell assumed that the measurement made by the cell was the fraction of time the receptor was occupied. Receptor activation starts a series of downstream reactions in the cell that leads to a cellular response. There is no reason to expect that the cellular response (however it is quantified) is simply proportional to the fraction of time the receptor is occupied. It is more likely that as a general case the measurement could be proportional to a linear combinations of powers of the fraction of time the receptor as $Q(T) = \sum_{i>0} a_i m(T)^i$, where a_i are constants dependent on the nature of the system being analyzed and the corresponding readout. Now,

$$\langle m(T) \rangle^{n} = \left[\frac{1}{T} \int_{0}^{T} p(t) dt \right]^{n} = \left[\frac{k_{+}c}{k_{+}c + k_{-}} \right]^{n} \left[1 - \frac{1}{T} \int_{0}^{T} e^{-(k_{+}c + k_{-})t} dt \right]^{n}$$

$$= \left[\frac{k_{+}c}{k_{+}c + k_{-}} \right]^{n} \left\{ 1 - \frac{n}{T} \int_{0}^{T} e^{-(k_{+}c + k_{-})t} dt + \frac{n(n-1)}{2!T^{2}} \left[\int_{0}^{T} e^{-(k_{+}c + k_{-})t} dt \right]^{2} \right\}$$

$$= \langle m(T) \rangle_{eq}^{n} - \left[\frac{k_{+}c}{k_{+}c + k_{-}} \right]^{n} \frac{n}{(k_{+}c + k_{-})T} + \mathcal{O}\left(\frac{1}{T^{2}} \right), \quad \text{if } (k_{+}c + k_{-})T \gg 1.$$

$$(15)$$

Also,

$$\langle m(T)^{n} \rangle = \frac{1}{T^{n}} \int_{0}^{T} dt_{1} \dots \int_{0}^{T} dt_{n} \langle O(t_{1}) \dots O(t_{n}) \rangle = \frac{n!}{T^{n}} \int_{0}^{T} dt_{n-1} \dots \int_{0}^{t_{1}} dt p(t) P(t, t_{1}) P(t_{1}, t_{2}) \dots P(t_{n-2}, t_{n-1})$$

$$= \langle m(T)^{n} \rangle_{eq} - \frac{n!}{T^{n}} \int_{0}^{T} dt_{n-1} \dots \int_{0}^{t_{1}} dt \frac{k_{+}c}{k_{+}c + k_{-}} e^{-(k_{+}c+k_{-})t} \left[\frac{k_{+}c}{k_{+}c + k_{-}} + \frac{k_{-}}{k_{+}c + k_{-}} e^{-(k_{+}c+k_{-})|t_{1}-t|} \right] \dots$$

$$\dots \left[\frac{k_{+}c}{k_{+}c + k_{-}} + \frac{k_{-}}{k_{+}c + k_{-}} e^{-(k_{+}c+k_{-})|t_{n-1}-t_{n-2}|} \right] = \langle m(T)^{n} \rangle_{eq} - \left[\frac{k_{+}c}{k_{+}c + k_{-}} \right]^{n} \frac{n!}{T^{n}} \int_{0}^{T} dt_{n-1} \dots \int_{0}^{t_{1}} dt e^{-(k_{+}c+k_{-})t} \left[e^{-(k_{+}c+k_{-})|t_{1}-t|} + \dots + e^{-(k_{+}c+k_{-})|t_{n-1}-t_{n-2}|} \right]$$

$$+ terms higher order in \frac{1}{T}$$

$$\sim \langle m(T)^{n} \rangle_{eq} - \left[\frac{k_{+}c}{k_{+}c + k_{-}} \right]^{n} \frac{n}{(k_{+}c + k_{-})T} + \mathcal{O}\left(\frac{1}{T^{2}}\right); \quad \text{if } (k_{+}c + k_{-}) \gg 1.$$

$$(16)$$

The factor of n! in the second equation has a origin similar to the factor of 2 in Eq. (12). Now,

$$\langle m(T)^{n} \rangle_{eq} - \langle m(T) \rangle_{eq}^{n} = \frac{n!}{T^{n}} \int_{0}^{T} dt_{n-1} \dots \int_{0}^{t_{1}} dt \frac{k_{+}c}{k_{+}c_{+}k_{-}} \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} + \frac{k_{-}}{k_{+}c_{+}k_{-}} e^{-(k_{+}c_{+}k_{-})|t_{1}-t_{1}|} \right] \dots$$

