

# Strong and electroweak supersymmetric corrections to single top processes at the Large Hadron Collider

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We present the one-loop corrections originating from QCD and electroweak interactions of supersymmetric origin within the minimal supersymmetric standard model to the single-top processes  $bq \rightarrow tq'$  and  $q\bar{q}' \rightarrow t\bar{b}$ . We illustrate their impact on top quark observables accessible at the LHC in the “ $t + \text{jet}$ ” final state, such as total cross section, several differential distributions, and left-right plus forward-backward asymmetries. We find that in many instances these effects can be observable for planned LHC energies and luminosities, quite large as well as rather sensitive to several minimal supersymmetric standard model parameters.

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## I. INTRODUCTION

Top quark processes at the LHC will be an ideal laboratory where it will be possible to profile with great accuracy the heaviest particle of the standard model (SM) and eventually extract possible effects of new physics (NP) beyond the SM. The manifestation of NP at the LHC will in fact depend on the actual NP mass scale,  $M_{\text{NP}}$ . If the available energy is less than  $M_{\text{NP}}$ , then NP will manifest itself through virtual effects. Conversely, if the available energy is greater than  $M_{\text{NP}}$ , NP will appear through the production and decay of new real particle resonances. A mixture of both scenarios may of course also occur.

If NP is identified with supersymmetry (SUSY), then it is fair to say, based on the available literature, that much has been done in the second respect. Various methods to extract the presence of new resonant SUSY particles have in fact been devised over the years and applied to several possible model realizations of a SUSY theory, particularly for the case of the simplest SUSY scenario, the minimal supersymmetric standard model (MSSM). In contrast, much less has been done so far in the first respect—primarily because the LHC is perceived not to be a precision machine, due to both the undefined partonic energy and the intrinsically large hadronic background, so that virtual effects (of order  $\alpha_s$  and  $\alpha$  or indeed smaller) are expected to be not easily discernible in the data. However, there are certain processes where the experimental precision is expected to eventually become comparable to the size of the virtual effects of NP.

Among these processes, one can certainly list top quark production, both in double- and single-top mode. In fact, with typical hadroproduction cross sections at the LHC of order 800 pb and 300 pb, respectively (at 14 TeV), and collider luminosities that can reach the  $300 \text{ fb}^{-1}$  level, several hundred million top events will altogether be pro-

duced during the lifetime of the CERN machine. Clearly, on the one hand, this renders the statistical error applicable on the experimental side to typical top quark observables negligible. Furthermore, on the other hand, the main systematics affecting the latter [both theoretical, coming from the parton distribution function (PDF) dependence, and experimental, coming from the jet energy scale uncertainty] will be understood at the percent level after the first hadronic data samples have been collected and studied in detail.

Therefore, both double- and single-top production processes at the LHC lend themselves to being precision physics laboratories, where, in particular, virtual effects of SUSY (when  $M_{\text{NP}} \equiv M_{\text{SUSY}} \gg m_t$ , where  $m_t$  is the top quark mass) can possibly be extracted, in the ultimate attempt to understand the underlying dynamics of SUSY breaking. We intend to illustrate here the potential afforded in this respect by two single-top production processes, namely, the  $t$  channel,

$$bq \rightarrow tq', \quad (1)$$

and the  $s$  channel,

$$q\bar{q}' \rightarrow t\bar{b} \quad (2)$$

(plus, of course, their charge conjugated channels), where  $q^{(\prime)}$  represents a light quark ( $u, d, s$ , or  $c$ )

The plan of the paper is as follows. The next section describes what is available in current literature concerning single-top processes at the LHC in terms of their higher order corrections within the MSSM. Section III illustrates how we performed the calculation of the one-loop SUSY QCD and electroweak (EW) corrections to the two single-top channels (1) and (2). The following section presents our numerical results. We conclude in Sec. V.

## II. SINGLE-TOP PROCESSES AT THE LHC

Because of their relevance for LHC physics, single- as well as double-top processes have extensively been studied in previous years, with a twofold aim: on the one hand, in order to have a precise SM prediction, at least at the complete next-to-leading order (NLO) (and possibly next-to-NLO for the QCD corrections), and, on the other hand, in order to investigate possible deviations from SM predictions due to the presence of NP. These two goals are obviously related, the former being a precondition for the second: in order to extract meaningful information from NP processes entering at one-loop level top quark hadro-production at the LHC, clearly, all similar SM effects should be well under control.

One-loop SM corrections to single- and double-top production in hadron-hadron collisions have been known for some time. These include the SM QCD [1] and SM EW [2] ones to  $gg$ ,  $q\bar{q} \rightarrow t\bar{t}$  (double-top production) and the corresponding ones [3,4] for  $bq \rightarrow tq'$ ,  $q\bar{q}' \rightarrow t\bar{b}$ , and  $bg \rightarrow tW^-$  (single-top production). Both SM QCD and EW corrections can be large, although in complementary energy regimes  $\sqrt{s} \approx m_t$  and  $\sqrt{s} \gg m_t$ , respectively. However, they all are rather stable against variations of the factorization and/or renormalization scales (these in turn quantifying the systematic uncertainties related to the unknown two-loop, or even higher order, corrections), so that one can conceivably attempt to investigate virtual effects of some NP, particularly of SUSY origin, induced onto single- and double-top processes. To stay with the MSSM, QCD and EW corrections to double-top production have been quantified in [5,6], respectively, while only part of these are known to date for the case of single-top reactions [7].

Our paper is thus the first one where both MSSM QCD and EW corrections to single-top processes inducing a “ $t$  plus jet” signature in the final state are computed (i.e.,  $bq \rightarrow tq'$  and  $q\bar{q}' \rightarrow t\bar{b}$ ). It is in fact from this perspective that we decided to postpone to another publication the presentation of similar corrections to the third single-top channel (i.e.,  $bg \rightarrow tW^-$ ), in the sense that this process originates an altogether different signature in the final state, “ $t$  plus lepton and missing (transverse) energy” or “ $t$  plus two jets,” depending on whether the  $W^\pm$  boson produced in association with the top quark decays leptonically or hadronically. In fact, both of the  $tW^-$  generated signatures require different triggers and selection procedures with respect to the  $t$  + jet one, so that, from a phenomenological point of view, they deserve a separate treatment.

## III. CALCULATION

The values of the SM input parameters considered for the numerical evaluation of the one-loop corrections can be found in Table I.

TABLE I. Numerical values of SM inputs.

Coupling constants	$\alpha(M_Z) = 1/127.934\,006\,52$ , $\alpha_s(M_Z) = 0.1176$
Gauge boson masses	$M_W = 80.424$ GeV, $M_Z = 91.1876$ GeV
Heavy quark masses	$m_t = 170.9$ GeV, $m_b^{\overline{\text{MS}}}(m_b) = 4.2$ GeV

We used the CTEQ6 2006 PDF set [8], in particular, the CTEQ6L fit, which is leading order (LO) in QCD, choosing renormalization and factorization scales both equal to the top mass,  $m_t$ .

As intimated in the introduction, due to the relatively large single-top production cross sections (in particular, for the  $t$ -channel process) and the consequent large available statistics, we reckon that it will be possible at the LHC to perform “precision physics” of top samples produced singly. It is therefore a worthy task trying to understand systematically if sizable virtual effects could modify the SM predictions, providing an indirect but encouraging hint of NP. As a first step we performed an adaptive scan over both the MSSM and minimal-super-gravity (mSUGRA) parameter spaces in order to investigate whether there exist regions with large one-loop corrections. We have done so limitedly to the inclusive production cross sections in either channel.

