## Determination of the $\eta$ - $\eta$ / mixing angle

Fu-Guang Cao

Institute of Fundamental Sciences, Massey University, Private Bag 11 222, Palmerston North, New Zealand (Received 2 December 2011; published 22 March 2012)

We extract  $\eta$ - $\eta'$  mixing angle and the ratios of decay constants of light pseudoscalar mesons  $\pi^0$ ,  $\eta$ , and  $\eta'$  using recently available *BABAR* measurements on  $\eta$ -photon and  $\eta'$ -photon transition form factors and more accurate experimental data for the masses and two-photon decay widths of the light pseduoscalar mesons.

DOI: 10.1103/PhysRevD.85.057501

PACS numbers: 14.40.Be, 13.40.Gp, 14.40.Df

Determining the composition of  $\eta$  and  $\eta'$  mesons has attracted continuous interest in hadronic physics. The idea of  $\eta$  and  $\eta'$  containing gluonic and intrinsic  $|c\bar{c}\rangle$  components has long been employed in explaining many experimental results, including recent observations of large branching ratios for some decay processes of  $J/\psi$  and *B* mesons into pseudoscalar mesons [1].

There are three charge neutral states in the nonet of pseudoscalar mesons in the SU(3)<sub>*F*</sub> quark model:  $\pi^0$ ,  $\eta_8$  and  $\eta_1$ . The latter two mix to give the physical particles  $\eta$  and  $\eta'$ ,

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \eta_8 \\ \eta_1 \end{pmatrix}.$$
(1)

Alternatively, one could use the quark-flavor basis mixing scheme [2],

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi \\ \sin\phi & \cos\phi \end{pmatrix} \begin{pmatrix} \eta_q \\ \eta_s \end{pmatrix},$$
(2)

with  $|\eta_q\rangle = \frac{1}{\sqrt{2}}(|u\bar{u}\rangle + |d\bar{d}\rangle)$  and  $|\eta_s\rangle = |s\bar{s}\rangle$ . The mixing angles in the two schemes are related via  $\theta = \phi$  –  $\arctan\sqrt{2} \simeq \phi - 54.7^{\circ}$ . A two-mixing-angle scheme has also been suggested in the study of the mixing of decay constants [3].

The mixing angle can neither be calculated from first principles in QCD nor measured directly—it has to be determined phenomenologically. There are a lot of studies on this subject using different methods and a number of different processes, including various decay processes involving the light pseudoscalar mesons [1–6].

One important source of information in determining the mixing angle is the transition processes,  $\gamma \gamma^* \rightarrow \eta$ ,  $\eta'$ , for which the transition from factors (TFFs),  $F_{\eta\gamma}(Q^2)$  and

 $F_{\eta'\gamma}(Q^2)$ , with  $Q^2$  being the virtuality of the off-mass-shell photon, are defined. The usual procedure [7] using the TFFs to evaluate the  $\eta$ - $\eta'$  mixing angle is to calculate the  $Q^2$  dependence of these transition form factors and compare with the experimental data, which are given at a certain range of  $Q^2$  [8–10]. However, theoretical calculations for the TFFs at finite  $Q^2$  suffer sizable corrections and are sensitive to the nonperturbative model used for the distribution amplitude of the mesons, which results in large uncertainties in determining the mixing angle.

Two analytical constraints on the  $\eta - \eta'$  mixing were obtained in Ref. [11] by considering the two-photon decays of  $\eta$  and  $\eta'$  and the asymptotic behavior of the  $\eta$  and  $\eta'$ TFFs in the limit  $Q^2 \rightarrow \infty$ , together with the fact that the asymptotic behavior of the meson TFFs is firmly predicted by QCD [12]. Newly available *BABAR* data [10] extend the measurements for the  $\eta$  and  $\eta'$  TFFs to higher  $Q^2$  and to a much larger range of  $Q^2$ , and thus provide new information for the  $\eta$  and  $\eta'$  TFFs at  $Q^2 \rightarrow \infty$ . At the same time, experimental information on the masses and two-photon decay widths of mesons involved have improved considerably over the last decade. These new experimental data shall have an impact on the determining of the  $\eta - \eta'$  mixing parameters.