$$\dots \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} + \frac{k_{-}}{k_{+}c_{+}k_{-}} e^{-(k_{+}c_{+}k_{-})|t_{n-1}-t_{n-2}|} \right] - \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{n}$$

$$= \frac{k_{-}}{k_{+}c_{+}k_{-}} \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{n-1} \frac{n!}{T^{n}} \int_{0}^{T} dt_{n-1} \dots \int_{0}^{t_{1}} dt \left[e^{-(k_{+}c_{+}k_{-})|t_{1}-t_{1}|} + \dots + e^{-(k_{+}c_{+}k_{-})|t_{n-1}-t_{n-2}|} \right]$$

$$+ \text{ terms higher order in } \frac{1}{T}$$

$$= \frac{k_{-}}{k_{+}c_{+}k_{-}} \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{n-1} \frac{n(n-1)}{(k_{+}c_{+}k_{-})T} + \mathcal{O}(\frac{1}{T^{2}}), \quad \text{if } (k_{+}c_{+}k_{-})T \gg 1.$$

$$(17)$$

Hence,

$$\begin{aligned} \langle Q(T)^{2} \rangle &- \langle Q(T) \rangle^{2} \\ &= \sum_{i,j>0} a_{i} a_{j} [\langle m(T)^{i+j} \rangle_{eq} - \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{i+j} \frac{i+j}{(k_{+}c_{+}k_{-})T} + \mathcal{O}\left(\frac{1}{T^{2}}\right) \\ &- \left\{ \langle m(T)^{i} \rangle_{eq} - \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{i} \frac{i}{(k_{+}c_{+}k_{-})T} + \mathcal{O}\left(\frac{1}{T^{2}}\right) \right\} \left[\langle m(T)^{j} \rangle_{eq} - \left[\frac{k_{+}c}{k_{+}c_{+}k_{-}} \right]^{j} \frac{j}{(k_{+}c_{+}k_{-})T} + \mathcal{O}\left(\frac{1}{T^{2}}\right) \right] \right] \\ &, if (k_{+}c_{+}k_{-})T \gg 1 \\ &= \sum_{i,j>0} a_{i} a_{j} [\langle m(T)^{i+j} \rangle_{eq} - \langle m(T)^{i} \rangle_{eq} \langle m(T)^{j} \rangle_{eq}] + \mathcal{O}\left(\frac{1}{T^{2}}\right), \quad if (k_{+}c_{+}k_{-})T \gg 1 \\ &= \langle Q(T)^{2} \rangle_{eq} - \langle Q(T) \rangle_{eq}^{2} + \mathcal{O}\left(\frac{1}{T^{2}}\right), \quad if (k_{+}c_{+}k_{-})T \gg 1 \\ &= \langle Q(T)^{2} \rangle_{eq} - \langle Q(T) \rangle_{eq}^{2}, \quad if (k_{+}c_{+}k_{-})T \gg 1 \end{aligned}$$
(18)

This is because $\sum_{i,j>0} a_i a_j [\langle m(T)^{i+j} \rangle_{eq} - \langle m(T)^i \rangle_{eq} \langle m(T)^j \rangle_{eq}]$ is $\mathcal{O}(\frac{1}{T})$ if $(k_+c+k_-)T \gg 1$, as it contains terms of the form $\langle m(T)^n \rangle_{eq} - \langle m(T) \rangle_{eq}^n$ which are $\mathcal{O}(\frac{1}{T})$ in the same limit. For a specially chosen combination of $\{a_i\}'$ s, we could have that $\sum_{i,j>0} a_i a_j [\langle m(T)^{i+j} \rangle_{eq} - \langle m(T)^i \rangle_{eq} \langle m(T)^j \rangle_{eq}]$ may be of $\mathcal{O}(\frac{1}{T^2})$, but for a generic combination of $\{a_i\}'$ s, this is not the case and hence $\langle Q(T)^2 \rangle - \langle Q(T) \rangle^2 = \langle Q(T)^2 \rangle_{eq} - \langle Q(T) \rangle_{eq}^2$, if $(k_+c+k_-)T \gg 1$. Also,

