

Randall-Sundrum braneworld in modified gravityHiroshi Nakada^{1,*} and Sergei V. Ketov^{1,2,3,4,†}¹*Department of Physics, Tokyo Metropolitan University,
Minami-ohsawa 1-1, Hachioji-shi, Tokyo 192-0397, Japan*²*Kavli Institute for the Physics and Mathematics of the Universe (IPMU),
The University of Tokyo, Chiba 277-8568, Japan*³*Department of Physics, Faculty of Science, Chulalongkorn University,
Thanon Phayathai, Pathumwan, Bangkok 10330, Thailand*⁴*Institute of Physics and Technology, Tomsk Polytechnic University,
30 Lenin Avenue, Tomsk 634050, Russian Federation*

(Received 23 July 2016; published 3 November 2016)

We modify the Randall-Sundrum model of braneworld (with two branes) by adding the scalar curvature squared term in five dimensions. We find that it does not destabilize the Randall-Sundrum solution to the hierarchy problem of the standard model in particle physics.

DOI: [10.1103/PhysRevD.94.103503](https://doi.org/10.1103/PhysRevD.94.103503)**I. INTRODUCTION**

The basic idea of the braneworld is an assumption that our world is confined to a brane in a higher-dimensional space-time. To the best of our knowledge, this idea was first proposed in 1983 by Rubakov and Shaposhnikov [1] who studied a confinement of matter on a domain wall in a five-dimensional space-time. In 1998, Arkani-Hamed *et al.* [2] found that the braneworld may be the alternative to (Kaluza-Klein) space-time compactification in solving the hierarchy problem of particle physics.

Two elegant and simple braneworld models were proposed by Randall and Sundrum in 1999 [3]. The first of them (RSI) has two branes embedded in a five-dimensional space-time with a negative cosmological constant. One brane is identified with our visible Universe, whereas another brane is hidden, similarly to the Horava-Witten model in string theory [4]. The RSI braneworld model generates a large (exponential) hierarchy between the electroweak scale and the Planck scale by using the geometrical warp factor in a higher-dimensional space-time metric, with a small extra dimension forming an orbifold S^1/Z_2 (see, e.g., [5] for a review).

In braneworld scenarios, extra dimensions may be large and even infinite, while only gravity can be present there. The braneworld idea survives all experimental checks. However, there is no experimental evidence for the existence of extra dimensions either.

The success of braneworld poses the question of its stability against possible modifications of Einstein gravity. Such modifications are inevitable because the Einstein gravity is known to be nonrenormalizable, either in four or higher space-time dimensions. In this paper we answer this question by modifying gravity in five space-time

dimensions via adding the simplest higher-order term proportional to the scalar curvature squared, in the context of the RSI braneworld model.

We use the natural units, $c = \hbar = 1$, and the five-dimensional space-time signature $\eta_{\mu\nu} = \text{diag}(-, +, +, +, +)$. Capital latin indices refer to five space-time dimensions, while lowercase greek indices refer to four dimensions of a brane.

Our gravitational action in five dimensions reads

$$S_{\text{gr}} = \frac{1}{2\kappa^2} \int d^5x \sqrt{-g} (R + \alpha R^2 - \Lambda), \quad (1)$$

where we have introduced the gravitational coupling constant κ , the space-time scalar curvature R , the modified gravity coupling constant α , and a cosmological constant Λ .

As in the RSI model [3], we introduce a visible brane (where SM fields are confined) located at $y = 0$ with tension τ_{RS} , and a hidden brane located at $y = L$ and having the negative tension $-\tau_{\text{RS}}$, where the coordinate y parametrizes the extra dimension with the orbifold identification $-y \sim y$. Accordingly, our full action is given by

$$S = S_{\text{gr}} + S_{\text{vis}} + S_{\text{hid}}, \quad (2)$$

where

$$S_{\text{hid}} = -\tau_{\text{RS}} \int d^4x dy \sqrt{-g} \delta(y - L), \quad (3)$$

$$S_{\text{vis}} = \tau_{\text{RS}} \int d^4x dy \sqrt{-g} \delta(y). \quad (4)$$

Our paper is organized as follows. In Sec. II we perform the duality transformation converting the modified gravity into a scalar-tensor gravity, which greatly simplifies our treatment, because it allows us to get rid of the higher derivatives in the action (1). The scalar potential is studied in Sec. III. In Sec. IV we derive the braneworld equations

*nakada-hiroshi1@ed.tmu.ac.jp

†ketov@tmu.ac.jp

for the warp factor coupled to the extra dynamical scalar in five dimensions, and study their solutions. Section V is our conclusion.