While in the MSSM case we have input the relevant parameters directly at the EW scale, in the mSUGRA scenario the low energy spectra have been obtained evolving the input parameters from the grand unification theory scale down to the EW scale through the code SUSPECT [9].

We have developed a dedicated C++ code to compute the one-loop corrections, and as an internal check, we have tested the cancellation of ultraviolet divergences appearing in the loop integrations. To quantify the one-loop corrections, we will mainly focus on  $K$  factors, defined in general as the ratio NLO/LO of a given observable, with all relevant quantities (QCD and EW coupling constants, etc.) evaluated at the given order consistently. The PDFs, on the other hand, have been kept at LO since their NLO effects are of SM origin, whereas we are looking here at purely SUSY corrections; therefore using NLO QCD PDF sets would introduce spurious SM effects from higher orders.

Guided by the inclusive results, we have then looked at differential distributions for some “benchmark points.” In particular, alongside the typical kinematical observables [invariant masses, transverse momenta, (pseudo)rapidities], due to the possibility of measuring spin-related observables in top samples [the top (anti)quark decays before hadronizing in fact, thus efficiently transmitting its polarization state onto the ensuing decay products], we have also focused our attention on the following observables:

- (i) The left-right asymmetry ( $A_{LR}$ ).
- (ii) The forward-backward asymmetry ( $A_{FB}$ ), only for the  $s$  channel.

$A_{LR}$  is defined as usual by the following ratio:

$$A_{LR} = \frac{\sigma_{pp \rightarrow t_L + X} - \sigma_{pp \rightarrow t_R + X}}{\sigma_{pp \rightarrow t_L + X} + \sigma_{pp \rightarrow t_R + X}}. \quad (3)$$

$A_{FB}$  deserves further details, since the definition of “forward” and “backward” regions in a  $pp$  collider is not straightforward. In the  $s$ -channel process the initial state can essentially be  $u\bar{d}$  or  $c\bar{s}$ . In the first case  $u$  is a

TABLE II. Input parameters for the mSUGRA benchmark points (all values with mass dimension are in GeV).

mSUGRA scenario	$m_0$	$m_{1/2}$	$A_0$	$\tan\beta$	$\text{sign}\mu$
LS2	300	150	-500	50	+
SPS1a	100	250	-100	10	+

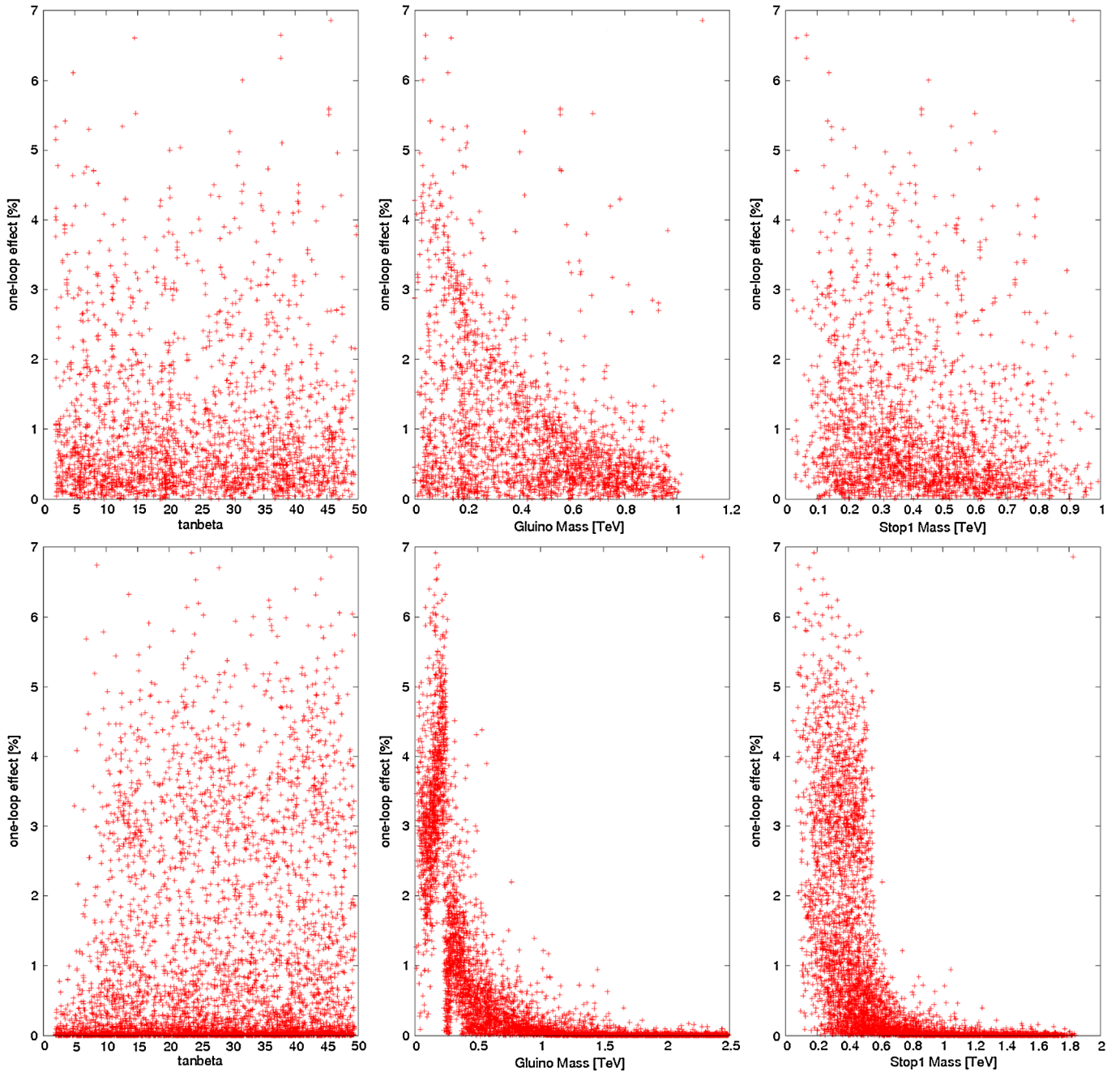


FIG. 1 (color online). Scan over three relevant SUSY parameters ( $\tan\beta$ ,  $m_{\tilde{g}}$ , and  $m_{\tilde{t}}$ ) in the context of the MSSM (upper panels) and mSUGRA (lower panels) to search for maxima of one-loop effects.

valence quark, and thus the momentum fraction it carries is much bigger than that of the  $\bar{d}$  sea quark. It is then possible to reconstruct here the direction of the incoming valence quark just by identifying the direction of the boosted top in the final state. This is a statistical process, since  $s$ -channel

single-top production can be initiated by two sea quarks too, such as in the  $c\bar{s}$  case, and the weight of the various events is given by the relevant PDF (hereafter denoted as  $P$ ). The definition of  $A_{\text{FB}}$  is then given by the following relation [10]:

$$A_{\text{FB}} = \frac{\int [dx_1 dx_2 \hat{\sigma}(pp \rightarrow u\bar{d}t\bar{b} + X)(P_u(x_1, \mu)P_{\bar{d}}(x_2, \mu) - P_{\bar{d}}(x_1, \mu)P_u(x_2, \mu))\text{sign}(x_1 - x_2)]}{\int dx_1 dx_2 \sum_{q=u,d,c,s} (P_q(x_1, \mu)P_{\bar{q}}(x_2, \mu) + P_{\bar{q}}(x_1, \mu)P_q(x_2, \mu))}, \quad (4)$$

where  $\mu$  is the factorization/renormalization scale.