$$\begin{aligned} \langle Q(T) \rangle &= \sum_{i>0} a_i \langle m(T)^i \rangle \\ &= \sum_{i>0} a_i \left[\frac{k_+ c}{k_+ c + k_-} \right]^i + \mathcal{O}\left(\frac{1}{T}\right), \text{ if } (k_+ c + k_-)T \gg 1 \\ &= \langle Q(T) \rangle_{\text{eq}} + \mathcal{O}\left(\frac{1}{T}\right), \text{ if } (k_+ c + k_-)T \gg 1 \\ &= \langle Q(T) \rangle_{\text{eq}}, \text{ if } (k_+ c + k_-)T \gg 1. \end{aligned}$$
(19)

We hence see that for generic $\{a_i\}'$ s, the evaluation of $\frac{\delta c}{c}$ is the same in the limit $(k_+c + k_-)T \gg 1$, whether we

consider the Berg-Purcell like equilibrium calculations which assume that the receptor has already equilibrated with the ligands before measurement starts or we consider the realistic nonequilibrium case where the measurement commences the moment the receptor is bought in contact with the ligands.

III. LOWER BOUND ON GRADIENT MEASUREMENT TIMES IN CELLS

To consider the problem of time required to measure a chemotactic gradient by the cell, consider the following setup illustrated in Fig. 2. A sphere representing the cell is covered with receptors, such that the probability of the receptor at position (R, θ) being occupied is $p(\theta, t)$ which at equilibrium becomes $p(\theta)$. The source of chemoattractants is at a distance *a* from the center of the sphere. Let us say the concentration of chemoattractants at a position (r, θ) , is $c(r, \theta, t)$ which at equilibrium becomes $c(r, \theta)$. We hence have that

$$\frac{dp(\theta,t)}{dt} = k_+ c(R,\theta,t) [1 - p(\theta,t)] - k_- p(\theta,t) \quad (20)$$



FIG. 2. (i) A cell getting exposed to a point source of chemoattractants. Experiments have shown [20] that for stimulation with 5-s pulses for a range of cAMP concentrations, no evidence of chemotaxis was observed at any distance from the cAMP pipette. Calculations in this paper suggest a measurement time bound of 4 s, below which it is not possible for cells to discern gradients irrespective of concentration of chemoattractants involved. (Images created with BioRender.com.) (ii) The various quantities entering into the calculation in the text. The cell is represented by a sphere of radius R, and the source of chemoattractant is represented by in green color at a distance of a from the center of the sphere.

and

$$-D\frac{\partial c(r,\theta,t)}{\partial r}|_{r=R} = k_{+}c(R,\theta,t)[1-p(\theta,t)] - k_{-}p(\theta,t).$$
(21)

At equilibrium we have

$$k_{-}p(R,\theta) = k_{+}c(R,\theta)[1-p(R,\theta)]$$
$$D\frac{\partial c(r,\theta)}{\partial r}|_{r=R} = 0.$$
 (22)

The symmetry of the problem implies no ϕ dependence. Since the source of chemoattractants is located at r = a, we have the equation

$$D\nabla^2 c(r,\theta,\phi) = C\delta(r-a)\delta(\theta), \qquad (23)$$

whose solution as described in the Appendix is

$$c(R,\theta) = \frac{C}{4\pi Da} + \frac{C}{4\pi Da} \sum_{n=1,\infty} \frac{R^n}{a^n} P_n(\cos\theta). \quad (24)$$

Since

$$p(R, \theta) = \frac{k_+ c(R, \theta)}{k_+ c(R, \theta) + k_-},$$
 (25)

we have the error in concentration measurement [1] for large times

$$\frac{\Delta c_{\rm rms}(R,\theta)}{c(R,\theta)} = \sqrt{\frac{2}{k_- T p(R,\theta)}} = \sqrt{\frac{2[k_+ c(R,\theta) + k_-)]}{k_- T k_+ c}}.$$
(26)

However, including terms to all orders in $\frac{1}{T}$, we have

$$\frac{\Delta c_{\rm rms}(R,\theta)}{c(R,\theta)} = \sqrt{\frac{2([k_+c(R,\theta)T + k_-T] - \{1 - e^{-[k_+c(R,\theta) + k_-]T}\})}{(k_-T)[k_+c(R,\theta)T]}}.$$
(27)