II. DUALITY TRANSFORM TO SCALAR-TENSOR GRAVITY

Let us replace $R + \alpha R^2$ by $(1 + 2\alpha B)R - \alpha B^2$ in the gravitational action (1), by using a new scalar field B in five dimensions, as

$$S_{\text{gr}} = \frac{1}{2\kappa^2} \int d^5x \sqrt{-g} [(1 + 2\alpha B)R - \alpha B^2 - \Lambda]. \quad (5)$$

The equation of motion of the field B is algebraic and reads $B = R$. Hence, after substituting it back to the action (5) we get the equivalent action (1).

The action (5) can be brought to the Einstein frame by a Weyl transformation in *five* dimensions (cf., the similar transformation familiar in four dimensions, as was introduced in [6–8]),

$$g_{AB} = \Omega^{-2} \tilde{g}_{AB}, \quad \sqrt{-g} = \Omega^{-5} \sqrt{-\tilde{g}}, \quad (6)$$

with a suitably chosen factor Ω . By using the induced relation

$$R = \Omega^2 [\tilde{R} + 8\tilde{\square}f - 12\tilde{g}^{AB} f_{,A} f_{,B}], \quad (7)$$

where the tildes refer to the transformed quantities, and the definitions

$$f = \ln \Omega \quad \text{and} \quad f_{,A} = \frac{\partial_A \Omega}{\Omega}, \quad (8)$$

we find the gravitational action as follows:

$$S_{\text{gr}} = \frac{1}{2\kappa^2} \int d^5x \sqrt{-\tilde{g}} \Omega^{-5} [(1 + 2\alpha\phi)\Omega^2 (\tilde{R} + 8\tilde{\square}f - 12\tilde{g}^{AB} f_{,A} f_{,B}) - \alpha B^2 - \Lambda]. \quad (9)$$

Hence, to get the transformed gravitational action in the Einstein frame, we set

$$\Omega^3 = e^{3f} = 1 + 2\alpha B. \quad (10)$$

It yields

$$f = \frac{1}{3} \ln(1 + 2\alpha B) \quad (11)$$

and

$$S_{\text{gr}} = \frac{1}{2\kappa^2} \int d^5x \sqrt{-\tilde{g}} [\tilde{R} - 12\tilde{g}^{AB} \partial_A f \partial_B f - e^{-5f} (\alpha B^2 + \Lambda)]. \quad (12)$$

A canonically normalized scalar kinetic term is obtained after rescaling

$$\phi = 2\sqrt{3}f/\kappa, \quad (13)$$

so that we find

$$B = \frac{\exp(\frac{\sqrt{3}\kappa\phi}{2}) - 1}{2\alpha} \quad (14)$$

and the scalar potential

$$V(\phi) = \frac{1}{8\alpha\kappa^2} \left[\exp\left(\frac{\kappa\phi}{2\sqrt{3}}\right) - 2 \exp\left(-\frac{\kappa\phi}{\sqrt{3}}\right) + (1 + 4\alpha\Lambda) \exp\left(-\frac{5\kappa\phi}{2\sqrt{3}}\right) \right]. \quad (15)$$

As a result, the gravitational action takes the form

$$S_{\text{gr}} = \frac{1}{2\kappa^2} \int d^5x \sqrt{-\tilde{g}} \tilde{R} + \int d^5x \sqrt{-\tilde{g}} \left[-\frac{1}{2} \tilde{g}^{AB} \partial_A \phi \partial_B \phi - V(\phi) \right]. \quad (16)$$

Our main result of this section is the scalar potential (15) induced by the modified $(R + R^2)$ gravity in five space-time dimensions.

III. SCALAR DYNAMICS IN RSI BRANEWORLD

As is demonstrated in the previous section, the net effect of adding the R^2 term to the gravitational action amounts to the presence of the extra dynamical scalar field ϕ minimally coupled to gravity and having the scalar potential (15) in five dimensions.

Note that the five-dimensional cosmological constant Λ enters both (15) and (16) via the factor

$$\beta = 1 + 4\alpha\Lambda, \quad (17)$$

while the scalar potential is bounded from below provided that

$$\beta > 0. \quad (18)$$

In addition, demanding the five-dimensional cosmological constant to be negative, as is needed in the RSI model, we get

$$\beta \leq 1. \quad (19)$$

The profile of the scalar potential at $\beta = 1/2$ is given in Fig. 1.

