

Long range azimuthal correlations of multiple gluons in gluon saturation limit

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We calculate the inclusive gluon correlation function for arbitrary number of gluons with full rapidity and transverse momentum dependence for the initial glasma state of the p-p, p-A, and A-A collisions. The formula we derive via superdiagrams generates cumulants for any number of gluons. Higher order cumulants contain information on correlations between multiple gluons, and they are necessary for calculations of higher dimensional ridges as well as flow coefficients from multiparticle correlations.

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I. INTRODUCTION

Hadrons produced in high multiplicity events of A-A, p-A, and p-p collisions show collective behavior [1,2]. This collectivity becomes apparent in the dihadron correlation function, which quantify the collective behavior of the hadrons in terms of the rapidity difference $\Delta\eta = \eta_1 - \eta_2$ and azimuthal angle difference $\Delta\phi = \phi_1 - \phi_2$. The origin of some of the contributions to the dihadron correlations are known. These are jet fragmentation around $\Delta\eta \sim \Delta\phi \sim 0$, resonance decays, and momentum conservation around $\Delta\phi \sim \pi$. When these are subtracted from the dihadron correlation function, the “double ridge” structure, $\cos(2\Delta\phi)$, becomes apparent. This shows that the correlation between hadron pairs becomes maximum when the hadrons are azimuthally collimated in the same direction or when they are back-to-back. Furthermore, these correlations in the near and away side are elongated in rapidity difference; the collimation and anticollimation effects are robust even though the hadrons are separated in rapidity for multiple units [3–18].

That the correlations are of long range in $\Delta\eta$ is attributed to the boost invariance; the gluons are produced at different rapidities, which is properly understood in the Regge limit of the QCD parton evolution. This is in contrast to the Bjorken limit of QCD, where gluons are emitted while being local in rapidity ($\Delta\eta \sim 0$), but they evolve in k_T during branching. As for the “double ridge” structure of the azimuthal correlations, there are currently two major attempts to explain the collectivity either as a final state (hydrodynamics) or as an initial state effect (glasma state by gluon saturation). One of them is a possible hydrodynamical evolution of the hadrons where the hadrons are affected by radial flow, and they come out around the relative angles $\Delta\phi \sim 0$ or $\Delta\phi \sim \pi$ [19–24]. The other approach searches for the origin of the collectivity of hadrons in the very early stages of collisions. Saturation of the gluons is expected to create a semihard mean transverse momentum in the target and projectile, which causes the emitted gluons to be azimuthally correlated. This is studied

by convolving unintegrated gluon distribution functions (UGD) from both the target and projectile, and this gives rise to ridge-type azimuthal correlations in the inclusive double gluon distribution. The initial correlation of the gluons are preserved when fragmentation functions are used to obtain final state hadrons from the correlated gluons [25–29].

The measured dihadron correlations alone are not enough to settle the dispute regarding the origin of the collectivity in high multiplicity hadron or nucleus collisions. Also, calculations based on the gluon saturation suggest that multigluon correlations exhibit strong non-gaussianity [30]. Therefore, examination of higher order inclusive distribution functions C_n becomes necessary to learn better about the true nature of the hadronic collectivity. The C_n 's for hadrons are measurable, and in an earlier study, we predicted that they would reveal higher-dimensional ridges [31]. C_n 's are also related to the observable flow moments $v_m(nPC)$, where nPC suggests that v_m is measured from n -particle correlations [7,32,33]. On the experimental side, measurements of multiple hadron correlations (trihadron, quadrohadron, etc.) in high multiplicity p-p, p-A, and A-A collisions are needed.

To study the hadronic correlations at a greater resolution, we derived triple and quadruple gluon correlations at arbitrary transverse momentum and rapidity dependence in [31] in the Gaussian white noise approximation [34–37]. The purpose of this paper is to generalize these calculations to arbitrary number of gluons and provide a formula that generates inclusive n -gluon distribution with full transverse momentum and rapidity dependency. Knowing all cumulants of a distribution is tantamount to knowing the distribution of the correlated random gluon production, and this distribution contains a complete information on the system. Hence, this work provides the solution to the problem of the gluon production from hadrons or nuclei at the saturation scale.

The outline of the paper is as follows. We first introduce the technology of superdiagrams that are used to obtain explicit expression for the inclusive gluon correlations with

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full momentum and rapidity dependence for any number of gluons. Then we give examples of how superdiagrams work for triple- and quadruple-gluon correlations, which were already calculated in an earlier work via regular glasma diagrams. We also derive, for the first time, the quintuple-gluon cumulant, C_5 , via the superdiagram technique. Finally, we list the superdiagrammatic rules and provide a formula for C_n .