Hence, in order to be able to discern a gradient we should have that

$$c(R, \theta = 0) - c(R, \theta = \pi) > \Delta c_{\rm rms}(R, \theta = 0) + \Delta c_{\rm rms}(R, \theta = \pi)$$

$$= \sqrt{\frac{2c(R, \theta = 0)(\{k_+ c(R, \theta = 0)T + (k_-T - 1) + e^{-[k_+ c(R, \theta = 0) + k_-]T}\})}{k_- k_+ T^2}}$$

$$+ \sqrt{\frac{2c(R, \theta = 0)(\{k_+ c(R, \theta = \pi)T + (k_-T - 1) + e^{-[k_+ c(R, \theta = \pi) + k_-]T}\})}{k_- k_+ T^2}}.$$
(28)

Assume that the measurement time T is such that $k_T > 1$, then

$$c(R, \theta = 0) - c(R, \theta = \pi) > \sqrt{\frac{2c(R, \theta = 0)[k_{+}c(R, \theta = 0)]}{k_{-}k_{+}T}} + \sqrt{\frac{2c(R, \theta = \pi)[k_{+}c(R, \theta = \pi)]}{k_{-}k_{+}T}}$$

$$\frac{C}{4\pi Da} \frac{2R/a}{1 - R^{2}/a^{2}} > \sqrt{\frac{2}{k_{-}T}} [c(R, \theta = 0) + c(R, \theta = \pi)]$$

$$\frac{C}{2\pi Da} \frac{R/a}{1 - R^{2}/a^{2}} > \sqrt{\frac{2}{k_{-}T}} \left[\frac{C}{2\pi Da} + \frac{C}{2\pi Da} \frac{R^{2}/a^{2}}{1 - R^{2}/a^{2}}\right]$$

$$T > \frac{2}{k_{-}} \left[\frac{1 + \frac{R^{2}/a^{2}}{1 - R^{2}/a^{2}}}{\frac{R/a}{1 - R^{2}/a^{2}}}\right]^{2} = \frac{2}{k_{-}} \frac{a^{2}}{R^{2}}.$$
(29)

We hence get a bound on the measurement time, independent of the nature of the chemoattractant or its concentration, below which it is not possible for the cell to decipher the concentration gradient. The smallest value this bound can take is when a = R or $T_{\text{smallest}} = \frac{2}{k}$. Since $k_{-}T \ge 2$, the assumption $k_T \ge 1$ used to derive above implies consistency of the derivation above. Experiments [19] have evaluated the detachment rates of Cy3-cAMP molecules from receptors on the surface of dictyostelium and found that k_{-} was in the range of $1.1-0.39 \text{ s}^{-1}$ in the anterior pseudopod region of the cell and $0.39-0.1 \text{ s}^{-1}$ in the posterior tail region. So assuming an rate of $k_{-} \sim 0.5 \text{ s}^{-1}$, gives a bound on $T_{\text{smallest}} = \frac{2}{k} = 4 \text{ s}$. This is in line with Ref. [20] where it was seen that the simulation of Dictyostelium with 5-s pulses using a wide range of cAMP concentrations, shows no evidence of chemotaxis by cells at any distance from the pipette. A possible reason for this could be the fact that the time period for which the signal was active, was not sufficient for discerning the gradient by the cell, leading to no response.

IV. CONCLUSION

Any calculation that evaluates the error in measurement of concentation should consider the cell from the time when it first gets in touch with the ligand concentration. In the first part of the paper we evaluated the error in such a scenario. We showed that the error is the same, in the limit of large measurement time, as evaluated in the Berg-Purcell calculation which assumed that the cell was already in equilibrium with the ligand concentration when measurement commenced. In the nonequilibrium case and for large measurement times $T \gg (k_+c + k_-)^{-1}$, the receptor will have equilibrated with the ligand concentration for most of its measurement history and hence it is expected that $\langle m(T) \rangle$ as well $\langle m(T)^2 \rangle$ will yield result similar to the case in which the receptor was already in equilibrium with the surrounding concentration when the measurement commenced. However, there is no reason apriori that the difference between the two $\langle m(T)^2 \rangle$ – $\langle m(T) \rangle^2$ will yield similar values. We however see an equality between the equilibrium and nonequilibrium estimates. The cellular response would generally be proportional to some linear combination of powers of the fraction of time the receptor is occupied. By considering the measured quantity to go as $Q(T) = \sum_{i>0} a_i m(T)^i$, where the $\{a_i\}'$ s are generically chosen, we found that the error in concentration measurement in the equilibrium as well as nonequilibrium cases is the same for large measurement times.