The minimum of the scalar potential (15) is achieved at

$$\tilde{\phi}_0 = \frac{\kappa}{2\sqrt{3}} \phi_0 = \frac{1}{3} \ln(-2 + \sqrt{4 + 5\beta}) \quad (20)$$

so that, in particular, $\tilde{\phi}_0 = \phi_0 = 0$ for $\beta = 1$.

The value of the scalar potential at its minimum for any $0 < \beta \leq 1$ is given by

$$8\alpha\kappa^2 V(\phi_0) = \tilde{V}(\tilde{\phi}_0) = \frac{-6(\sqrt{4 + 5\beta} - 2 - \beta)}{(-2 + \sqrt{4 + 5\beta})^{5/3}}. \quad (21)$$

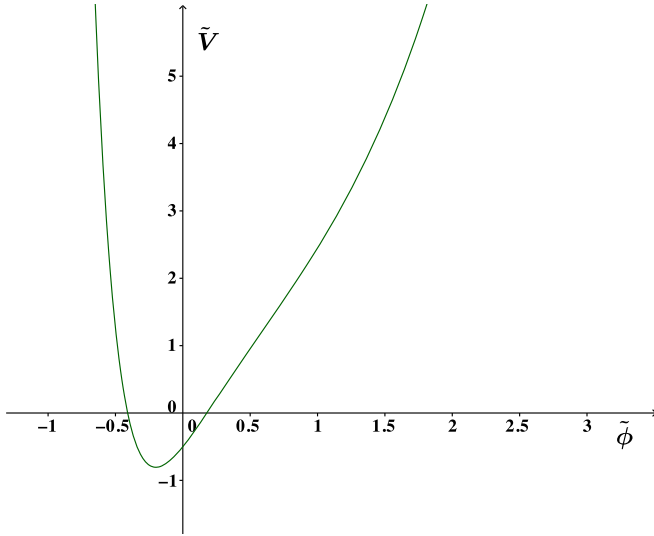


FIG. 1. The profile of the scalar potential $8\alpha\kappa^2 V(\phi) = \tilde{V}(\tilde{\phi})$ for $\tilde{\phi} = \frac{\kappa}{2\sqrt{3}}\phi$ and $\beta = 0.5$. The value of the scalar potential at its minimum is given by $\tilde{V}(\tilde{\phi}_0) \approx -0.81$. There are two solutions to $V = 0$: one at $\tilde{\phi}_{\text{left}} \approx -0.4$ and another at $\tilde{\phi}_{\text{right}} \approx 0.17$, on the left and on the right of the anti-de Sitter (AdS) minimum, respectively.

The $\tilde{V}_0(\tilde{\phi}_0)$ is an increasing negative function of β for $0 < \beta < 1$, and it vanishes at $\beta = 1$.

Since $\beta = 1$ implies $\Lambda = 0$, we assume that $0 < \beta < 1$ in what follows. Then the value of the scalar potential at its minimum, defined by (21), is always negative, which corresponds to an AdS vacuum.

IV. MODIFIED RSI MODEL

The Einstein equations for the action (16) in five dimensions read

$$\tilde{G}_{AB} = \tilde{R}_{AB} - \frac{1}{2}\tilde{g}_{AB}\tilde{R} = \kappa^2 T_{AB}, \quad (22)$$

where we have introduced the total energy-momentum tensor

$$T_{AB} = T_{AB}^{\phi} + T_{AB}^{\text{vis}} + T_{AB}^{\text{hid}} \quad (23)$$

as a sum of three contributions,

$$T_{AB}^{\phi} = \partial_A \phi \partial_B \phi + \tilde{g}_{AB} \left(-\frac{1}{2} \tilde{g}^{MN} \partial_M \phi \partial_N \phi - V(\phi) \right), \quad (24)$$

$$T_{AB}^{\text{hid}} = e^{-\frac{5\kappa}{2\sqrt{3}}\tilde{\phi}} \tilde{g}_{AB} \tau_{RS} \delta(y-L), \quad (25)$$

$$T_{AB}^{\text{vis}} = -e^{-\frac{5\kappa}{2\sqrt{3}}\tilde{\phi}} \tilde{g}_{AB} \tau_{RS} \delta(y). \quad (26)$$

The Randall-Sundrum (RS) ansatz for the five-dimensional space-time metric with the four-dimensional Poincaré symmetry is given by [3]

$$ds^2 = e^{-2A} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2, \quad (27)$$

where the warp factor $A(y)$ only depends upon the coordinate y of the hidden (fifth) dimension,¹ and $\eta_{\mu\nu}$ is the Minkowski metric in four other dimensions.

