

Chiral edge modes in evolutionary game theory: A kagome network of rock-paper-scissors cycles

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We theoretically demonstrate the realization of a chiral edge mode in a system beyond natural science. Specifically, we elucidate that a kagome network of rock-paper-scissors (K-RPS) hosts a chiral edge mode of the population density which is protected by the nontrivial topology in the bulk. The emergence of the chiral edge mode is demonstrated by numerically solving the Lotka-Volterra (LV) equation. This numerical result can be intuitively understood in terms of the cyclic motion of a single rock-paper-scissors cycle, which is analogous to the cyclotron motion of fermions. Furthermore, we point out that a linearized LV equation is mathematically equivalent to the Schrödinger equation describing quantum systems. This equivalence allows us to clarify the topological origin of the chiral edge mode in the K-RPS; a nonzero Chern number of the payoff matrix induces the chiral edge mode of the population density, which exemplifies the bulk-edge correspondence in two-dimensional systems described by evolutionary game theory.

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I. INTRODUCTION

Notion of topology plays a central role in condensed matter physics [1–7]. One of the remarkable properties of topological system is the emergence of edge states [8] protected by the topology in the bulk which is a source of anomalous behaviors. For instance, integer quantum Hall systems show the quantized Hall conductance with extremely high accuracy [1,9–11] due to the chiral edge mode (i.e., one-way propagating modes localized around the edge).

So far, insulators and superconductors have been extensively analyzed as platforms of topological physics. However, recently, it turned out that topological phenomena¹ [12,13] extend beyond the quantum systems [14–23]. In particular, the chiral edge modes protected by topological properties have been reported for various systems, such as photonic crystals [24–28], mechanical metamaterials [29–44], equatorial waves [45], active matter [46–48], and so on. This progress is significant not only from an academic viewpoint but also from an engineering viewpoint because such topologically protected chiral edge modes may result in new inventions, e.g., the topological insulator laser [49,50] and potential application to a novel energy transfer system of extremely low transmission loss.

In spite of the above progress, the chiral edge modes are still restricted to systems of natural science. Discovery of chiral edge modes beyond the natural science is crucial as it may provide a new perspective.

The aim of this paper is to report the discovery of a chiral edge mode in a system of evolutionary game theory which is

beyond natural science² [51–53]. Specifically, we elucidate that a chiral edge mode of the population density³ emerges in a kagome network of rock-paper-scissors (K-RPS)⁴ (see Fig. 1) due to the nontrivial topology in the bulk. On each site of the K-RPS, one of the strategies, rock (R), paper (P), and scissors (S), is assigned. The payoff of a player on a given site is described by the arrows of bonds connecting the sites. The emergence of the chiral edge mode is demonstrated by numerically solving the Lotka-Volterra (LV) equation. This result can be intuitively understood by focusing on cyclic motion of the single rock-paper-scissors (RPS) cycle which is analogous to the cyclotron motion of fermions under a uniform magnetic field. Furthermore, we elucidate the topological origin of the chiral edge modes by pointing out the mathematical equivalence of a linearized LV equation of the K-RPS and the Schrödinger equation of a fermionic kagome lattice model with the nontrivial topology (see Fig. 4). This equivalence elucidates the bulk-edge correspondence for two-dimensional systems described by evolutionary game theory; a nonzero Chern number of the payoff matrix induces the chiral edge modes observed in the time evolution of the population density of players located at the sites of the K-RPS.

Topological band structure and a zero mode have been discussed for a RPS chain [23]. We would like to stress, however, that the emergence of the chiral edge modes has remained an open question because Ref. [23] analyzes the one-dimensional system. Novelty of our work is discovering the topological origin of the one-way mode by pointing out the mathematical equivalence of the linearized LV equation and the Schrödinger equation.

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¹References [12,13] mentions topology of game theory. We note, however, that topology discussed in these references differs from the topology of eigenvectors (or eigenstates). One can characterize the topology of the eigenstates by topological invariant which is computed from eigenvectors.

²Evolutionary game theory can be applied not only to natural science but also to social science [51–53].

³For definition of the population density, see just below of Eq. (2).

⁴Namely, players choose the given strategy assigned at each site as Fig. 1.

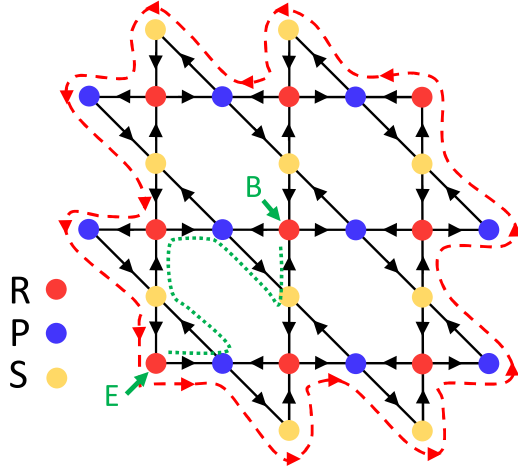


FIG. 1. Kagome network of rock-paper-scissors. At site I strategy s_I is assigned as illustrated by a colored dot; red, blue, and yellow dots denote strategies, rock (R), paper (P), scissors (S), respectively. In this panel, the arrows connecting sites illustrate the payoffs; R beats S; S beats P; P beats R. [The explicit form of the payoffs is shown in Eq. (1) for a single RPS cycle.] This kagome network of rock-paper-scissors is described by evolutionary game theory rather than by solid state physics. However, it can be mapped to a spinless fermion model (see Fig. 4).