In this paper, we extract the  $Q^2 \rightarrow \infty$  behavior of the  $\eta$ and  $\eta'$  TFFs from the *BABAR* data. Using this new information and updated experimental data about the twophoton decays  $\eta \rightarrow \gamma \gamma$  and  $\eta' \rightarrow \gamma \gamma$  [13], we determine the  $\eta$ - $\eta'$  mixing angle and the ratios of decay constants in the two mixing schemes [see Eqs. (1) and (2)] using the method of Ref. [11].

The analytical expressions obtained in Ref. [11] for the mixing angle  $\theta$  and the ratio of the decay constants of the  $\eta_1$  and  $\eta_8$  states  $r = f_1/f_8$  are

$$\tan\theta = \frac{-(1+c^2)(\rho_1+\rho_2) + \sqrt{(1+c^2)^2(\rho_1+\rho_2)^2 + 4(c^2-\rho_1\rho_2)(1-c^2\rho_1\rho_2)}}{2(c^2-\rho_1\rho_2)},$$
(3)

$$r = \frac{(1+c^2)(\rho_1 - \rho_2) + \sqrt{(1+c^2)^2(\rho_1 - \rho_2)^2 + 4c^2(1+\rho_1\rho_2)^2}}{2c(1+\rho_1\rho_2)},$$
(4)

where  $c = \sqrt{8}$  and

$$\rho_1 = \left[\frac{\Gamma_{\eta \to \gamma\gamma}}{\Gamma_{\eta' \to \gamma\gamma}} \frac{m_{\eta'}^3}{m_{\eta}^3}\right]^{1/2},\tag{5}$$

$$\rho_2 = \frac{F_{\eta\gamma}(Q^2 \to \infty)}{F_{\eta'\gamma}(Q^2 \to \infty)}.$$
(6)

One advantage of determining the mixing parameters from Eqs. (3)–(6) is that both the theoretical uncertainty incurred in calculating the TFFs at finite  $Q^2$  and the experimental uncertainty are minimized by considering the ratios of the decay widths for the two-photon decay processes and the ratios of the transition form factors at large  $Q^2$ .

Furthermore, considering the ratio of the decay widths for the  $\pi^0 \rightarrow \gamma \gamma$  and  $\eta \rightarrow \gamma \gamma$  processes, we can also determine the ratios  $f_8/f_{\pi}$  and  $f_1/f_{\pi}$ ,

$$\frac{f_8}{f_{\pi}} = \rho_0 \left[ \frac{c_8}{c_{\pi}} \cos\theta - \frac{1}{r} \frac{c_1}{c_{\pi}} \sin\theta \right],\tag{7}$$

$$\frac{f_1}{f_{\pi}} = \rho_0 \bigg[ \frac{c_8}{c_{\pi}} r \cos\theta - \frac{c_1}{c_{\pi}} \sin\theta \bigg], \tag{8}$$

where  $c_{\pi} = 1$ ,  $c_8 = 1/\sqrt{3}$ ,  $c_1 = 2\sqrt{2}/\sqrt{3}$ , and

$$\rho_0 = \left[\frac{\Gamma_{\pi^0 \to \gamma\gamma}}{\Gamma_{\eta \to \gamma\gamma}} \frac{m_\eta^3}{m_{\pi^0}^3}\right]^{1/2}.$$
(9)

The above analysis can be easily applied to the quark-flavor basis mixing scheme [see Eq. (2)] by replacing the parameters  $c = c_1/c_8$ ,  $r = f_1/f_8$ ,  $c_8$ , and  $c_1$  with  $c' = c_s/c_q = \sqrt{2}/5$ ,  $r' = f_s/f_q$ ,  $c_q = 5/3$ , and  $c_s = \sqrt{2}/3$ , respectively [11].