Reference [2] utilized a fluctuation dissipation framework to talk about how noisy ligand attachment or detachment events add limitations to concentration measurements. Their work first considers the case without diffusion, producing an rms error in the estimate of receptor occupancy n given by

$$\delta n_{\rm rms} = \sqrt{\frac{2\bar{n}(1-\bar{n})}{(k_+\bar{c}+k_-)\tau}}$$
 (30)

including diffusion increases this error to

$$\delta n_{\rm rms} = \sqrt{\frac{2\bar{n}(1-\bar{n})}{(k_+\bar{c}+k_-)T}} + \frac{(\bar{n}(1-\bar{n}))^2}{\pi D a \bar{c} T},\tag{31}$$

where \bar{c} , \bar{n} are averages of measured concentration and receptor occupancy, T is the measurement time, D is the diffusion coefficient, and a is the receptor size. Hence, including the effect of diffusion only serves to increase the error. Since Ref. [2] considered the situation where the ligands have equilibrated with the receptor, they could use the fluctuation dissipation frameworks to do their calculations. Our work in part one of the paper is considering the case when the receptor is suddenly got in touch with the ligands implying a nonequilibrium framework, so fluctuation dissipation frameworks to do calculations are not possible here. Hence, including diffusion effects if possible would be more involved and we do not attempt the same in our paper. However from the lessons of equilibrium calculations as done in Ref. [2], we can say that diffusion effects will only add to measurement errors. How this addition will look like or whether it will yield the same form as equation above, given the similarity in the equilibrium and nonequilibrium case when diffusion is not considered, is an open question.

It is also quite interesting that the observation that the bound on measurement times below which gradients cannot be detected, being independent of the chemoattractant properties such as concentration and diffusion was not noticed in theoretical works up until now. For example, in Ref. [16] the authors consider the noise in receptor ligand binding or unbinding to evaluate the limits to detection of gradients. The only information of the chemotactic gradient that appears in their analysis is the ligand concentration at points on the surface of the cell. For example, the section "Two Receptors" evaluates the variance in the difference between concentration between two points on the cell surface and relates the same to the concentration at these two points. The work does not calculate any relation between the ligand concentration of various points on the cell surface but assumes the same to be given. In our work we consider a point source that produces the chemotactic gradient around the cell itself and calculate the concentration on various points on the cell surface. Using this crucial information we could evaluate a bound on the measurement time itself below which no gradient could be discerned irrespective of concentrations involved, an observation that was missed or not obtained in works such as Ref. [16].

The FRET measurements for example in Ref. [21] point to various signaling elements in the cell getting activated; however, they do not imply or suggest that any chemotaxis in direction of the source is attained for a pulse of magnitude below 5 s, since as is known from Ref. [20] there is no chemotactic response in direction of the gradient for pulses of time duration 5 s. The part two of the paper provides a reason as to why for all possible ligand concentrations there is no chemotactic response observed for the 5-s pulse, tracing this reason to the limitations imposed by the ligand binding unbinding noise.

APPENDIX

1. I

We elaborate on some calculations from the section on concentration measurements. We know that

$$\langle m(T)^2 \rangle = \frac{1}{T^2} \int_0^T \int_0^T G(t, t') dt dt'$$

$$= \frac{1}{T^2} \int_0^T dt \int_0^t dt' p(t') \left[\frac{k_+ c}{k_+ c + k_-} + \frac{k_-}{k_+ c + k_-} e^{-(k_+ c + k_-)(t-t')} \right]'$$

$$+ \frac{1}{T^2} \int_0^T dt' \int_0^{t'} dt p(t) \left[\frac{k_+ c}{k_+ c + k_-} + \frac{k_-}{k_+ c + k_-} e^{-(k_+ c + k_-)(t'-t)} \right]$$