II. CYCLOTRON MOTION IN A SINGLE RPS CYCLE

As a first step, we show that “cyclotron motion” can be observed in a single RPS cycle [54], which plays a key role in the emergence of chiral edge modes beyond natural science.

Consider two players who choose one of the strategies $(s_1, s_2, s_3) = (R, P, S)$. The rule of the game is illustrated in Fig. 2(a); R beats S; S beats P; P beats R. In this case, the payoff of a player is A_{IJ} if the player chooses the strategy s_I and the other player choosing s_J ($I, J = 1, 2, 3$). Here the payoff matrix of this game is given by

$$A = \begin{pmatrix} 0 & -1 & 1 \\ 1 & 0 & -1 \\ -1 & 1 & 0 \end{pmatrix}. \quad (1)$$

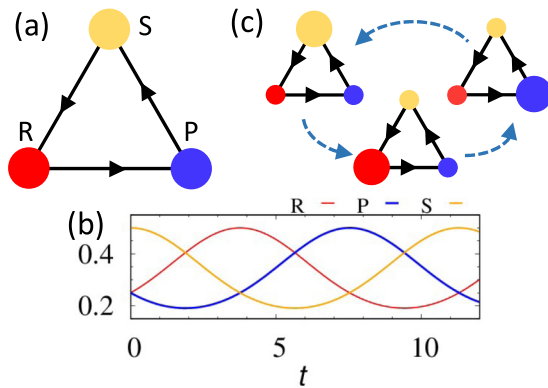


FIG. 2. (a) Sketch of a RPS cycle. The arrows in (a) denote the dominance relationship between the strategies; R beats S; S beats P; and P beats R. (b) Dynamics of the RPS cycle with an initial condition $\mathbf{x}_{\text{ini}} = (\frac{1}{4}, \frac{1}{4}, \frac{1}{2})$. (c) Sketch of the dynamics of the RPS cycle. The size of the circle indicates the population density.

Now, let us consider the case where a large number of players repeat the game. It is known that the dynamics of this game is described by the LV equation⁵

$$\partial_t x_I = x_I \mathbf{e}_I^T A \mathbf{x}, \quad (2)$$

with $\mathbf{x} = (x_1, x_2, x_3)^T$ and x_1, x_2 , and x_3 being the population density⁶ of players who choose the strategies R, P, and S, respectively [23,55]. We suppose that the vector \mathbf{x} is normalized $\sum_I x_I = 1$ unless otherwise noted. The vector \mathbf{e}_I is the unit vector whose I th element takes one; $[\mathbf{e}_I]_J = \delta_{IJ}$. When a player chooses the strategy s_I , the expectation value of the payoff is written as $\mathbf{e}_I^T A \mathbf{x}$. Thus Eq. (2) indicates that the population density x_I increases in order to enhance the payoff.

We note that the vector $\mathbf{c} = (1, 1, 1)^T/3$ satisfies

$$A \mathbf{c} = 0, \quad (3)$$

which means that $\mathbf{x} = \mathbf{c}$ is a stationary state. This vector describes a Nash equilibrium [56,57]⁷ for a classical game theory where the strategy of players does not change over time and players play the game only once. For the relation $A \mathbf{c} = 0$, the following fact is essential: At each site, the number of bonds with the out-going arrow is equal to the number of bonds with the in-coming arrow. Because of this fact, for arbitrary I , the expectation value $\mathbf{e}_I^T A \mathbf{c}$ is zero, which results in the relation $A \mathbf{c} = 0$. Noting these facts, we can also find the stationary state for the K-RPS.

A slight deviation from the stationary state \mathbf{c} results in cyclic motion which is analogous to the cyclotron motion of fermions in the Landau levels. The time evolution with the initial state $\mathbf{x}_{\text{ini}} = (\frac{1}{4}, \frac{1}{4}, \frac{1}{2})$ is shown in Fig. 2(b). Here, we choose $A_{21} = 1$ as a unit of time [see Eq. (2)]. As illustrated in Fig. 2(c), the data shown in Fig. 2(b) indicate the cyclic motion, which can be intuitively understood by noticing that the population density propagates along the arrows illustrated in Fig. 2(a). Namely, with the given initial state, $\mathbf{x}_{\text{ini}} = (\frac{1}{4}, \frac{1}{4}, \frac{1}{2})$, players who choose strategy R gain the highest payoff (i.e., $\mathbf{e}_I^T A \mathbf{x}$ become maximum for $I = 1$). Thus, after time evolution, the population density of players choosing strategy R increases. In a similar manner, we can see that the population density of players choosing strategy P increases in the next step.

The above results elucidate that the players of the single RPS cycle mimic the “cyclotron motion” as if they were fermions in the Landau levels. We note that analogy of cyclotron motion can be mathematically shown (see Appendix B). As we see in Sec. III, the “cyclotron motion”

⁵If the players know the prediction and try to behave in a different way, the population density may not follow this equation. However, so far, it is assumed that Eq. (2) holds. Namely, it has been assumed that the increase/decrease of the population density corresponds to the payoff [55].

⁶As mentioned in the main text, we normalize the vector \mathbf{x} so that $\sum_I x_I = 1$ holds. In this sense, x_I denotes the population density

⁷It should be noted that the Nash equilibrium does not necessarily corresponds to evolutionary stable strategy. Indeed, it is known that there is no evolutionary stable strategy in the single-rock-paper-scissors cycle [56,57].

in the single RPS plays an important role to search a system exhibiting a chiral edge mode.