The parameters  $\rho_0$  and  $\rho_1$  can be fixed by using the masses and two-photon decay widths of  $\pi^0$ ,  $\eta$  and  $\eta'$ . We employ the data given by the 2010 Particle Data Group [13],

$$\Gamma_{\pi^0 \to \gamma\gamma} = 7.74 \pm 0.46 \text{ eV},$$
  

$$\Gamma_{\eta \to \gamma\gamma} = 0.510 \pm 0.026 \text{ keV},$$
(10)  

$$\Gamma_{\eta' \to \gamma\gamma} = 4.28 \pm 0.19 \text{ keV},$$

$$m_{\pi^0} = 134.9766 \pm 0.0006 \text{ MeV},$$
  
 $m_{\eta} = 547.853 \pm 0.024 \text{ MeV},$  (11)  
 $m_{\eta'} = 957.78 \pm 0.06 \text{ MeV}.$ 

We will use the CLEO [9] and *BABAR* [10] data for the TFFs at large  $Q^2$  to determine the parameter  $\rho_2$ . The CLEO Collaboration [9] measured  $F_{\eta\gamma}(Q^2)$  and  $F_{\eta'\gamma}(Q^2)$ in the  $Q^2$  regions up to 20 and 30 GeV<sup>2</sup>, respectively, and presented the data in a monopole form proposed in Ref. [14]:

$$|F_{P\gamma}(Q^2)|^2 = \frac{1}{(4\pi\alpha)^2} \frac{64\pi\Gamma_{P\to\gamma\gamma}}{m_P^3} \frac{1}{(1+Q^2/\Lambda_P^2)^2},$$
 (12)

where  $\alpha \simeq 1/137$  is the QED fine coupling constant and  $\Lambda_P$  is the pole mass parameter.

The *BABAR* Collaboration [10] recently measured the  $\eta$ -photon and  $\eta'$ -photon transition form factors in the  $Q^2$  range from 4 to 40 GeV<sup>2</sup>. The results were not presented in the monopole form [Eq. (12)], partially because their results for the pion-photon transition from factor exhibit a very quick growth for  $Q^2 > 15$  GeV<sup>2</sup> [15], which is very hard to explain in QCD [16]. However, this trend of fast growth is noticeably missing from the *BABAR* data for the  $\eta$ -photon and  $\eta'$ -photon transition form factors, and thus the *BABAR* data for the  $\eta$ -photon and  $\eta'$ -photon transition form factors are consistent with perturbative QCD calculations for the form factors and shall be described with the monopole form as given by Eq. (12).

We use QCD-motivated monopole form Eq. (12) to fit experimental data. The values of  $\Lambda_{\eta}$  and  $\Lambda'_{\eta}$  in Eq. (12) that were determined using the CLEO data, *BABAR* data, and the combined data are presented in Table I. We have combined the statistical and systematic errors for the CLEO data in quadrature, since the *BABAR* data are presented with only combined errors. The values of  $\chi^2/d.o.f$  given in the table provide further justification for the use of Eq. (12) in describing these data. The values of  $\Lambda_{\eta}$  and  $\Lambda'_{\eta}$  determined with the CLEO and *BABAR* data agree within their uncertainties, but the *BABAR* data greatly improve the accuracy in determining the values of  $\Lambda_{\eta}$  and  $\Lambda'_{\eta}$ . Using the combined data in the fitting changes the results slightly.

The parameter  $\Lambda_P$  has a natural explanation as the pole mass of vector meson in the vector meson–dominated

	η		$\eta'$	
	$\Lambda_{\eta}$ (MeV)	$\chi^2/d.o.f$	$\Lambda_{\eta'}$ (MeV)	$\chi^2/d.o.f$
CLEO	$775 \pm 12$	0.95	856 ± 10	0.88
BABAR	$787 \pm 7$	0.99	$861 \pm 4$	1.04
CLEO + BABAR	$784 \pm 6$	0.96	849 ± 6	0.88

TABLE I.  $\Lambda_P$  and  $\chi^2/d.o.f$  in fitting the data for the TFFs with Eq. (12).