The corresponding (nonvanishing) Christoffel symbols, Ricci tensor, and Ricci scalar are

$$\Gamma_{\mu\nu}^5 = A' e^{-2A} \eta_{\mu\nu}, \quad (28)$$

$$R_{\mu\nu} = (A'' - 4A'^2) e^{-2A} \eta_{\mu\nu}, \quad (29)$$

$$R_{55} = 4A'' - 4A'^2, \quad (30)$$

$$R = 8A'' - 20A'^2, \quad (31)$$

respectively, where the primes denote the derivatives with respect to y .

Therefore, the Einstein equations take the form

$$G_{55} = 6(A')^2 = \kappa^2 T_{55} = \kappa^2 \left[\frac{1}{2} (\phi')^2 - V(\phi) \right], \quad (32)$$

$$G_{\mu\nu} = -3(A'' + 2A'^2) e^{-2A} \eta_{\mu\nu} = \kappa^2 T_{\mu\nu}, \quad (33)$$

$$T_{\mu\nu} = \tilde{g}_{\mu\nu} \left[-\frac{1}{2} \tilde{g}^{55} (\phi')^2 - V(\phi) \right] \quad (34)$$

$$+ e^{-\frac{5\kappa}{2\sqrt{3}}\tilde{\phi}} e^{-2A} \eta_{\mu\nu} \tau_{RS} \delta(y-L) \quad (35)$$

$$- e^{-\frac{5\kappa}{2\sqrt{3}}\tilde{\phi}} e^{-2A} \eta_{\mu\nu} \tau_{RS} \delta(y), \quad (36)$$

where we have assumed that the scalar field ϕ only depends upon y , because it is also required by the four-dimensional Poincaré invariance (on the visible brane).

The equation of motion of the scalar field $\phi(y)$ coupled to gravity reads

$$\phi'' + 4A'\phi' + \frac{V(\phi)}{d\phi} = 0. \quad (37)$$

As is already seen in Fig. 1, the scalar potential is very (exponentially) steep, whereas the right-hand side of Eq. (32) must be positive, so that a slow roll is impossible. Therefore, the scalar field is quickly *stabilized* at its minimum (AdS vacuum) ϕ_0 in five dimensions.

We numerically verified the scalar stabilization with the potential (15) by using the Runge-Kutta method and the Maxima software in application to a system of two coupled ordinary differential equations (32) and (37). When choosing the parameters as $\beta = 1/2$ (as in Fig. 1) and $\alpha = 1/96$ [it fixes the scale of V in Eq. (37)], the initial conditions $\phi' = V = A' = 0$ are consistent with (32). The corresponding numerical solutions to the functions $\tilde{\phi}(y)$ and $A'(y)$ are given in Fig. 2. We conclude that the functions $\tilde{\phi}(y)$ and

¹We denote by x^μ the remaining (four) space-time coordinates.

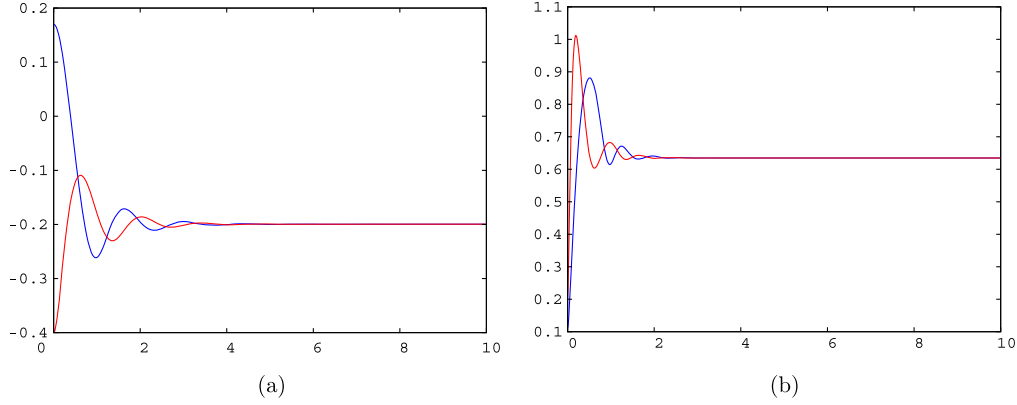


FIG. 2. (a) the functions $\tilde{\phi}(y)$, and (b) the functions $A'(y)$, with the left (red) and right (blue) initial conditions, respectively.