III. DYNAMICS OF THE K-RPS

A typical example of quantum systems exhibiting chiral edge modes is an integer quantum Hall system. Fermions in this two-dimensional system show the cyclotron motion which breaks time-reversal symmetry. Keeping this fact in mind, one can expect the emergence of chiral edge modes (i.e., one-way propagating the population density localized around the edge) in a two-dimensional network constructed from the RPS cycles which mimic the cyclotron motion. Specifically, we consider the K-RPS illustrated in Fig. 1 which indeed hosts a chiral edge modes.

A. Numerical results

By solving the LV equation (2) numerically, we demonstrate the presence of chiral edge modes in the K-RPS. The payoff matrix $A \in M(N_{\text{tot}}, \mathbb{R})$ can be read off from the arrow assigned to each bond (see Fig. 1). Here N_{tot} denotes the number of sites. In order to analyze the time evolution, we employ a fourth order Runge-Kutta method [58]. We discretize time as $t_n = n\Delta t$ with $\Delta t = 0.1$ and $n = 0, 1, 2, \dots$

We analyze the dynamics with an initial state,⁸ which slightly deviates from a stationary state $\mathbf{c}^{(K)} = (1, 1, 1, \dots, 1)^T / N_{\text{tot}}$ ($\mathbf{x} = \mathbf{c}^{(K)} + \delta\mathbf{x}$). The relation $A\mathbf{c}^{(K)} = 0$ holds because the number of bonds with the in-coming arrow is equal to the number of bonds with the out-going arrow for each site of the K-RPS [see Fig. 1 and the argument below Eq. (3)].

Figure 3 clearly indicates the presence of the chiral edge mode; sites of the high population density (red dots) propagate along the edge in the counterclockwise direction. Our numerical results also imply topological stability of the chiral edge mode. In Appendix A, we demonstrate the following two facts. (i) The system hosts the chiral modes even with an initial condition significantly deviating from $\mathbf{c}^{(K)}$ (see Fig. 7). (ii) Even in the presence of a defect on the edge, the chiral edge mode propagates by detouring around the defect (see Fig. 8). The latter result is particularly counterintuitive. The above results indicate the robustness of the chiral edge mode.

B. Intuitive discussion

The dynamics obtained in Fig. 3 is intuitively understood as follows.

Firstly, we recall that the K-RPS is composed of the single RPS cycle. This fact means that the population density propagates along the out-going arrows in order to maximize the payoff. For instance, when the population density at site denoted with “E” is higher than the other sites (see Fig. 1), it propagates around the path illustrated in the red dashed line in Fig. 1, which implies the presence of the chiral edge modes.

⁸We have also analyzed evolution of an initial state which describes a population density deviating from $N_{\text{tot}}\mathbf{c}^{(K)}$ only at the center of the system. In this case, we observe that the deviation propagates homogeneously.

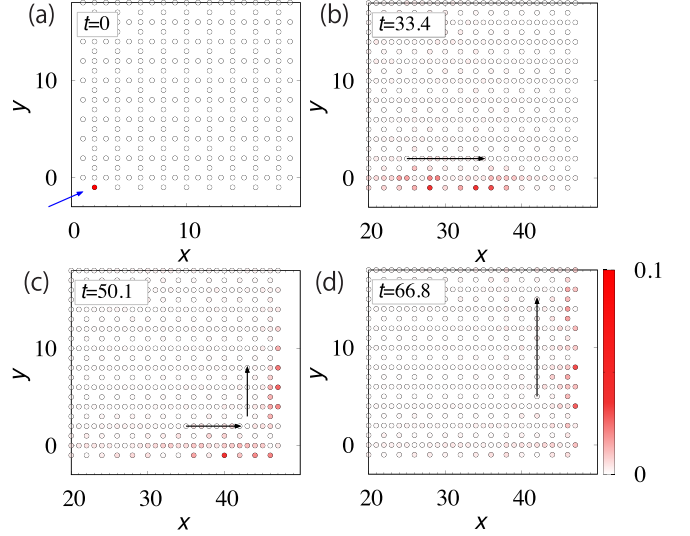


FIG. 3. Time evolution of the population density for the K-RPS. In these figures, the absolute value of the deviation $|\delta\mathbf{x}|$ ($\delta\mathbf{x} = \mathbf{x} - N_{\text{tot}}\mathbf{c}^{(K)}$) is plotted; the vector \mathbf{x} satisfies $\sum_i x_i = N_{\text{tot}} + 0.1$. For $t = 0$, δx_i takes 0.1 at the site denoted by the blue arrow in (a); otherwise δx_i is zero. The deviation of the population density propagates along the edge in the counterclockwise direction as illustrated by the black arrows in (b)–(d). We have simulated the time evolution up to $t_{\text{max}} = 200$. For $0 \leq t \leq t_{\text{max}}$, the chiral edge mode is observed. More detailed data are shown in Fig. 6 of Appendix A. The scale of color plot in (a)–(c) is the same as the one in (d).

The localization of the chiral edge mode can be deduced as follows. Firstly, we note that in the bulk, sites are connected by four bonds; out-going arrows are assigned to two of the bonds, and in-coming arrows are assigned to the other two bonds. Thus, in contrast to the players on the edge sites, those in the bulk have two options of out-going arrows which results in the localization of the chiral mode around the edge. For instance, when the population density at site “E” on the edge propagates to site “B” in the bulk along the path denoted with green dashed-line in Fig. 1, it passes through five branches. Because the propagation is suppressed at each of these branches, the deviation of the population density is localized around the edges.