We also have computed differential distributions of these quantities over the invariant mass of the final state, the transverse momentum of the top as well as the rapidity of both the top and the light quark in the final state.

Coming back to the scans, in the general MSSM scenario, we have explored over all soft parameters involved in the generation of the spectra, while in the mSUGRA scenario we have scanned over the four standard mSUGRA input parameters. The scans have been performed interfacing the tool ADSCAN[11] to our code.

The main result of this preliminary analysis, both for MSSM and mSUGRA, is that the relative effect of the SUSY one-loop corrections seems to be rather small over the whole range of the independent parameters. Moreover, the effects show a very smooth and mild dependence on the parameters. A meaningful sample of our results is shown in Fig. 1. Each point in the diagrams corresponds to a full computation of NLO effects for a point in the parameter space, plotted versus a low energy parameter, such as a mass or mixing angle evaluated at the EW scale. We have chosen to show just three of them, corresponding to the parameters which are expected to affect most heavily the NLO corrections:  $m_{\tilde{g}}$  (the gluino mass),  $\tan\beta$  (the ratio of the two doublet Higgs vacuum expectation values), and  $m_{\tilde{t}_1}$  (the lightest stop mass).

From the  $\tan\beta$  plot it is easy to see that the density of the points is larger for very small values of the corrections and rapidly decreases for larger values of the effect. Furthermore, and possibly contrary to naive expectations, the density of the points is almost  $\tan\beta$  independent. In essence, for the largest part of the parameters' space, the corrections are very small, below the 2% level, while it is unlikely to find parameters' configurations leading to corrections as high as 5%, and this result holds independently of  $\tan\beta$  (i.e., there are no visible regions in the  $\tan\beta$  range where one-loop effects have a maximum). Trivially, the plots of the corrections as a function of the gluino or stop masses simply state that it is more likely to find "larger" (6% or so at the most) corrections in parameters' regions where these masses are small.

Despite the results obtained from the scan never being very large, it should be noted that these refer to inclusive corrections only. However, while scanning over the MSSM

and mSUGRA parameter spaces, we have also tested for the size of SUSY QCD and EW corrections to differential distributions (including asymmetries) and found them to be rather large, albeit in limited (and at times disfavored) regions of phase space, thereby explaining their smallness in the inclusive results. In order to illustrate the typical pattern emerging at a differential level, we have defined a couple of representative benchmark mSUGRA points, whose parameters are given in Table II.

The SPS1a point [12] is the typical mSUGRA "standard candle," while the LS2 point has been introduced in [13] and is characterized by an interestingly light spectrum

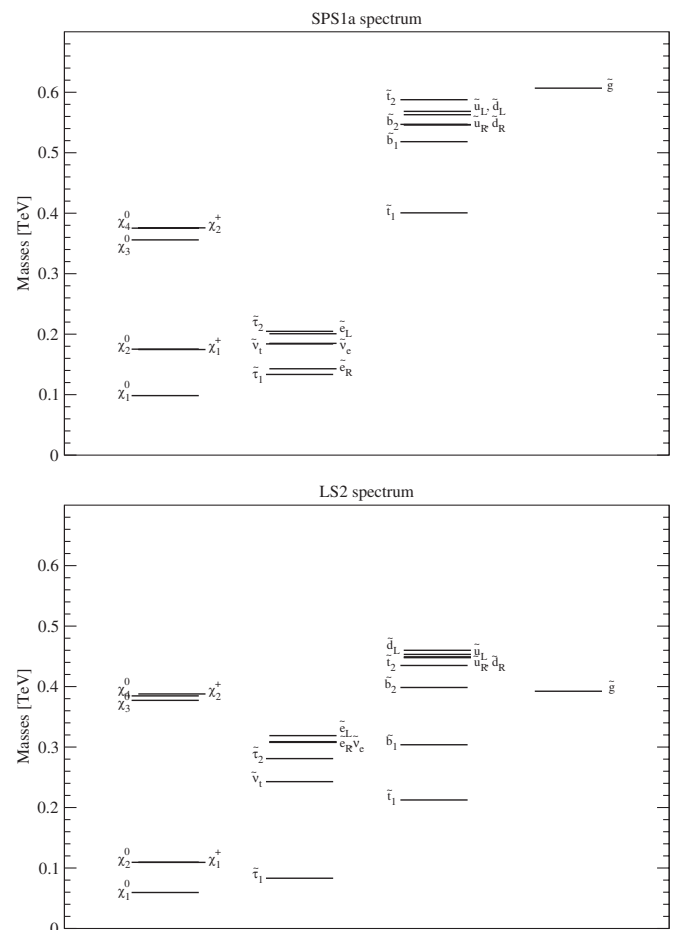


FIG. 2. The SPS1a and LS2 mass spectra.

TABLE III. Numerical results for  $t$ -channel production at 7 TeV.

	$\sigma$ (pb)	L-R Asymmetry ( $1 - \delta$ )
Born	35.44	$\delta = 4.7327 \times 10^{-7}$
LS2	35.29 ( $K = 0.996$ )	$\delta = 4.7515 \times 10^{-7}(1 - K = 1.87 \times 10^{-9})$
SPS1	35.33 ( $K = 0.997$ )	$\delta = 4.7353 \times 10^{-7}(1 - K = 2.6 \times 10^{-10})$

TABLE IV. Numerical results for  $s$ -channel production at 7 TeV.

	$\sigma$ (pb)	L-R Asymmetry	F-B Asymmetry
Born	2.061	0.6777	0.598 049
LS2	2.086 ( $K = 1.0121$ )	0.6796 ( $K = 1.0028$ )	0.598 018 ( $K = 0.999 948$ )
SPS1	2.062 ( $K = 1.0005$ )	0.6783 ( $K = 1.0009$ )	0.598 046 ( $K = 0.999 995$ )