	The many parameters determined for the 7/8 7/1 many seneme.			
	θ	$f_{1}/f_{8}$	$f_8/f_{\pi}$	$f_1/f_{\pi}$
CLEO	$-16.26 \pm 0.86$	$1.162 \pm 0.053$	$0.955 \pm 0.042$	$1.109 \pm 0.053$
BABAR	$-16.54 \pm 0.71$	$1.146\pm0.045$	$0.966 \pm 0.041$	$1.107\pm0.050$
CLEO + BABAR	$-16.84 \pm 0.72$	$1.128\pm0.044$	$0.979\pm0.042$	$1.105 \pm 0.050$

TABLE II. The mixing parameters determined for the  $\eta_8$ - $\eta_1$  mixing scheme.

TABLE III.	The mixing parameters	determined for the quark-flavor	basis mixing scheme.
------------	-----------------------	---------------------------------	----------------------

	$\phi$	$f_s/f_q$	$f_q/f_{\pi}$	$f_s/f_{\pi}$
CLEO	$38.11\pm0.79$	$1.197\pm0.065$	$1.076\pm0.044$	$1.29\pm0.10$
BABAR	$37.90 \pm 0.70$	$1.177 \pm 0.054$	$1.077 \pm 0.044$	$1.268\pm0.088$
CLEO + BABAR	$37.66\pm0.70$	$1.156\pm0.054$	$1.078\pm0.044$	$1.246 \pm 0.087$

model for the TFFs. The values we obtained,  $\Lambda_{\eta} \sim$  780 MeV and  $\Lambda_{\eta} \sim$  860 MeV, are very close to the masses of  $\rho$ (770 MeV) and  $K^*$ (890 MeV).

The results for the mixing angle and decay constants determined using the CLEO data, BABAR data, and the combined data, together with the two-photon decay widths, are presented in Tables II and III for the  $\eta_8$ - $\eta_1$  mixing scheme and quark-flavor basis mixing scheme, respectively. The mixing angle obtained in this work,  $\phi \simeq 37^{\circ} \sim$ 38°, is slightly smaller than the central value of 39.8° obtained in Ref. [11]. This is mainly due to an increase in the estimation for the  $\Gamma_{\eta \to \gamma \gamma}$  by the 2010 Particle Data Group. This increase also affects slightly the results for the ratios of decay constants. The uncertainties for the mixing angles and the ratios  $f_1/f_8$  and  $f_s/f_q$  obtained in this work are considerably smaller than that given in Ref. [11] due to the new, more accurate experimental data for the meson masses, the two-photon decay widths, and the mesonphoton transition form factors. The uncertainties for the other ratios of decay constants,  $f_8/f_{\pi}$  and  $f_1/f_{\pi}$  in the  $\eta_8$ - $\eta_1$  mixing scheme and  $f_q/f_{\pi}$  and  $f_s/f_{\pi}$  in the quarkflavor basis mixing scheme, are also generally smaller than that estimated in Ref. [11].

Our results for the mixing angle are in agreement with recent results of  $\phi \simeq 37^{\circ} \sim 42^{\circ}$  obtained with other methods [6]. The value of  $f_8/f_{\pi}$  obtained in this work is smaller

than that obtained with Chiral Perturbation Theory at nextto-leading order [4] and some phenomenological analyses [6], but is larger than the result reported in Ref. [5]. We note that the Chiral Perturbation Theory result may be alerted by higher-order corrections. As it has been pointed out in Ref. [11], in the previous studies either the questionable assumption that the decay constants and the particle states share the same mixing scheme or the twomixing-angle scheme was adopted. The relations between the mixing parameters involved in the two-mixing-angle scheme and those that appear in our model remain to be studied further.

In summary, understanding the composition of the light pseudoscalar mesons  $\eta$  and  $\eta'$  is of great importance in the study of many hadron processes involved these mesons. Employing the two analytical constraints on the  $\eta$ - $\eta'$ mixing proposed by us in a previous work, we extracted the  $\eta$ - $\eta'$  mixing angle and the ratios of decay constants in two widely used mixing schemes using recently available *BABAR* measurements on the  $\eta$ -photon and  $\eta'$ -photon transition form factors and more accurate experimental data for the masses and two-photon decay widths of the light pseduoscalar mesons.