$$= \frac{2}{T^2} \int_0^T dt' \int_0^{t'} dt p(t) \left[\frac{k_+ c}{k_+ c + k_-} + \frac{k_-}{k_+ c + k_-} e^{-(k_+ c + k_-)(t'-t)} \right]$$

$$= \frac{2}{T^2} \int_0^T dt' \int_0^{t'} dt \left[\frac{k_+ c}{k_+ c + k_-} \right] (1 - e^{-(k_+ c + k_-)t}) \left[\frac{k_+ c}{k_+ c + k_-} + \frac{k_-}{k_+ c + k_-} e^{-(k_+ c + k_-)(t'-t)} \right]$$

$$= \frac{1}{T^2} \left[-\frac{2ck_- k_+}{(ck_+ + k_-)^4} + \frac{2ck_- k_+ e^{-T(ck_+ + k_-)}}{(ck_+ + k_-)^4} + \frac{c^2k_+^2}{(ck_+ + k_-)^4} - \frac{c^2k_+^2 e^{-T(ck_+ + k_-)}}{(ck_+ + k_-)^4} \right]$$

$$+ \frac{1}{T} \left[-\frac{c^3k_+^3}{(ck_+ + k_-)^4} + \frac{c^2k_-k_+^2 e^{-T(ck_+ + k_-)}}{(ck_+ + k_-)^4} + \frac{c^2k_-k_+^2 e^{-T(ck_+ + k_-)}}{(ck_+ + k_-)^4} \right]$$

$$+ \left[\frac{c^4k_+^4}{2(ck_+ + k_-)^4} + \frac{c^3k_-k_+^3}{(ck_+ + k_-)^4} + \frac{c^2k_-^2k_+^2}{2(ck_+ + k_-)^4} \right],$$
(A1)

where we plugged into the second last equation [Eq. (2)]. Also we have

$$\langle m(T) \rangle^{2} = \left[\int_{0}^{T} \frac{k_{+}c}{k_{+}c+k_{-}} \right] \left[1 - e^{-(k_{+}c+k_{-})t} \right]^{2}$$

$$= \frac{1}{T^{2}} \left[\frac{c^{2}k_{+}^{2}}{(ck_{+}+k_{-})^{4}} + \frac{c^{2}k_{+}^{2}e^{-2T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} - \frac{2c^{2}k_{+}^{2}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} \right]$$

$$+ \frac{1}{T} \left[\frac{2c^{3}k_{+}^{3}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} - \frac{2c^{3}k_{+}^{3}}{(ck_{+}+k_{-})^{4}} + \frac{2c^{2}k_{-}k_{+}^{2}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} - \frac{2c^{2}k_{-}k_{+}^{2}}{(ck_{+}+k_{-})^{4}} \right]$$

$$+ \left[\frac{c^{4}k_{+}^{4}}{(ck_{+}+k_{-})^{4}} + \frac{2c^{3}k_{-}k_{+}^{3}}{(ck_{+}+k_{-})^{4}} \right] \right]$$

$$(A2)$$

Hence

$$\langle m(T)^{2} \rangle - \langle m(T) \rangle^{2} = \frac{1}{T^{2}} \left[-\frac{c^{2}k_{+}^{2}e^{-2T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} + \frac{c^{2}k_{+}^{2}}{(ck_{+}+k_{-})^{4}} + \frac{4ck_{-}k_{+}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} - \frac{4ck_{-}k_{+}}{(ck_{+}+k_{-})^{4}} \right] + \frac{1}{T} \left[-\frac{2c^{3}k_{+}^{3}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} + \frac{2c^{2}k_{-}k_{+}^{2}}{(ck_{+}+k_{-})^{4}} + \frac{2ck_{-}^{2}k_{+}e^{-T(ck_{+}+k_{-})}}{(ck_{+}+k_{-})^{4}} + \frac{2ck$$

which in large T limit is

$$\langle m(T)^2 \rangle - \langle m(T) \rangle^2 = \left(\frac{k_+ c}{k_+ c + k_-}\right)^2 \frac{2}{(k_+ c + k_-)T} \frac{k_-}{k_+ c}.$$
 (A4)