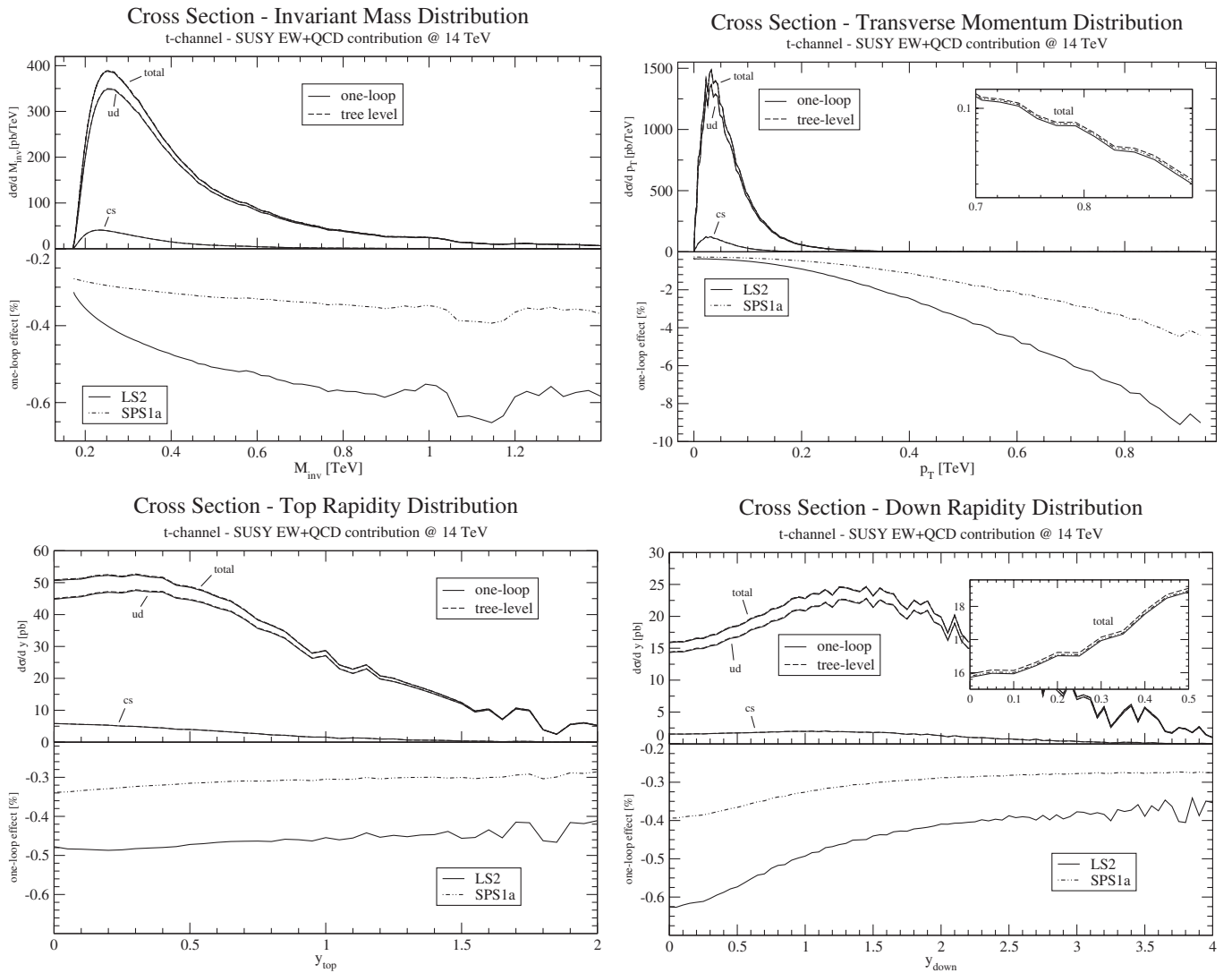


FIG. 3. Differential distributions of the cross section in the  $t$  channel at 14 TeV. In this and in the following figures, the small panels focus on the range where the one-loop effects are higher, and they are shown only when the corrections from LS2 and SPS1a are distinguishable.

TABLE V. Numerical results for  $t$ -channel production at 14 TeV.

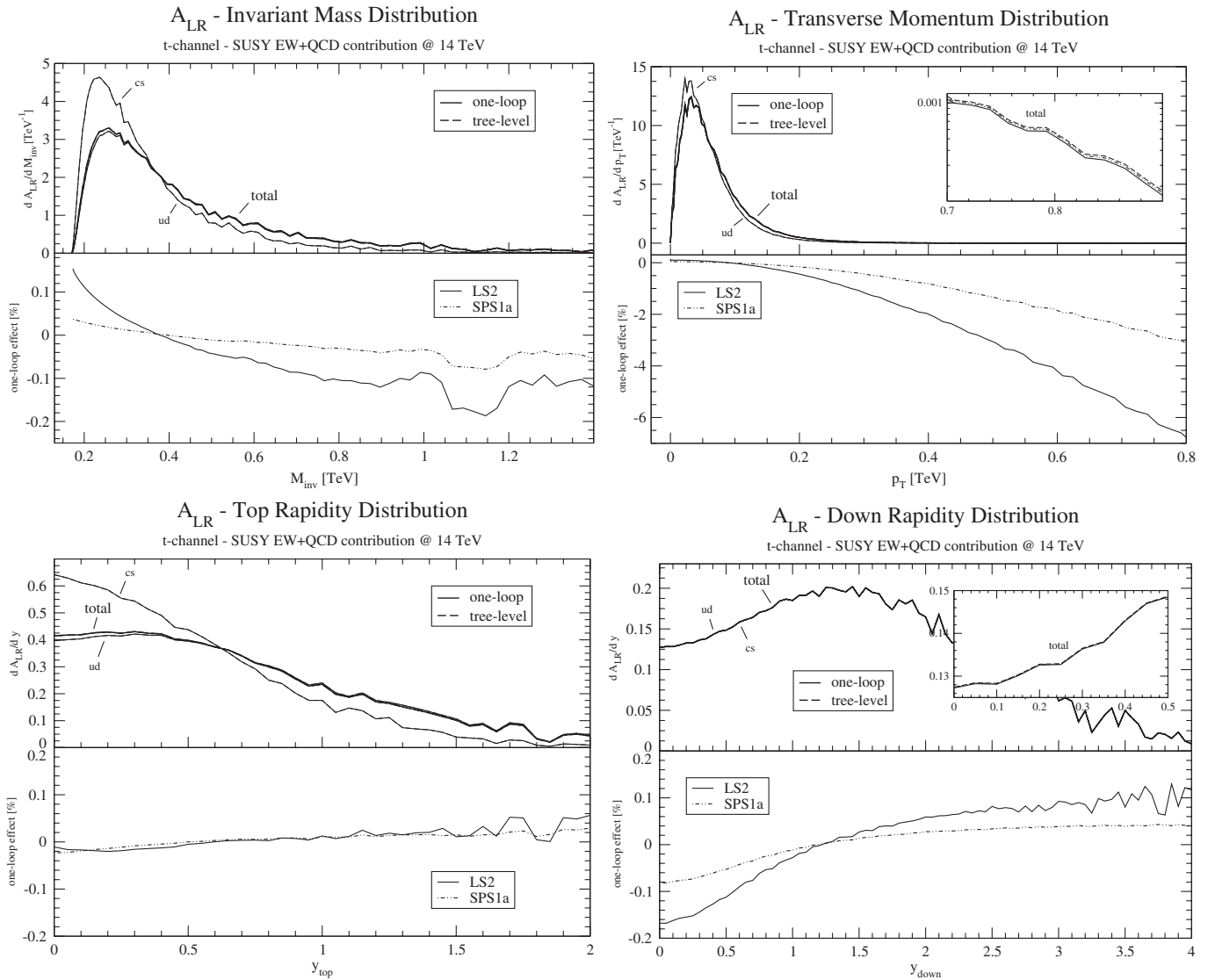
	$\sigma$ (pb)	L-R Asymmetry ( $1 - \delta$ )
Born	122.5	$\delta = 4.3491 \times 10^{-7}$
LS2	122.0 ( $K = 0.996$ )	$\delta = 4.3668 \times 10^{-7}$ ( $1 - K = 1.77 \times 10^{-9}$ )
SPS1	122.1 ( $K = 0.997$ )	$\delta = 4.3516 \times 10^{-7}$ ( $1 - K = 2.5 \times 10^{-10}$ )

(yet compliant with current experimental and theoretical constraints), which—if realized in nature—could prompt for easily detectable signals at the LHC. More details about the two spectra can be found in Fig. 2.