F.G. thanks X.-H. Guo for his hospitality at Beijing Normal University where part of this work was done.

- C. Di Donato, G. Ricciardi, and I. I. Bigi, Phys. Rev. D 85, 013016 (2012); Y.-D. Tsai, H.-N. Li, and Q. Zhao, Phys. Rev. D 85, 034002 (2012); R. Fleischer, R. Knegjens, and G. Ricciardi, Eur. Phys. J. C 71, 1798 (2011).
- [2] T. Feldmann, P. Kroll, and B. Stech, Phys. Rev. D 58, 114006 (1998); Phys. Lett. B 449, 339 (1999).
- [3] See, e.g., J. Schechter, A. Subbaraman, and H. Weigel, Phys. Rev. D 48, 339 (1993); T. Feldmann and P. Kroll, Eur. Phys. J. C 5, No. 2, 327 (1998).
- [4] J. F. Donoghue, B. R. Holstein, and Y. C. R. Lin, Phys. Rev. Lett. 55, 2766 (1985); R. Kaiser and H. Leutwyler, in *Proceedings of Workshop on Nonperturbative Methods* in *Quantum Field Theory* (University of Adelaide, Adelaide, Australia, 1998), p. 15.
- [5] A. V. Kisselev and V. A. Petrov, Z. Phys. C 58, 595 (1993).
- [6] For recent studies, see, e.g., C. E. Thomas, J. High Energy Phys. 10 (2007) 026; D. M. Asner *et al.* (CLEO Collaboration), Phys. Rev. D **79**, 072007 (2009);

F. Ambrosino *et al.* (KLOE Collaboration), J. High Energy Phys. 07 (2009) 105; S. D. Bass, Acta Phys. Pol. B Proc. Suppl. **2**, 11 (2009); R. Escribano, Eur. Phys. J. C **65**, 467 (2009); V. Mathieu and V. Vento, Phys. Rev. D **81**, 034004 (2010); S. Kiesewetter and V. Vento, Phys. Rev. D **82**, 034003 (2010); H.-W. Ke, X.-H. Yuan, and X.-Q. Li, Int. J. Mod. Phys. A **26**, 4731 (2011).

- [7] See, e.g., J. Cao, F.-G. Cao, T. Huang, and B.-Q. Ma, Phys. Rev. D 58, 113006 (1998).
- [8] H.-J. Behrend *et al.* (CELLO Collaboration), Z. Phys. C 49, 401 (1991).
- [9] J. Gronberg *et al.* (CLEO Collaboration), Phys. Rev. D 57, 33 (1998).
- [10] P. A. Sanchez *et al.* (*BABAR* Collaboration), Phys. Rev. D 84, 052001 (2011).
- [11] F.G. Cao and A.I. Signal, Phys. Rev. D 60, 114012 (1999).

- [12] G.P. Lepage and S.J. Brodsky, Phys. Lett. 87B, 359 (1979); Phys. Rev. D 22, 2157 (1980).
- [13] K. Nakamura *et al.* (Particle Data Group), J. Phys. G 37, 075021 (2010).
- [14] S.J. Brodsky and G.P. Lepage, Phys. Rev. D 24, 1808 (1981).
- [15] B. Aubert *et al.* (*BABAR* Collaboration), Phys. Rev. D 80, 052002 (2009).
- [16] See e.g., S. J. Brodsky, F.-G. Cao, and G. F. de Teramond, Phys. Rev. D 84, 033001 (2011); S. J. Brodsky, F.-G. Cao, and G. F. de Teramond, Phys. Rev. D 84, 075012 (2011); A. P. Bakulev, S. V. Mikhailov, A. V. Pimikov, and N. G. Stefanis, Phys. Rev. D 84, 034014 (2011); P. Kroll, Eur. Phys. J. C 71, 1623 (2011); X.-G. Wu and T. Huang, Phys. Rev. D 82, 034024 (2010); H. L. L. Roberts, C. D. Roberts, A. Bashir, L. X. Gutierrez-Guerrero, and P.C. Tandy, Phys. Rev. C 82, 065202 (2010), and references therein.