II. SUPERDIAGRAMS FOR TRIPLE- AND QUADRUPLE-GLUON GLASMA DIAGRAMS

The inclusive gluon distribution functions C_n 's are calculated via connected diagrams. Therefore C_n 's are cumulants, not moments, and they contain information of the genuine multiparticle gluon correlations as cumulants do not contain disconnected diagrams. The double-, triple-, and quadruple-gluon cumulants are found by calculating 4, 16, and 96 connected glasma diagrams, respectively [31,38]. Using glasma diagrams to calculate even higher cumulants is impractical. Already for C_5 , the number of connected rainbow glasma diagrams to be calculated becomes 448.

In this section, we introduce the machinery of superdiagrams; a handful of diagrams that does the job of hundreds of connected glasma diagrams. With superdiagrams, one needs to calculate only $2(n-2)$ diagrams for the n -gluon cumulant, C_n . For example, C_4 can be easily obtained by four superdiagrams instead of calculating 96 connected glasma diagrams. For C_5 , one needs only six superdiagrams rather than 448 connected glasma diagrams.

Below we show how superdiagrams work for C_3 and C_4 . The triple-gluon cumulant is given by [31]

$$C_3(\mathbf{p}, \mathbf{q}, \mathbf{l}) = \frac{\alpha_s^3 N_c^3 S_\perp}{\pi^{12} (N_c^2 - 1)^5 \mathbf{p}_\perp^2 \mathbf{q}_\perp^2 \mathbf{l}_\perp^2} \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} (\mathcal{T}_1 + \mathcal{T}_2), \quad (1)$$

where N_c is the gluon color factor, S_\perp is the transverse area of the overlap between the target and projectile, \mathbf{p}_\perp , \mathbf{q}_\perp , and \mathbf{l}_\perp are transverse momentum variables of the produced gluons. The strong coupling constant α_s can be taken as constant, or it can be running with respect to the gluon

transverse momenta. The terms \mathcal{T}_1 and \mathcal{T}_2 that contain the UGDs are given by

$$\mathcal{T}_1 = \Phi_{1,p}^2(\mathbf{k}_\perp) [2 \times \Phi_{1,q}(\mathbf{k}_\perp)] \Phi_{2,p}(\mathbf{p}_\perp - \mathbf{k}_\perp) \mathcal{T}_{A_2}, \quad (2)$$

$$\mathcal{T}_2 = \Phi_{2,l}^2(\mathbf{k}_\perp) [2 \times \Phi_{2,q}(\mathbf{k}_\perp)] \Phi_{1,p}(\mathbf{p}_\perp - \mathbf{k}_\perp) \mathcal{T}_{A_1}, \quad (3)$$

$$\begin{aligned} \mathcal{T}_{A_1, A_2} = & [\Phi_{1(2),q}(\mathbf{q}_\perp - \mathbf{k}_\perp) + \Phi_{1(2),q}(\mathbf{q}_\perp + \mathbf{k}_\perp)] \\ & \times [\Phi_{1(2),l}(\mathbf{l}_\perp - \mathbf{k}_\perp) + \Phi_{1(2),l}(\mathbf{l}_\perp + \mathbf{k}_\perp)]. \end{aligned} \quad (4)$$

Here the indices 1 and 2 of the UGDs (Φ) refer to nucleus 1 and nucleus 2, and the second letter indices refer to the rapidity variables. The argument of the UGDs in the parentheses correspond to the transverse momentum of the gluons, and \mathbf{k}_\perp is the integration variable. The transverse momentum dependence of the UGDs at any order ($n = 3, 4, \dots$) follows a simple pattern [compare Eqs. (2)–(4) with Eqs. (6)–(8)]. On the other hand, the rapidity indices of the UGDs are nontrivial. The main power of the glasma superdiagrams is to readily determine the rapidity dependence of the cumulant C_n at any order.

According to the conventions we use, p is the momentum and rapidity index of the gluon closest to nucleus 1 (A_1) in rapidity evolution, whereas l is the index of the gluon which is closest to nucleus 2 (A_2). From Eq. (2), one observes that the UGD with the rapidity index p is squared, whereas in Eq. (3), the UGD with the rapidity index l is squared since it is the closest one to A_2 . The rapidity structure of Eq. (2) can be summarized by the glasma superdiagrams in Fig. 1. These two superdiagrams represent 16 connected glasma diagrams that is necessary to calculate C_3 , and they give \mathcal{T}_1 and \mathcal{T}_2 .