#### IV. RESULTS

A long LHC run at 7 TeV has recently been scheduled; it will probably last two years and will hopefully lead directly to the final expected energy of 14 TeV. In view of this

recent development, we present our results for both such planned energies. Our numerical results will be illustrated in turn in the following two subsections. Notice that we will not show all the differential distributions analyzed for the two benchmark points; rather we will illustrate a meaningful sample of these (and only in the 14 TeV case) and focus on the one-loop corrections they receive. We will also show the contribution of the SUSY EW and QCD parts separately, in some instances at least, to better explain the features of the underlying dynamics, like,

FIG. 4. Differential distributions of the left-right asymmetry in the  $t$  channel at 14 TeV.

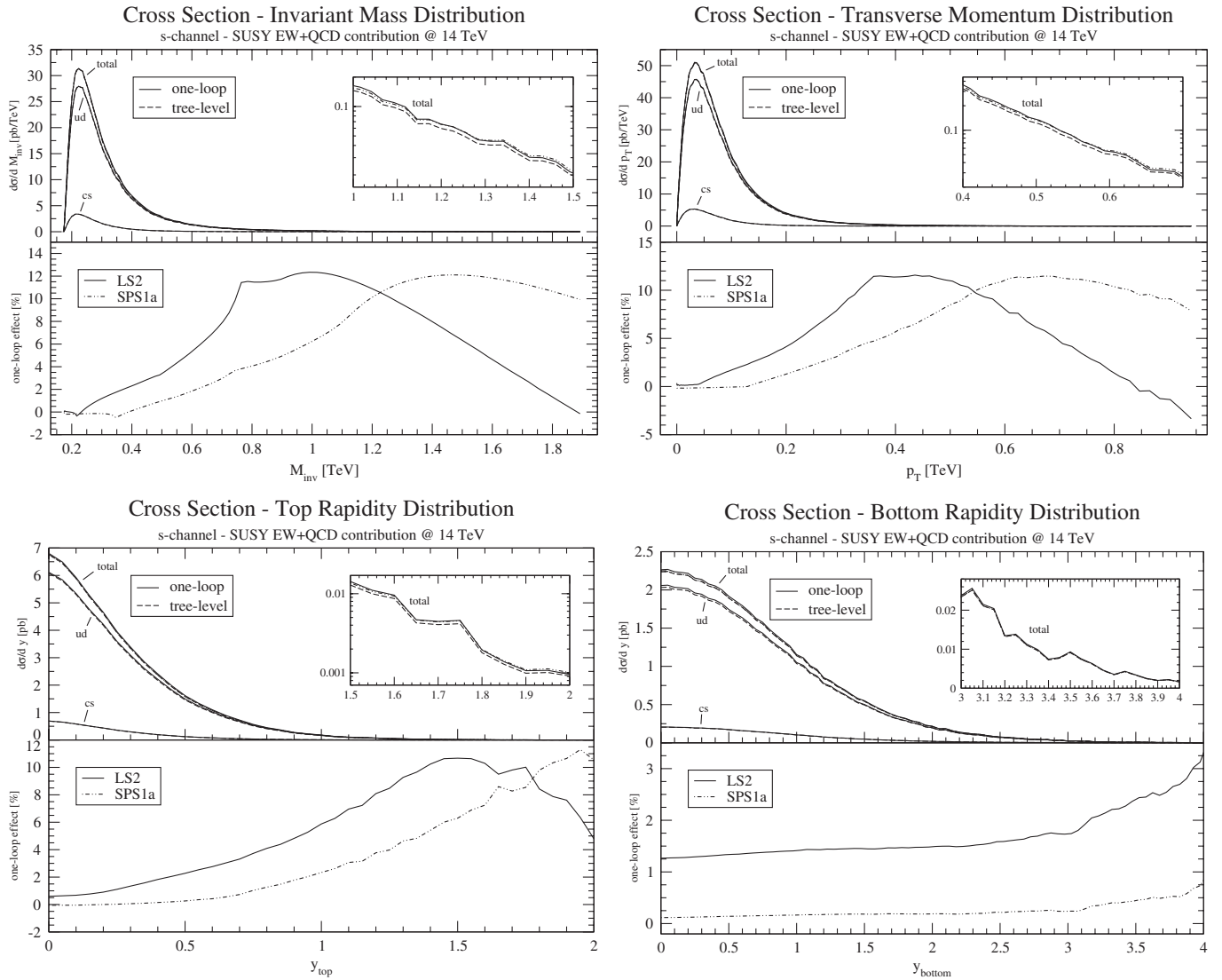


FIG. 5. Differential distributions of the cross section in the  $s$  channel at 14 TeV.

e.g., dominant contributions, asymptotic behaviors, and threshold effects.

### A. LHC at 7 TeV

The total cross section and asymmetries for  $t$  and  $s$  channels are given in Tables III and IV respectively.

As it is possible to see, in both the channels, the NLO corrections to the total cross sections are opposite in sign: in the  $t$ -channel case one-loop corrections reduce the cross

section, while in the  $s$  channel there is an enhancement, but the corrections are 1% or less in both the benchmarks considered (in line with the results from the mSUGRA scans). The main difference between the two channels is of course the value of the cross section,  $\sigma$ . Focusing on the  $t$  channel, which has a greater  $\sigma$ , we have found that, with an integrated luminosity of  $1 \text{ fb}^{-1}$ , the LHC will generate  $\approx 350\,00$  events per year; therefore the detection of any deviation from the SM prediction due to the SUSY EW and

TABLE VI. Numerical results for  $s$ -channel production at 14 TeV.

	$\sigma$ (pb)	L-R Asymmetry	F-B Asymmetry
Born	5.138	0.6861	0.536 65
LS2	5.206 ( $K = 1.0132$ )	0.6883 ( $K = 1.0032$ )	0.536 95 ( $K = 1.000 56$ )
SPS1	5.145 ( $K = 1.0014$ )	0.6869 ( $K = 1.0012$ )	0.536 76 ( $K = 1.000 20$ )

QCD one-loop corrections, at least in the early stages of the experiment, will be indeed very challenging. Also  $A_{LR}$  shows different behaviors in the two channels, but unfortunately, they are much greater in the  $s$  channel, which, as we have just verified, has NLO corrections well beyond observability at 7 TeV. The same observability considerations can be applied to  $A_{FB}$  in the  $s$  channel.

## B. LHC at 14 TeV

### 1. The $t$ -channel case

The total cross section and asymmetries for this production mechanism are given in Table V.

In this case the total cross section is significantly bigger than in the 7 TeV case, and therefore we expect that it will be possible to extract much more information from experimental data. The  $K$  factor for the integrated cross section is

admittedly small, less than 1% for both benchmarks, but with the expected luminosity of  $10 \text{ fb}^{-1}$ , the production of single-top through this  $t$ -channel process will count  $\sim 1$  M events, and therefore it should be possible to detect a discrepancy in the expected number of SM events that could be interpreted as a hint of NP and thus boost the search of a SUSY resonance. As far as the  $A_{LR}$  is concerned, the top quark will emerge almost completely left-polarized, yet the detection of a one-loop induced discrepancy in the asymmetry will be again unlikely, the  $K$  factor being close to 1 part in a billion or less.