The above results explicitly demonstrate the emergence of the chiral edge mode in the K-RPS. As we see in Sec. IV B, the topology in the bulk governs the chiral edge modes. This is supported by a simulation shown in Fig. 9 of Appendix A where the payoff matrix is flipped, $A \rightarrow -A$, (Namely, we impose the opposite rule instead of the ordinary one; R beats P; P beats S; S beats R). The obtained data show the edge mode propagating the opposite direction.

IV. TOPOLOGICAL CHARACTERIZATION OF THE CHIRAL EDGE MODES OF THE K-RPS

So far, we have seen that the K-RPS hosts a chiral edge mode. Here, we elucidate the topological origin by pointing out a relation between the K-RPS (Fig. 1) and the fermionic quantum model (Fig. 4) with the nontrivial topology.

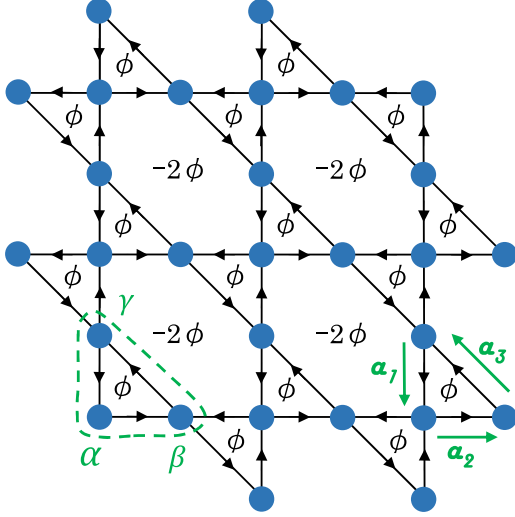


FIG. 4. Sketch of fermionic lattice model with magnetic fluxes introduced in Ref. [59]. a spinless fermion acquires a phase ϕ when it hops around a triangle along the arrows. The K-RPS shown in Fig. 1 is mapped to the fermionic model with $\phi = 3\pi/2$.

A. A linearized LV equation and the Schrödinger equation

In order to see the relation between the K-RPS and the fermionic kagome lattice model, we linearize the LV equation around the stationary state $\mathbf{c}^{(K)}$.

For \mathbf{x} slightly deviating from $\mathbf{c}^{(K)} = (1, 1, 1, \dots, 1)^T / N_{\text{tot}}$, the LV equation is rewritten as

$$\begin{aligned} \partial_t \mathbf{e}_I \cdot \delta \mathbf{x} &= \mathbf{e}_I \cdot (\mathbf{c}^{(K)} + \delta \mathbf{x}) \mathbf{e}_I^T A(\mathbf{c} + \delta \mathbf{x}) \\ &= (\mathbf{c}^{(K)} + \delta \mathbf{x}) \cdot \mathbf{e}_I \mathbf{e}_I^T A(\delta \mathbf{x}) \\ &\sim (\mathbf{c}^{(K)})^T P_I A \delta \mathbf{x}, \end{aligned} \quad (4)$$

with $P_I = \mathbf{e}_I \mathbf{e}_I^T$. From the first to the second line, we have used the relation $A\mathbf{c}^{(K)} = 0$. In the last line, we have discarded the second order term of $(\delta \mathbf{x})^2$.

Noting the relation $N_{\text{tot}} P_I \mathbf{c}^{(K)} = \mathbf{e}_I$, we can write the above linearized equation as

$$i \partial_t \delta \mathbf{x} = \frac{1}{N_{\text{tot}}} H \delta \mathbf{x}, \quad (5)$$

with a Hermitian matrix $H = iA$, which is mathematically equivalent to the Schrödinger equation up to the prefactor.

Taking into account the explicit form of the payoff matrix A , we can see that the K-RPS is mapped to the tight-binding model of spinless fermion in the kagome lattice which hosts a chiral edge mode due to the nontrivial topology in the bulk. To see this more clearly, let us discuss the following eigenvalue problem which governs the dynamics described by the linearized LV equation (5):

$$\sum_j A_{ij}(\mathbf{k}) \psi_{jn} = \psi_{in}(\mathbf{k}) \epsilon_n, \quad (6a)$$

$$A(\mathbf{k}) = \begin{pmatrix} 0 & -(1 + e^{i2\mathbf{k} \cdot \mathbf{a}_2}) & (1 + e^{-i2\mathbf{k} \cdot \mathbf{a}_1}) \\ (1 + e^{-i2\mathbf{k} \cdot \mathbf{a}_2}) & 0 & -(1 + e^{i2\mathbf{k} \cdot \mathbf{a}_3}) \\ -(1 + e^{i2\mathbf{k} \cdot \mathbf{a}_1}) & (1 + e^{-i2\mathbf{k} \cdot \mathbf{a}_3}) & 0 \end{pmatrix}, \quad (6b)$$

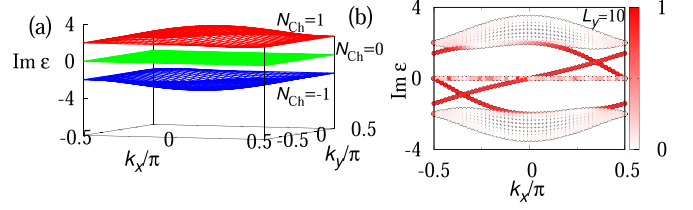


FIG. 5. (a) Band structure of “Bloch Hamiltonian” $A(\mathbf{k})$. The Chern number takes $N_{\text{Ch}} = -1, 0$, and 1 for the bottom, middle, and top bands, respectively. (b) Spectrum of the K-RPS under the cylinder geometry. Color of data points represents the absolute value of the polarization $|P_n|$ defined in Eq. (8). Polarization for each eigenstate is plotted in Fig. 10 of Appendix C.