The quadruple-gluon cumulant that requires calculation of 96 glasma diagrams can be produced with four superdiagrams. The fourth cumulant is given by [31]

$$C_4(\mathbf{p}, \mathbf{q}, \mathbf{l}, \mathbf{w}) = \frac{\alpha_s^4 N_c^4 S_\perp}{\pi^{16} (N_c^2 - 1)^7 \mathbf{p}_\perp^2 \mathbf{q}_\perp^2 \mathbf{l}_\perp^2 \mathbf{w}_\perp^2} \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} (\mathcal{Q}_1 + \mathcal{Q}_2), \quad (5)$$

where

$$\mathcal{Q}_1 = \Phi_{1,p}^2(\mathbf{k}_\perp) \Phi_{1,q}(\mathbf{k}_\perp) [4 \times \Phi_{1,l}(\mathbf{k}_\perp) + 2 \times \Phi_{1,q}(\mathbf{k}_\perp)] \Phi_{2,p}(\mathbf{p}_\perp - \mathbf{k}_\perp) \mathcal{Q}_{A_2}, \quad (6)$$

$$\mathcal{Q}_2 = \Phi_{2,w}^2(\mathbf{k}_\perp) \Phi_{2,l}(\mathbf{k}_\perp) [4 \times \Phi_{2,q}(\mathbf{k}_\perp) + 2 \times \Phi_{2,l}(\mathbf{k}_\perp)] \Phi_{1,p}(\mathbf{p}_\perp - \mathbf{k}_\perp) \mathcal{Q}_{A_1}, \quad (7)$$

$$\begin{aligned} \mathcal{Q}_{A_1(A_2)} = & [\Phi_{1(2),q}(\mathbf{q}_\perp - \mathbf{k}_\perp) + \Phi_{1(2),q}(\mathbf{q}_\perp + \mathbf{k}_\perp)] [\Phi_{1(2),l}(\mathbf{l}_\perp - \mathbf{k}_\perp) + \Phi_{1(2),l}(\mathbf{l}_\perp + \mathbf{k}_\perp)] \\ & \times [\Phi_{1(2),w}(\mathbf{w}_\perp - \mathbf{k}_\perp) + \Phi_{1(2),w}(\mathbf{w}_\perp + \mathbf{k}_\perp)]. \end{aligned} \quad (8)$$

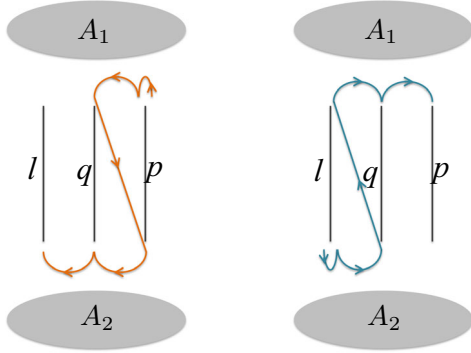


FIG. 1. Two glasma superdiagrams to determine the rapidity indices of UGDs in Eqs. (2) and (3). The superdiagram on the left produces the rapidity indices of the UGDs in \mathcal{T}_1 [see Eq. (2)], whereas the superdiagram on the right is for \mathcal{T}_2 [see Eq. (3)]. Whether one begins writing down UGDs starting from A_1 (left figure) or A_2 (right figure), the first UGD is always squared, which is $\Phi_{1,p}^2$ when one starts from A_1 and $\Phi_{2,l}^2$ when one starts from A_2 . C_3 contains multiplication of three UGDs from A_1 and another three from A_2 . In the superdiagram on the left, three hoppings on A_1 already brings three UGDs, so before reaching l , one has to move to the other nucleus. Another rule of the superdiagrams is that one starts from the rightmost index on A_1 (p) and ends at the leftmost index on A_2 as in the left superdiagram. There is also the same number of superdiagrams where one starts from the leftmost index on A_2 and end at the rightmost index on A_1 as in the right superdiagram. Since these two sets of superdiagrams are simultaneous horizontal and vertical reflections of each other, in the following figures, we will only draw the first set of the superdiagrams where one start from the rightmost index on A_1 . As for the hoppings, only next-to-nearing hoppings are allowed; hence, connecting p to l by skipping q is not allowed.

Figure 2 shows the two superdiagrams that contributes to \mathcal{Q}_1 ; the two mirror images of these contractions (not shown in the figure) gives \mathcal{Q}_2 .

III. RULES FOR SUPERDIAGRAMS

The rules of superdiagrams can be summarized as follows.