Since in limit of large T ($T^{-1} << k_+c, k_-$)

$$1 = \frac{k_{-}}{k_{+}c} \frac{\langle m(T) \rangle}{1 - \langle m(T) \rangle}$$

$$\frac{\delta c}{c} = \frac{\delta \langle m(T) \rangle}{\langle m(T) \rangle (1 - \langle m(T) \rangle)}$$

$$= \sqrt{\left(\frac{k_{+}c}{k_{+}c + k_{-}}\right)^{2} \frac{1}{(k_{+}c + k_{-})T} \left[2\frac{k_{-}}{k_{+}c}\right] \frac{(k_{+}c + k_{-})^{2}}{k_{+}ck_{-}}}{k_{+}ck_{-}}}$$

$$= \sqrt{\frac{2}{k_{+}cT} + \frac{2}{k_{-}T}}.$$
(A5)

2. II

We next, consider calculations involved in evaluating the bound on gradient measuring time. We have to solve

$$D\nabla^2 c(r,\theta,\phi) = C\delta(r-a)\delta(\theta)$$
 (A6)

with the boundary condition

$$D\frac{\partial c(r,\theta)}{\partial r}|_{r=R} = 0.$$
 (A7)

The symmetry of the problem implies no ϕ dependence. One solution to above equation is

$$c(r,\theta) = \frac{C}{4\pi D\sqrt{r^2 + a^2 - 2ar\cos\theta}}.$$
 (A8)

In the region close to r = R, the above becomes

$$c(r,\theta) = \frac{C}{4\pi Da} \sum_{n=0,\infty} \frac{r^n}{a^n} P_n(\cos\theta), \qquad (A9)$$

where $P_n(\cos \theta)$ are the Legendre polynomials. $c(r, \theta) + b(r, \theta)$ is also a solution of Eq. (A6) if

$$\nabla^2 b(r,\theta) = 0. \tag{A10}$$

Solution of above equation for region r > R is

$$b(r,\theta) = \sum_{n=0,\infty} \frac{A_n}{r^{n+1}} P_n(\cos\theta).$$
(A11)

Hence for the region close to r = R, a solution is

$$c(r,\theta) = \frac{C}{4\pi Da} \sum_{n=0,\infty} \frac{r^n}{a^n} P_n(\cos\theta) + \sum_{n=0,\infty} \frac{A_n}{r^{n+1}} P_n(\cos\theta).$$
(A12)

The boundary condition

$$D\frac{\partial c(r,\theta)}{\partial r}|_{r=R} = 0$$
 (A13)

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implies

$$0 = \frac{C}{4\pi Da} \sum_{n=1,\infty} \frac{nR^{n-1}}{a^n} P_n(\cos\theta)$$
$$-\sum_{n=0,\infty} \frac{(n+1)A_n}{R^{n+2}} P_n(\cos\theta).$$
(A14)

Hence $A_0 = 0$ and for n > 0

$$A_n = \frac{C}{4\pi Da} \frac{nR^{2n+1}}{(n+1)a^n}.$$
 (A15)

Hence,

$$c(R,\theta) = \frac{C}{4\pi Da} \sum_{n=0,\infty} \frac{R^n}{a^n} P_n(\cos\theta)$$
$$+ \frac{C}{4\pi Da} \sum_{n=1,\infty} \frac{nR^{2n+1}}{(n+1)a^n} \frac{1}{R^{n+1}} P_n(\cos\theta)$$
$$= \frac{C}{4\pi Da} + \frac{C}{4\pi Da} \sum_{n=1,\infty} \frac{R^n}{a^n} P_n(\cos\theta).$$
(A16)

So

$$c(R, \theta = 0) = \frac{C}{4\pi Da} + \frac{C}{4\pi Da} \sum_{n=1,\infty} \frac{R^n}{a^n}$$

= $\frac{C}{4\pi Da} + \frac{C}{4\pi Da} \frac{R/a}{1 - R/a}$
= $\frac{C}{4\pi Da} \frac{1}{1 - R/a}$
 $c(R, \theta = \pi) = \frac{C}{4\pi Da} + \frac{C}{4\pi Da} \sum_{n=1,\infty} (-1)^n \frac{R^n}{a^n}$
= $\frac{C}{4\pi Da} - \frac{C}{4\pi Da} \frac{R/a}{1 + R/a}$
= $\frac{C}{4\pi Da} \frac{1}{1 + R/a}$. (A17)

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