where $A(\mathbf{k})$ with $\mathbf{k} = (k_x, k_y)$ is the Fourier transformed payoff matrix. Here, ψ_{jn} ($j, n = 1, 2, 3$) denotes the j th component of the eigenvector ψ_n with the eigenvalue ϵ_n . The vectors connecting the neighboring site are defined as $\mathbf{a}_1 := (0, -1)$, $\mathbf{a}_2 := (1, 0)$, and $\mathbf{a}_3 := (-1, 1)$ (see also Fig. 4). Because $A(\mathbf{k})$ is anti-Hermitian [$A(\mathbf{k}) = -A^\dagger(\mathbf{k})$], the eigenvalues are pure imaginary ($\epsilon_n \in i\mathbb{R}$). We note that the Hermitian matrix $H(\mathbf{k}) = iA(\mathbf{k})$ is identical to the Bloch Hamiltonian of spinless fermions in the kagome lattice with magnetic fluxes $\phi = 3\pi/2$ (for details of the fermionic system see Appendix B). The symmetry class of $H(\mathbf{k})$ is class D where particle-hole symmetry is preserved [60–62].

The above results clarify the mathematical equivalence of the linearized LV equation of the K-RPS and the Schrödinger equation of the fermionic kagome lattice model with the nontrivial topology.⁹ This equivalence and the bulk-edge correspondence [8] provides a topological perspective of the chiral edge modes of the K-RPS whose details are discussed below.

B. Topological characterization of the chiral edge modes of the K-RPS

Based on the above results, we address the topological characterization of the chiral edge mode in the K-RPS.

Specifically, we characterize the edge mode by the following steps. Firstly, we analyze the bulk band structure as well as the Chern number [see Eq. (7)]. Secondly, by diagonalizing the system under the cylinder geometry, we demonstrate that the topological properties in the bulk induce the chiral edge mode, which is known as the bulk-edge correspondence [8] in the context of topological insulators/superconductors. Here, the system under the cylinder geometry denotes the system where the periodic (open) boundary condition is imposed along the x (y) direction.

Firstly, we discuss the bulk properties by diagonalizing the Fourier transformed payoff matrix $A(\mathbf{k})$ which is anti-Hermitian. Figure 5(a) plots the band structure of $A(\mathbf{k})$. Because each band is separated by a gap in the

⁹We stress that the above argument elucidates just the mathematical equivalence of the linearized LV equation and the Schrödinger equation. Thus the equivalence does not mean that the population density itself obeys the Schrödinger equation.

two-dimensional Brillouin zone, the Chern number of each band is quantized which is defined as

$$N_{\text{Ch}} = \sum_{\mu\nu} \epsilon_{\mu\nu} \int \frac{dk_x dk_y}{2\pi i} \partial_{k_\mu} \mathcal{A}_{\nu\nu}, \quad (7a)$$

$$\mathcal{A}_{n\mu} = \sum_j \psi_{nj}^\dagger(\mathbf{k}) \partial_{k_\mu} \psi_{jn}(\mathbf{k}), \quad (7b)$$

Here, $\epsilon_{\mu\nu}$ ($\mu, \nu = x, y$) denotes the anti-symmetric tensor satisfying $\epsilon_{xy} = 1$.

Employing the method based on Ref. [63], we find that the Chern number of the bottom (top) band takes -1 (1) [see Fig. 5(a)].

The nontrivial topology characterized by the Chern number in the bulk predicts the chiral edge modes around the boundary. Indeed, the spectrum of the matrix A under the cylinder geometry exhibits the chiral edge modes [see Fig. 5(b)]. In Fig. 5(b), the color assigned to each eigenvalue ϵ_n denotes the absolute value of the polarization $|P_n|$ of the corresponding eigenvector $\psi_n(k_x)$

$$P_n = 1 - \frac{2}{L_y} \sum_{J_y} \psi_{nJ_y}^\dagger(k_x) J_y \psi_{J_y n}(k_x), \quad (8)$$

where J_y ($J_y = 1, 2, \dots, L_y$) labels the sites along the y direction. Figure 5(b) indicates that the chiral modes denoted by red-colored dots are localized around edges (for more details, see Fig. 10 of Appendix C). The replacement $A \rightarrow -A$ flips the sign of Chern number. Correspondingly, the direction of the chiral edge mode changes as mentioned in Sec. III A (see Fig. 9 in Appendix A). These behaviors correspond to the behaviors of the fermionic lattice model (Fig. 4) where the magnetic fluxes are flipped ($\phi \rightarrow -\phi$).

The above results elucidate that the chiral edge mode of the population density in the K-RPS is induced by the nontrivial topology of the bulk, which exemplifies the bulk-edge correspondence¹⁰ for two-dimensional systems described by evolutionary game theory [8]; the density profile of player propagates as a wave in the counter-clockwise direction along the edge due to $N_{\text{Ch}} = 1$.

We note that previous work [23] has studied the dynamics in a one-dimensional system. However, the chiral edge mode protected by the nontrivial topology in the bulk has not been reported so far.