- (1) On nucleus A_1 (A_2), one starts from the rightmost (leftmost) rapidity index and end at the leftmost (rightmost) rapidity index on A_2 (A_1). The first

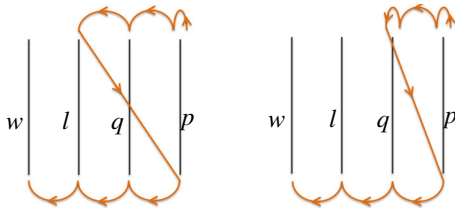


FIG. 2. Two glasma superdiagrams that contribute to \mathcal{Q}_1 given in Eq. (6), which is part of C_4 . The other two superdiagrams are not shown as they are simultaneous horizontal and vertical reflections of these two superdiagrams.

rapidity site is always visited twice; it is the rightmost (leftmost) site on A_1 (A_2).

- (2) For C_n , there must be n hoppings on each nucleus. For example, there are four hoppings on A_1 for C_4 as shown in Fig. 2. Also, the line connecting the rapidity sites should continue by visiting nearest neighbors without skipping any site. For example, going from p directly to l by skipping q is not allowed.

- (3) Each site can only be visited twice at maximum. In light of these rules, one can draw the possible six superdiagrams for C_5 . Figure 3 shows three of these superdiagrams.

Now a multigluon cumulant at any order for $n \geq 4$ can be constructed by the following recipe.

- (i) The prefactor and the integral part of C_n shall be in the form

$$C_n = \frac{\alpha_s^n N_c^n S_\perp}{\pi^{4n} (N_c^2 - 1)^{2n-1}} \left(\prod_{i=1}^n \frac{1}{p_{\perp i}^2} \right) \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} (\mathcal{N}_1 + \mathcal{N}_2). \quad (9)$$

- (ii) \mathcal{N}_1 is given by

$$\mathcal{N}_1 = \Phi_{1,p_1}^2(\mathbf{k}_\perp) \left[\prod_{j=1}^{n-3} \Phi_{1,p_{j+1}}^2(\mathbf{k}_\perp) \right] \times \left[\sum_{h=1}^{n-2} 2^h \Phi_{1,p_{h+1}}^2(\mathbf{k}_\perp) \right] \Phi_{2,p_n}(\mathbf{p}_{\perp 1} - \mathbf{k}_\perp) \mathcal{N}_{A_2}. \quad (10)$$

\mathcal{N}_2 can be obtained from \mathcal{N}_1 by making these changes: replace the nucleus index 1 with 2, and replace the momentum index p_i with p_{n+1-i} , where n is the order of the cumulant. For example, one should change the indices according to $p \leftrightarrow w$ and $q \leftrightarrow l$ for C_4 .

- (iii) $\mathcal{N}_{A_1(A_2)}$ that is contained in $\mathcal{N}_{2(1)}$ is given by

$$\mathcal{N}_{A_1(A_2)} = \prod_{m=2}^n [\Phi_{1(2),p_m}(\mathbf{p}_{\perp m} - \mathbf{k}_\perp) + \Phi_{1(2),p_m}(\mathbf{p}_{\perp m} + \mathbf{k}_\perp)], \quad (11)$$

where n is again the order of the cumulant C_n . For C_4 , $(p_2, p_3, p_4) = (q, l, w)$, and $(\mathbf{p}_{\perp 2}, \mathbf{p}_{\perp 3}, \mathbf{p}_{\perp 4}) = (\mathbf{q}_\perp, \mathbf{l}_\perp, \mathbf{w}_\perp)$ [see Eq. (8)].

Equation (9) together with the definitions given in Eqs. (10) and (11) is the main result of this paper. C_n gives inclusive n -gluon distribution, and it quantifies the correlation of n gluons in transverse momentum and rapidity. C_n can be used to calculate higher-dimensional ridges and flow moments $v_m(n\text{PC})$ from n -particle correlations. In principle, the cumulants C_n can be summed to

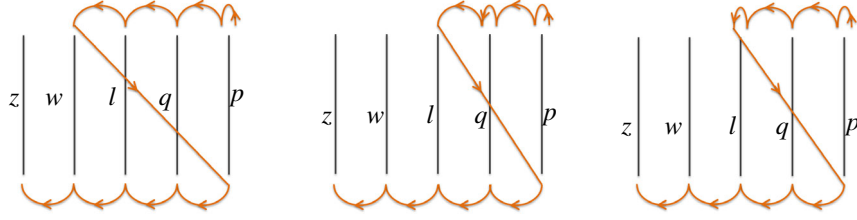


FIG. 3. Three glasma superdiagrams that contribute to C_5 . The other three superdiagrams are not shown as they are simultaneous horizontal and vertical reflections of these three superdiagrams.

find the cumulant generating function, and then the probability distribution via Laplace transform of this generating function. We will not make any attempt in this direction since the UGDs contained in C_n are complicated functions, and in practice, they are in the form of numerical tables.