The analysis of differential observables is now mandatory to better understand the origin of the corrections to the integrated cross section and left-right asymmetry. In Fig. 3 we have collected our results for the differential distributions of the cross section. It can be noticed that the corrections are always bigger in the LS2 scenario than in

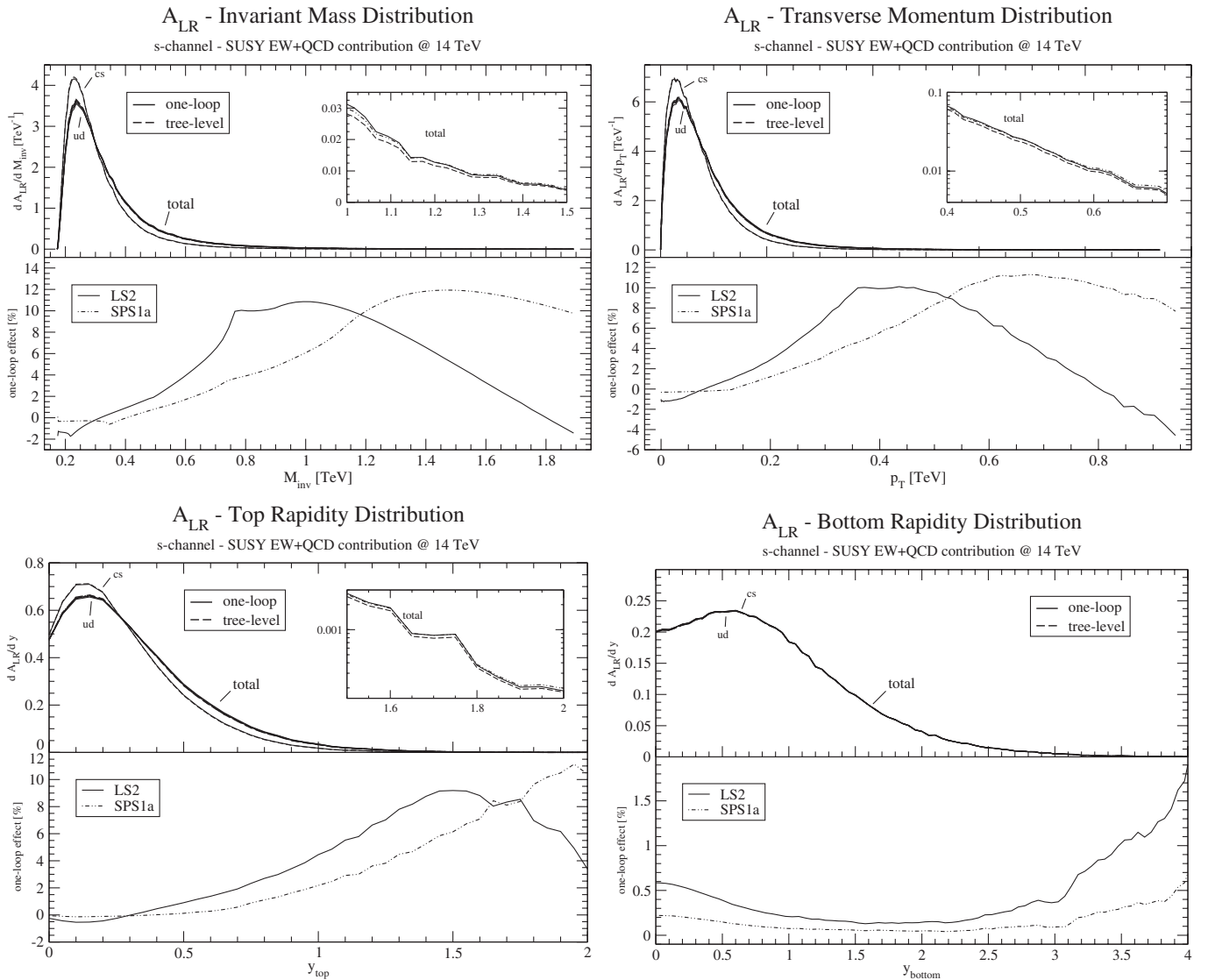


FIG. 6. Differential distributions of the left-right asymmetry in the  $s$  channel at 14 TeV.



the SPS1a case and that in the transverse momentum distribution they can reach the  $-10\%$  limit for high values of the top  $p_t$  (transverse momentum). However, the differential cross section is quite small in this limit, of the order of  $0.01$  pb/TeV, and therefore observation of large one-loop corrections for events with high top transverse momentum is probably possible though quite difficult. It is also possible to see that the one-loop corrections in the other differential distributions are always quite small, below the  $1\%$  or so limit.

In Fig. 4 the one-loop corrections to the differential distributions for the left-right asymmetry are shown. Again, the corrections are generally bigger in the LS2 scenario than in the SPS1a case and can reach significant values for high top  $p_t$ , of  $-8\%$  or so. The  $p_t$  distribution has another interesting feature: the asymmetry is bigger in

the high top  $p_t$  region, where, however, one-loop corrections are much less than  $1\%$ .

## 2. The $s$ -channel case

The total cross section and asymmetries for this production mechanism are given in Table VI.

Single-top events from the  $s$  channel at the LHC will be significantly more than at the Tevatron, and this will allow a much more precise analysis of the process and of deviations to related observables due to new physics. The inclusive one-loop corrections we have obtained are at most of  $\sim 1\%$  in the LS2 scenario, but may be enough to allow a small observable difference in the expected number of events. In contrast, the left-right asymmetry receives NLO corrections of the order of  $0.1\%$  in both the scenarios considered (here, unlike the general case, the effect of