V. SUMMARY

In this paper, we have discovered the emergence of the chiral edge mode beyond natural science.

Specifically, we have elucidated that the K-RPS hosts a chiral edge mode of the population density which is protected by the nontrivial topology in the bulk. The emergence of the chiral edge mode is demonstrated by numerically solving the LV equation. The dynamics is also intuitively understood by focusing on the “cyclotron motion” of the single RPS cycle.

¹⁰The Chern number computed for the periodic boundary conditions predicts the existence of the chiral edge modes with a boundary, which is known as the bulk-edge correspondence [8].

Furthermore, we have also elucidated the topological origin of the chiral edge mode by mapping the K-RPS to the fermionic kagome lattice model. The former (latter) is described by evolutionary game theory (quantum mechanics). By making use of the above mapping, we have found that due to the nonzero Chern number of the payoff matrix in the bulk, the chiral edge mode of the population density emerges regardless of the other details of the K-RPS, which exemplifies the bulk-edge correspondence in two-dimensional systems described by evolutionary game theory. We note that the chiral edge mode induced by bulk topology should also be observed in other systems because its emergence depends only on the nontrivial topology in the bulk [64].

We finish this paper with two remarks. We note that a topological band structure in a one-dimensional chain of RPS has been analyzed in Ref. [23]. As well as the discovery of the chiral edge mode, the novelty of this paper is elucidating the mathematical equivalence of the linearized LV equation and the Schrödinger equation clarifying the topological origin of the one-way propagating mode observed in the time evolution. We also note that the experimental observation of the chiral edge mode is a significant open question to be addressed. Because RPS cycles have been reported for a wide variety of systems, e.g., a system of bacteria [65], and human societies [53,66,67], we expect the observation of chiral modes in such systems.

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APPENDIX A: DETAILS OF TIME EVOLUTION FOR THE K-RPS

The time evolution of the population density. In Fig. 3, the sign of deviation is discarded. For the complementarily information, we plot the deviation $\delta\mathbf{x}$ ($\delta\mathbf{x} = \mathbf{x} - N_{\text{tot}}\mathbf{e}^{(K)}$) in Fig. 6.

Time evolution for the K-RPS with other conditions. Here, we discuss the dynamics of the K-RPS with three distinct cases; (i) the dynamics with an initial condition significantly deviating from $N_{\text{tot}}\mathbf{e}^{(K)}$, (ii) the dynamics of the system with a defect on the edge, and (iii) the dynamics of the system where the payoff matrix is replaced as $A \rightarrow -A$.

Firstly, we discuss the dynamics of the K-RPS [see Fig. 1(a)] with an initial condition significantly deviating from $N_{\text{tot}}\mathbf{e}^{(K)}$. In Fig. 7, we can find a one-way propagating mode localized at the edge even when the initial condition is significantly deviates from $N_{\text{tot}}\mathbf{e}^{(K)}$; the propagation is denoted by the arrows.

Secondly, we demonstrate that even in the presence of a defect on the edge, the chiral edge mode propagates along the edge by detouring around the defect. Figure 8 plots the

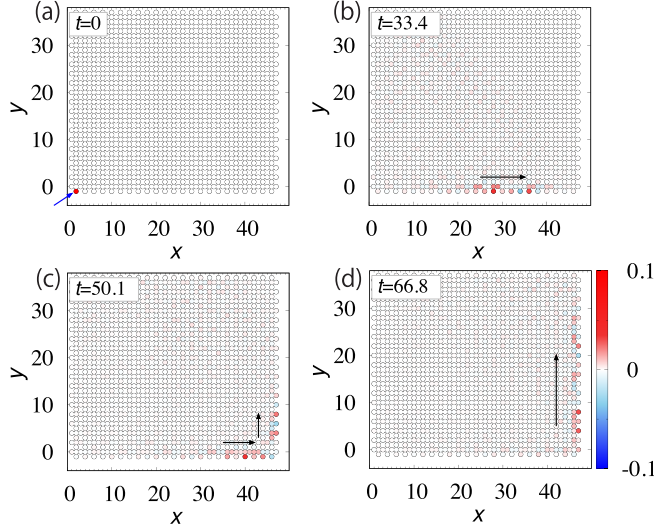


FIG. 6. Time evolution of the population density for the K-RPS. In these figures, the deviation from $N_{\text{tot}}\mathbf{e}^{(K)}$ ($\delta\mathbf{x} = \mathbf{x} - N_{\text{tot}}\mathbf{e}^{(K)}$) is plotted; the vector \mathbf{x} satisfies $\sum_i x_i = N_{\text{tot}} + 0.1$. For $t = 0$, δx_i takes 0.1 at the site denoted by the blue arrow in (a); otherwise δx_i is zero. The deviation of the population density propagates along the edge in the counterclockwise direction as illustrated by the black arrows in (b)–(d). The scale of color plot in (a)–(c) is the same as the one in (d). The data, which are more suited for printing in gray scale, are provided in Fig. 3.

deviation of the population density δx_i around the defect. This figure indicates that the population density detours around the defect, which supports the robustness of the chiral edge mode.