Now we shall check if the recipe given above yields the correct number of glasma diagrams for C_5 . At the order $n = 5$, \mathcal{N}_{A_2} will contain 2^4 separate terms, so \mathcal{N}_1 will contain $(2^3 + 2^2 + 2) \times 2^4 = 224$ terms. \mathcal{N}_2 contains the same number of diagrams, so the total number of connected diagrams becomes 448.

We shall now check if the cumulant expansion gives the same number of diagrams. The moments μ_n ¹ can be written in terms of the cumulants κ_n

$$\mu_n = \sum_{k=1}^n B_{n,k}(\kappa_1, \dots, \kappa_{n-k+1}), \quad (12)$$

where $B_{n,k}$ are partial Bell polynomials. The fifth cumulant κ_5 is same as C_5 , and it can be written in terms of the lower-order cumulants and the fifth moment μ_5

$$\begin{aligned} \kappa_5 = \mu_5 - 5\kappa_4\kappa_1 - 10\kappa_3\kappa_2 - 10\kappa_3\kappa_1^2 \\ - 15\kappa_2^2\kappa_1 - 10\kappa_2\kappa_1^3 - \kappa_1^5. \end{aligned} \quad (13)$$

Here, μ_5 includes all connected and disconnected glasma rainbow diagrams with five gluons. Subtracting from μ_5 the other terms in the rhs of Eq. (13) gives the connected five gluon diagrams, which is κ_5 . A word of caution regarding the term κ_2^2 is in order [31]. κ_2 includes two upper and two lower rainbow diagrams. However, the term κ_2^2 mixes upper and lower rainbow diagrams. Such mixings do not occur in μ_5 , where the disconnected diagrams are formed either by the combination of lower or upper rainbow diagrams. So, since μ_5 is already free from mixed diagrams, subtracting any mixed diagrams from μ_5 would result in wrong counting. We resolve this issue by modifying the term in Eq. (13) as such

¹ μ_n are moments, not ‘‘central moments.’’ In the literature, sometimes μ'_n are used for moments.

$$15\kappa_2^2\kappa_1 \rightarrow \frac{1}{2}15\kappa_2^2\kappa_1, \quad (14)$$

so that only the diagrams of the form upper \otimes upper and lower \otimes lower are subtracted from μ_5 , but not those of the form lower \otimes upper or upper \otimes lower.

The fifth moment μ_5 includes all connected and disconnected rainbow glasma diagrams at the five gluon level, and the number of such diagrams is given by $2(2n - 1)!! - 1$ [31]. The first factor of 2 accounts for both upper and lower rainbow diagrams, the term with the double factorial counts the number of pairings between gluons, and 1 is subtracted at the end not to double count the maximally disconnected glasma diagram (κ_1^n), which is shown as concentric circles [31]. At the order $n = 5$, there are $\mu_5 = 1889$ connected and disconnected rainbow diagrams in total. The numbers of diagrams that each cumulant for $n < 5$ contains have been given in Ref. [31]: $\kappa_1 = 1$, $\kappa_2 = 4$, $\kappa_3 = 16$, and $\kappa_4 = 96$. From Eq. (13) with the modification in Eq. (14), we find the number of connected diagrams for C_5

$$\begin{aligned} \kappa_5 = 1889 - 5 \times 96 - 10 \times 16 \times 4 - 10 \times 16 \\ - 15 \times \frac{1}{2} \times 4^2 - 10 \times 4 - 1 \\ = 448. \end{aligned} \quad (15)$$

This number is the same as what our recipe previously gave; see the discussion above Eq. (13). This completes the proof that our formulas given in the recipe Eqs. (9)–(11) produce the correct number of diagrams.

IV. SUMMARY AND OUTLOOK

We have developed a superdiagrammatic technique which allows calculating the inclusive gluon distributions C_n at the saturation limit easily bypassing the necessity of calculating thousands of glasma diagrams. Inclusive gluon distribution functions contain information on azimuthal and rapidity correlations between produced gluons in p-p, p-A, and A-A collisions. Multiple-hadron correlations are measured in these experiments. These hadronic correlation functions are used to measure the ridge-type azimuthal correlations between hadrons as well as flow moments

$v_m(nPC)$ from multiple hadrons. On the theory side, hadron correlations are calculated by convolving the gluon correlation functions C_n with fragmentation functions. Higher dimensional ridges from number of gluons $n > 2$ and flow moments v_m from multiparticle correlations are a testing ground for different approaches such as hydrodynamics and gluon saturation/glasma physics.

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