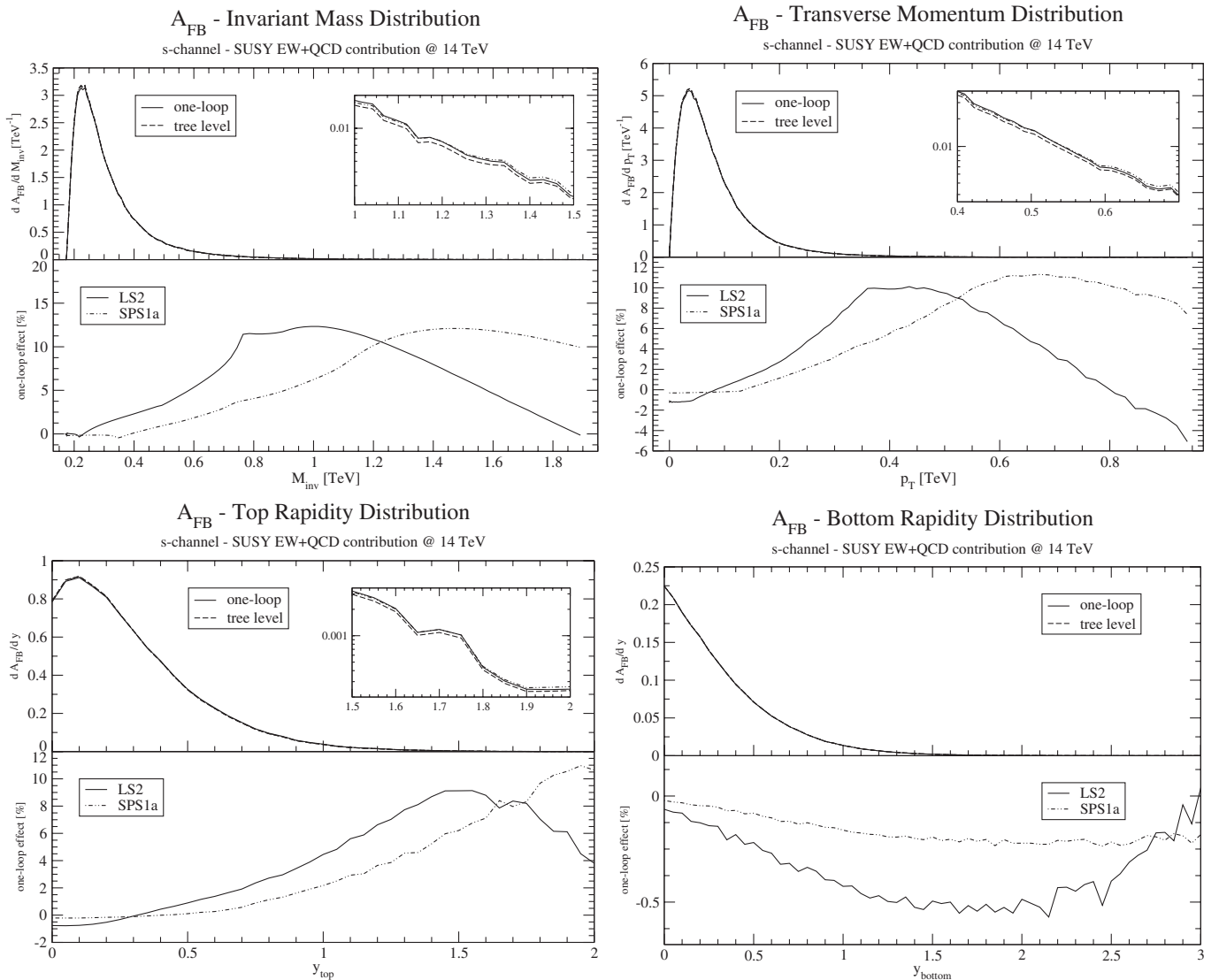


FIG. 7. Differential distributions of the forward-backward asymmetry in the  $s$  channel at 14 TeV.

pure SUSY corrections is almost completely negligible). The same result holds for the one-loop corrections to the forward-backward asymmetry: the  $K$  factor is very small and the effect practically undetectable with the  $s$ -channel cross section being of some pb.

The differential distributions for the cross section are shown in Fig. 5. In the  $s$ -channel case there are significant differences between the two scenarios considered: the corrections for LS2 are higher than those for SPS1a in the regions where the contributions to the total cross section are bigger, and this is the reason for the big difference in the  $K$  factors for the integrated cross section in the two scenarios.

The same feature can be observed for the left-right asymmetry distributions, Fig. 6, where, however, the enhancement of the one-loop corrections from the LS2 scenario is milder and therefore the differences in the integrated asymmetry are smaller.

The corrections to differential quantities in the forward-backward asymmetry are qualitatively and quantitatively very similar to those for the left-right asymmetry, as it is possible to see in Fig. 7; thus the same comments as above apply in this case.

Considering all the differential distributions for the  $s$ -channel process, we can notice that they do not differ too much from each other between the two benchmarks and, as a general feature, they reach their maximum value in the invariant mass distribution for  $M_{\text{inv}} \gtrsim 700$  GeV, in the transverse momentum distribution for  $p_T \gtrsim 300$  GeV, and in the top rapidity for  $y_t \gtrsim 1.5$ .

To better understand the origin of the corrections and why there are peaks for certain values of  $M_{\text{inv}}$ ,  $p_T$ , and  $y_t$ , we have computed (here, limitedly to the  $s$  channel) the SUSY EW and QCD corrections for the cross section separately and plotted the related differential quantities in Figs. 8 and 9 for the two benchmark scenarios. It is

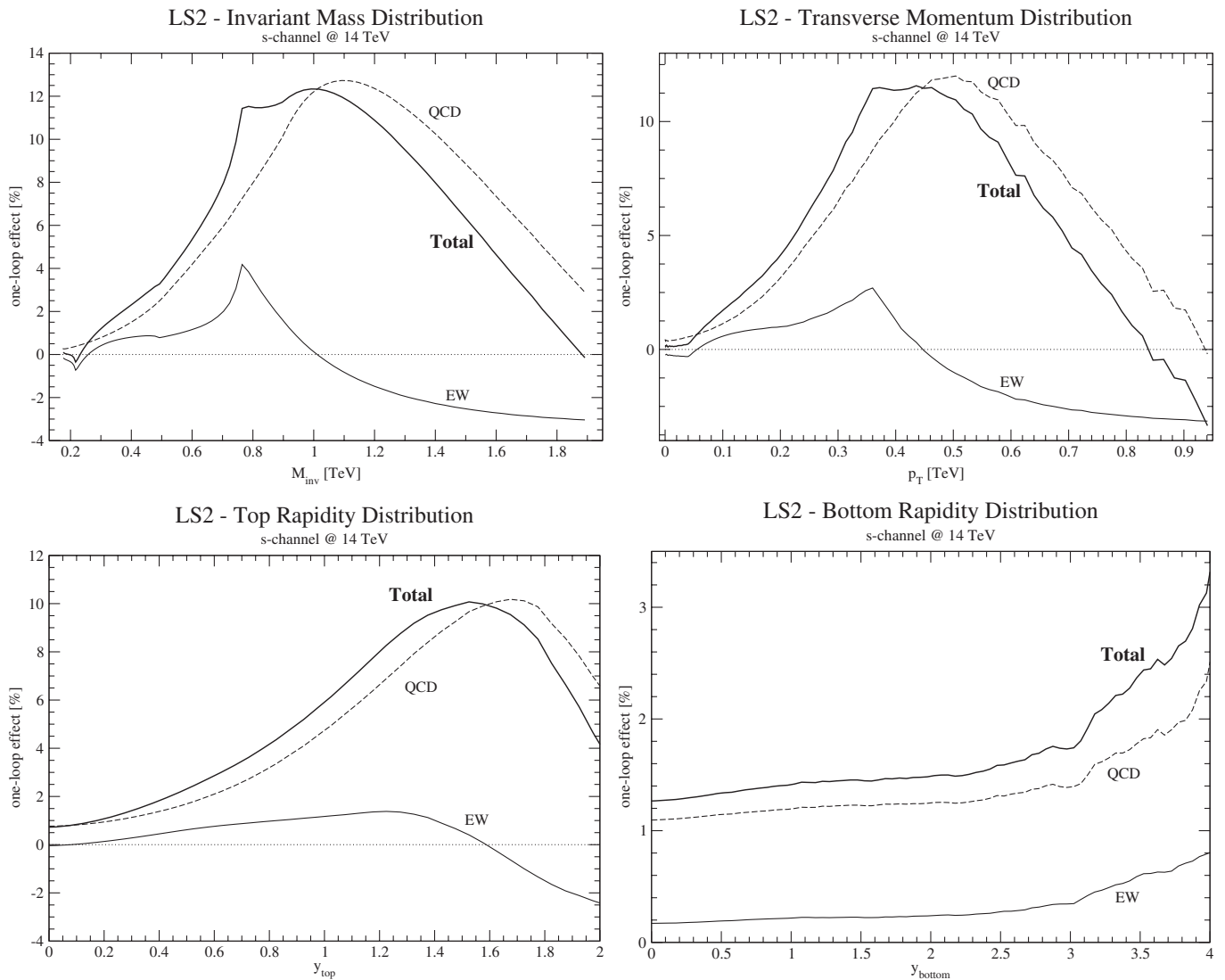


FIG. 8. SUSY EW and QCD corrections to differential distributions of the cross section in the  $s$  channel at 14 TeV for LS2.