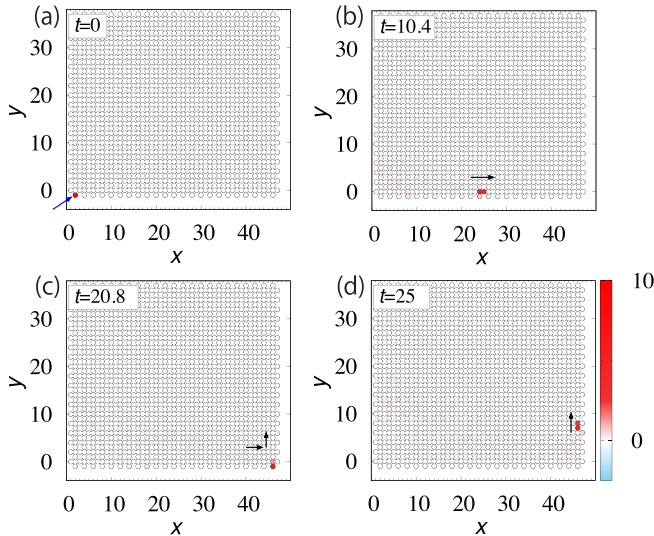


FIG. 7. Time evolution of the population density for the K-RPS. In each panel, the deviation from $N_{\text{tot}}\mathbf{e}^{(K)}$ ($\delta\mathbf{x} = \mathbf{x} - N_{\text{tot}}\mathbf{e}^{(K)}$) is plotted; the vector \mathbf{x} satisfies $\sum_i x_i = N_{\text{tot}} + 9$. For $t = 0$, δx_i takes 9 at the site denoted by the blue arrow in panel (a); otherwise, δx_i takes zero. The deviation of the population density is positive which propagates along the edge in the counter-clockwise direction as indicated by the black arrows in (b)–(d). The scale of color plot in (a)–(c) is the same as the one in (d).

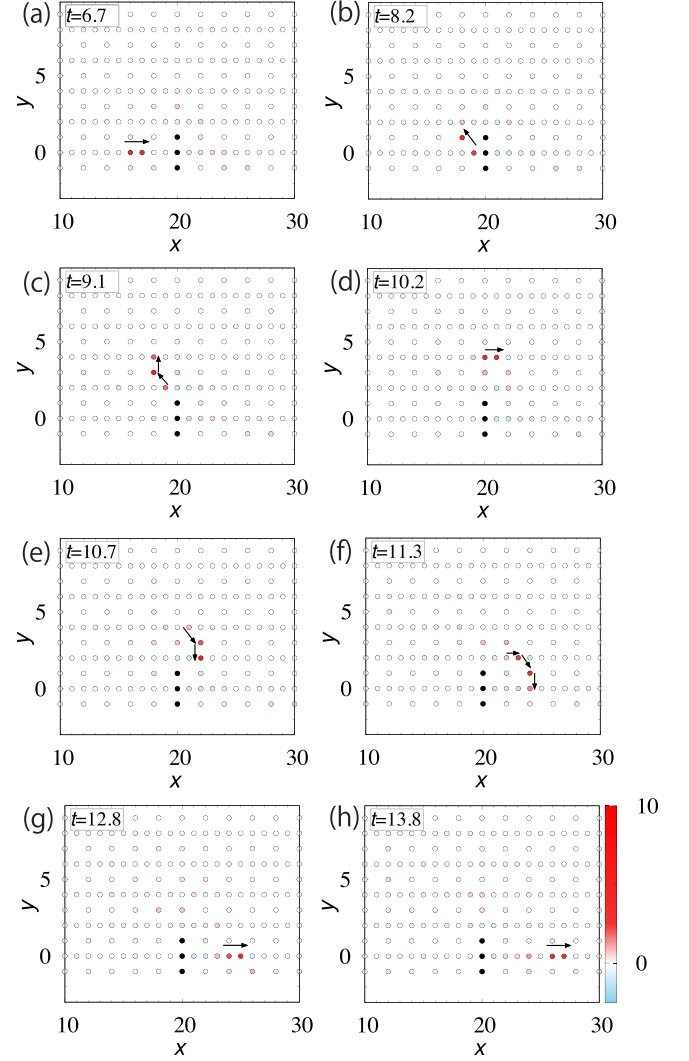


FIG. 8. Time evolution of the population density for the K-RPS with a defect on the edge. Sites denoted with black dots are isolated from the other sites, which serves as a defect. Except for the presence of the defect, the set up and the initial condition are the same as those of Fig. 7. The deviation of the population density is positive whose propagation is denoted by arrows. The scale of color plot in panels (a)–(g) is the same as the one in panel (e).

Thirdly, we discuss the effect of the replacement $A \rightarrow -A$ on the dynamics. As we can see in Fig. 9, the replacement $A \rightarrow -A$ flips the chirality of the edge mode, which is consistent with the fact that the replacement A to $-A$ flips the sign of the Chern number.

APPENDIX B: SPINLESS FERMIONS IN A KAGOME LATTICE WITH MAGNETIC FLUXES

The Hamiltonian of spinless fermions in the kagome lattice has been introduced in Ref. [59]. However, in order to make this paper self-contained, we briefly describe the model.