$A'(y)$ quickly approach constant values indeed, independently of their initial conditions.

Once the scalar field is at its minimum (AdS vacuum) ϕ_0 in five dimensions, the standard RSI scenario applies, being described by the solution [3]

$$A(y) = k|y|, \quad -L \leq y \leq L, \quad k > 0. \quad (38)$$

In particular, the size L of the fifth dimension is related to the Planck mass M_{Pl} in four dimensions as

$$\kappa^2 M_{\text{Pl}}^2 = \frac{1 - e^{-2kL}}{2k}. \quad (39)$$

All that is consistent with the other Einstein equations provided that

$$k^2 = -\kappa^2 V(\phi_0) \equiv -\kappa^2 V_{\text{min}} \quad \text{and} \quad \tau_{\text{RS}} = 6k/\kappa^2. \quad (40)$$

Then the negative cosmological constant in five dimensions is compensated in the visible brane by its tension, so that we get Minkowski geometry there.

It is worth mentioning that the effective value of τ_{RS} gets modified in our approach against the RSI one by the factor of $e^{-\frac{5\kappa}{2\sqrt{3}}\phi_0}$, because of Eqs. (25) and (26).

As in Ref. [3], the hierarchy between the electroweak scale and the Planck scale M_{Pl} in four dimensions (on the visible brane) is achieved via the presence of the exponential factor e^{-kL} induced by the extra dimension in the vacuum expectation value of the Higgs field, so that we must require $kL \approx \ln 10^{16} \approx 35$. Then the exponential term in Eq. (39) becomes very small and can be ignored. The Newtonian limit of RSI model leads to the similar exponentially small corrections to the Newton law of gravity, which are not in conflict with observations.

V. CONCLUSION

Given the phenomenological viability of the Randall-Sundrum braneworld, as the established and reasonable alternative to Kaluza-Klein compactification, it makes sense to analyze stability of the RS braneworld against possible modifications of gravity, as well as against quantum gravity corrections. In our investigation, we took a small step in this direction by proving stability of the RSI model against the simplest modification of the higher-dimensional gravity described by adding the scalar curvature squared term in five dimensions.

An impact of the R^2 -modified gravity on the RSI model can be simply described in terms of a single dynamical scalar field with the particular scalar potential (15). It is clear from our construction that this scalar has the gravitational origin as the spin-0 part of the five-dimensional space-time metric. We found that the value of the RSI parameter k is determined by dynamics of that scalar in the fifth dimension.

Our results and methods may be extendable to more general modified gravity models of braneworld as, e.g., those studied in Refs. [9,10].

ACKNOWLEDGMENTS

S. V. K. was supported by a Grant-in-Aid of the Japanese Society for Promotion of Science (JSPS) under Grant No. 26400252, a TMU President Grant of the Tokyo Metropolitan University in Japan, the World Premier International Research Center Initiative (WPI Initiative), MEXT, Japan, the CUiverse research promotion project by Chulalongkorn University (Grant No. CUAASC) in Bangkok, Thailand, and the Competitiveness Enhancement Program of Tomsk Polytechnic University in Russia.

- [1] V. A. Rubakov and M. E. Shaposhnikov, *Phys. Lett.* **125B**, 136 (1983).
- [2] N. Arkani-Hamed, S. Dimopoulos, and G. Dvali, *Phys. Lett. B* **429**, 263 (1998); **436**, 257 (1998).
- [3] L. Randall and R. Sundrum, *Phys. Rev. Lett.* **83**, 3370 (1999); **83**, 4690 (1999).
- [4] P. Hořava and E. Witten, *Nucl. Phys.* **B475**, 94 (1996).
- [5] R. Maartens and K. Koyama, *Living Rev. Relativ.* **13**, 5 (2010).
- [6] G. Bicknell, *J. Phys. A* **7**, 1061 (1974).
- [7] B. Whitt, *Phys. Lett.* **145B**, 176 (1984).
- [8] K.-I. Maeda, *Phys. Rev. D* **39**, 3159 (1989).
- [9] O. Bertolami and C. Carvalho, *Phys. Rev. D* **76**, 104048 (2007).
- [10] O. Bertolami, C. Carvalho, and J.N. Laia, *Nucl. Phys.* **B807**, 56 (2009).