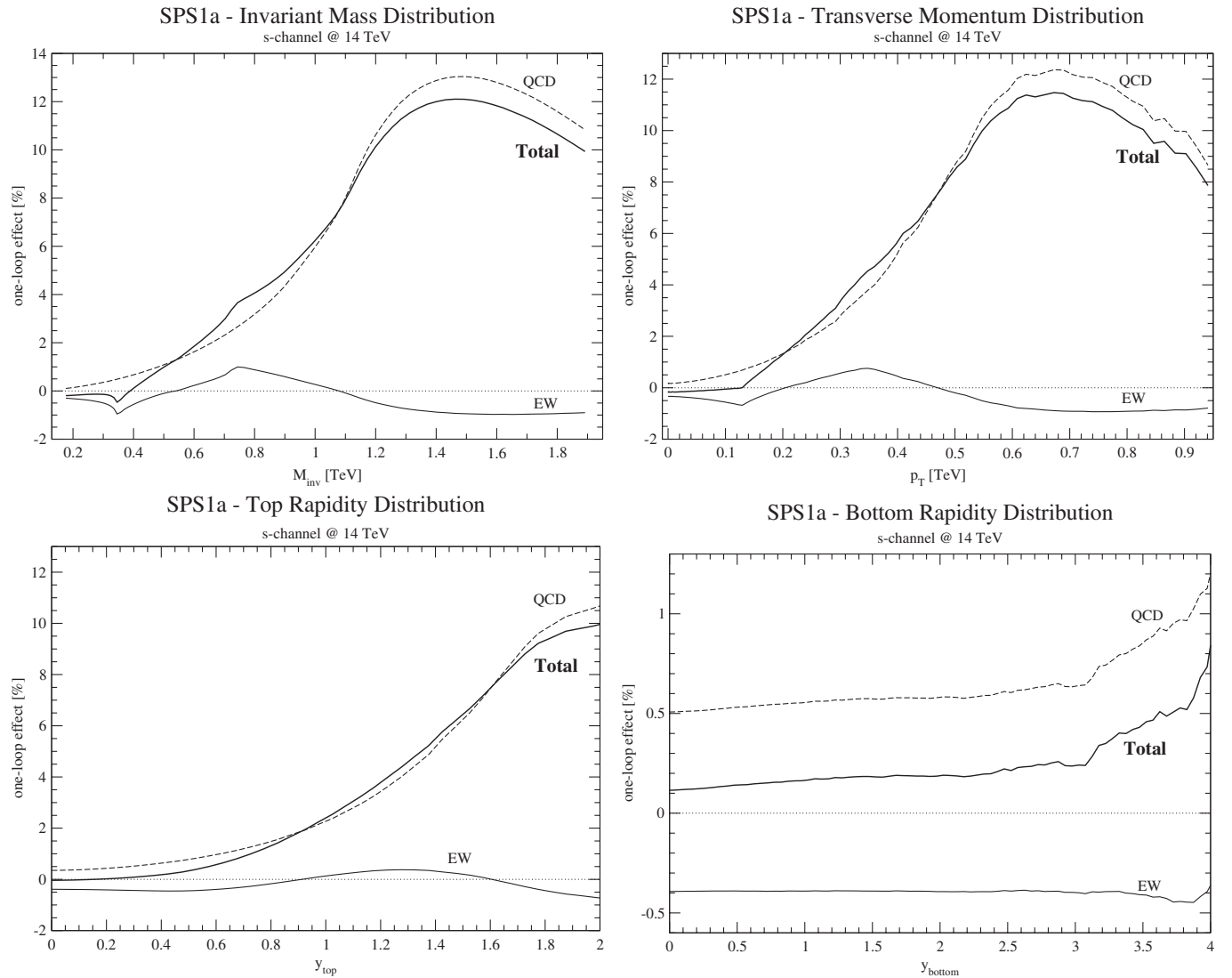


FIG. 9. SUSY EW and QCD corrections to differential distributions of the cross section in the  $s$  channel at 14 TeV for SPS1a.

possible to see that the corrections are dominated by the SUSY QCD contribution, while the EW counterpart provides small contributions to the total corrections. In both benchmark scenarios, however, the EW contribution shows peculiar features in the invariant mass distribution, given by peaks and troughs in the correction, and a closer inspection of these threshold effects reveals that they are situated in correspondence of  $M_{\chi_i^\pm}^2$ . Therefore charginos could play an interesting role in the determination of SUSY effects in  $s$ -channel single-top production processes. Going back to the dominant QCD contribution, the only parameter which enters the QCD correction alone is the gluino mass ( $m_{\tilde{g}} = 607$  GeV in SPS1a and  $m_{\tilde{g}} = 392$  GeV in LS2); thus we can argue that the one-loop corrections are sensitive mostly to this parameter. Smaller gluino masses shift the SUSY QCD peak in the corrections towards regions where the differential distributions have higher values and therefore the integrated quantities will be af-

fected more by NLO corrections. This means that observing higher pure SUSY one-loop corrections would point towards scenarios with lighter gluino and vice versa.

To conclude this section, we are aware that recent measurements have slightly changed the value of the top mass [14]. Nevertheless, we have verified that any change (within a reasonable range) in this value has practically no effect on one-loop corrections, as it is possible to see in Fig. 10.

## V. CONCLUSIONS

In summary, QCD and EW corrections through the one-loop level originating from SUSY and affecting the cross section for “ $t$ + jet” final states at the LHC are always small at the inclusive level, at energies of both 7 and 14 TeV. However, they appear sizable (up to 10%) and detectable [after  $\mathcal{O}(100)$  fb $^{-1}$  of luminosity] at the differential level, if the collider energy is 14 TeV, while for

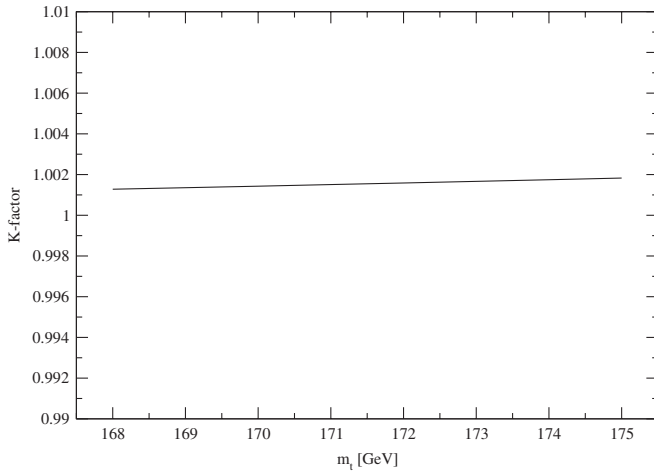


FIG. 10.  $K$ -factor dependence on the top mass in the  $s$  channel at 14 TeV for SPS1a.

7 TeV they will presumably not be accessible. Such large corrections also appear in the left-right and forward-backward (when definable) asymmetries, again, at the differential but not inclusive level. The bulk of the higher

order effects comes from QCD, with a much smaller, yet very sensitive to SUSY parameters, EW component. The corrections discussed here are typically larger for the  $s$ -channel mode than for the  $t$ -channel one, which renders them more difficult to access as the former is subleading with respect to the latter at the LHC. Altogether though, our results point to the relevance of SUSY one-loop effects on “ $t+$  jet” final states for the LHC running at design energy and luminosity. In the presence of high statistic samples, some typical SUSY parameter dependencies could possibly be extracted, in either mSUGRA or the MSSM.

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