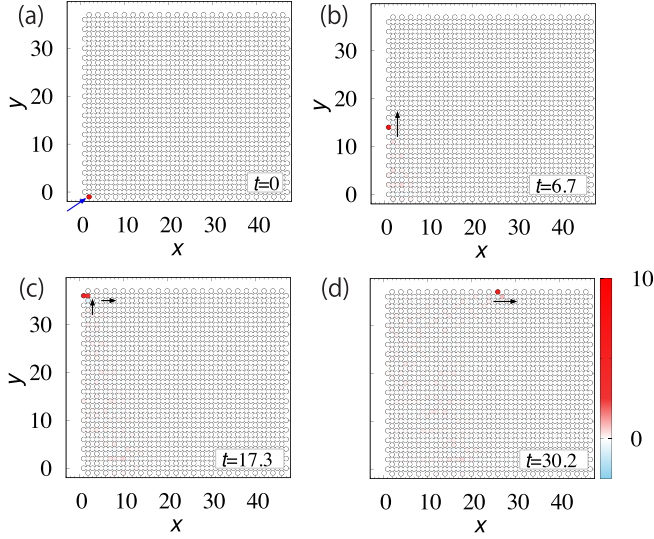


FIG. 9. Time evolution of the population density for the K-RPS. Each panel is plotted in a similar way as Fig. 7. The deviation of the population density is positive whose propagation in the clockwise direction is denoted by black arrows in (b)–(d). The scale of color plot in (a)–(c) is the same as the one in (d).

Consider spinless fermions in the kagome lattice with magnetic fluxes. The Hamiltonian reads

$$H_{\text{fermi}} = \sum_{\langle ij \rangle} t_{ij} d_i^\dagger d_j, \quad (\text{B1})$$

where d_j^\dagger (d_j) creates (annihilates) a spinless fermion at site j . We note that i appearing as subscripts specifies sites. The hopping integral t_{ij} ($t_{ij} = t_{ji}^*$) describes hopping from site j to site i which takes $t_{ij} = t_0 e^{i\phi}$ when it describes hopping apparet to arrows in Fig. 4. The summation is taken over neighboring sites.

Applying the Fourier transformation, Eq. (B1) is rewritten as

$$H_{\text{fermi}} = \sum_{\mathbf{k}} d_{\mathbf{k}}^\dagger h(\mathbf{k}) d_{\mathbf{k}}, \quad (\text{B2a})$$

with

$$h(\mathbf{k}) = \begin{pmatrix} 0 & e^{-i\phi/3}(1 + e^{2ik \cdot \mathbf{a}_2}) & e^{i\phi/3}(1 + e^{-2ik \cdot \mathbf{a}_1}) \\ \text{H.c.} & 0 & e^{-i\phi/3}(1 + e^{2ik \cdot \mathbf{a}_3}) \\ \text{H.c.} & \text{H.c.} & 0 \end{pmatrix}, \quad (\text{B2b})$$

$$\mathbf{d}_{\mathbf{k}} = (d_{\alpha\mathbf{k}} \quad d_{\beta\mathbf{k}} \quad d_{\gamma\mathbf{k}})^T, \quad (\text{B2c})$$

and $d_{s\mathbf{k}} = \frac{1}{\sqrt{N_{\text{UC}}}} \sum_{\mathbf{R}_j} e^{-ik \cdot (\mathbf{R}_j)} d_j$. Here, s labels sublattices ($s = \alpha, \beta, \gamma$), and \mathbf{R}_j denotes a position of a unit cell

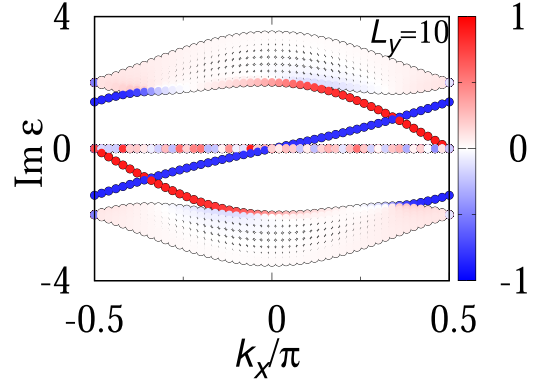


FIG. 10. Spectrum of the K-RPS under the cylinder geometry. Color of data points represents the polarization P_n defined in Eq. (8). The data, which are more suited for printing in gray scale, are provided in Fig. 5(b).

including site j . For $\phi = 3\pi/2$, Eq. (B2b) is reduced to $iA(\mathbf{k})$ with $A(\mathbf{k})$ defined in Eq. (6b).

We note that Eq. (B2b) is equivalent to Eq. (5) of Ref. [59], which can be seen by applying the unitary transformation,

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{-ik \cdot \mathbf{a}_2} & 0 \\ 0 & 0 & e^{ik \cdot \mathbf{a}_1} \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}. \quad (\text{B3})$$

We note that the K-RPS can be mapped to the above fermionic system regardless of the number of unit cells. Indeed a single-RPS can be mapped to a fermionic three-site system forming a loop. The Hamiltonian reads

$$H_{3\text{sites}} = -i \begin{pmatrix} 0 & -1 & 1 \\ 1 & 0 & -1 \\ -1 & -1 & 0 \end{pmatrix}, \quad (\text{B4})$$

describing cyclotron motion of a spinless fermion in this three-site system. The above result elucidate mathematical equivalence between the cyclic motion in the single-RPS and the cyclotron motion in the fermionic three-sites system. To be more specific, the linearized LV equation for the single-RPS is mathematically equivalent to the Schrödinger equation for the Hamiltonian $H_{3\text{sites}}$.

APPENDIX C: POLARIZATION OF EACH EIGENSTATE

The polarization P_n of each eigenstate is shown in Fig. 10. Noting that P_n takes $P_n = -1$ ($P_n = 1$) when the state $\psi_n(k_x)$ is localized at $J_y = 1$ ($J_y = L_y$) [see Eq. (8)], we can see that the chiral mode denoted by blue-colored (red-colored) dots is localized around $J_y = 1$ ($J_y = L_y$) (see Fig. 10).

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