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<sup>4</sup>R. Aaron, R. D. Amado, and Y. Y. Yam, Phys. Rev. 136, B650 (1964).

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<sup>8</sup>D. Harrington, Phys. Rev. 139, B691 (1965).

<sup>9</sup>C. W. Gardiner, Phys. Rev. D 4, 1391 (1971).

<sup>10</sup>S. Ali, M. Rahman, and D. Husain, Phys. Rev. D 6, 1178 (1972).

<sup>11</sup>D. Husain and S. Ali, Phys. Rev. C 2, 1587 (1970).

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## Errata

**Hadronic Final States of Deep-Inelastic Processes in Simple Quark-Parton Models**, John Kogut, D. K. Sinclair, and Leonard Susskind [Phys. Rev. D 7, 3637 (1973)].

1. Page 3655: At the end of Case 2 the following sentence should be added: The final mass-shell condition  $(k_n + p)^2 = m^2$  requires that  $\alpha = \pm 1$ .

2. Page 3656: In Case 2 the three parenthetic remarks “(or  $Q - k_{j+} \leq \sim Q^{-\alpha}$ ),” “(or  $Q - k_{i+} \leq \sim Q^{-\alpha}$ ),” and “[or  $k_{i+} = Q + O(Q^{-\alpha})$ ]” should be deleted.

3. Page 3650: Directly before Acknowledgments, the following sentence should be inserted: We have recently had brought to our attention related work in the context of the nonperturbative parton model.<sup>39,40</sup>

4. Page 3657: The following two references should be added to the end of the list:

<sup>39</sup>P. V. Landshoff and J. C. Polkinghorne, Nucl. Phys. B33, 221 (1971).

<sup>40</sup>C. F. A. Pantin, Nucl. Phys. B46, 205 (1972).

**Nonlinear Bootstrap in Dual Models**, J. Maharana and R. Ramachandran [Phys. Rev. D 7, 3670 (1973)]. There is a need for a clarifying note and an erratum in the above paper. For the process  $\pi^- K^+ \rightarrow \pi^- K^+$  it is possible to consider the corresponding inclusive reaction  $\pi^- + K^+ \rightarrow K + \text{anything}$ , rather than the reaction  $\pi^- + K^+ \rightarrow \pi^0 + \text{anything}$  described in Sec. IID. This then requires the use of strangeness-conservation sum rules, and the bootstrap equation (2.12), now applicable, takes the form

$$\sigma_K^{\text{Res}}(\pi^- K^+) = \sum_{K=K^+, K^0} \int d^3 p_K \frac{d\sigma^{\text{Res}}}{d^3 p_K}(\pi^- K^+ \rightarrow K + x) - \sum_{\bar{K}=K^-, \bar{K}^0} \int d^3 p_{\bar{K}} \frac{d\sigma^{\text{Res}}}{d^3 p_{\bar{K}}}(\pi^- K^+ \rightarrow \bar{K} + x).$$

The main advantage in using this sum rule rather than the one used in text is that it justifies the neglect of the contribution arising from Fig. 2(d) to the right-hand side of Eq. (2.12). Apart from a different set of Chan-Paton factors, this introduces no change in the body of the paper. As a result there is a factor  $\frac{1}{2}$  which multiplies the right-hand side of Eq. (3.9), and a factor 2 multiplies the right-hand side of Eq. (3.10).

This changes the result  $G_{K^* K^* \pi^2}/4\pi$  from the value of 1.99 to

$$\frac{G_{K^* K^* \pi^2}}{4\pi} = 3.98,$$

which is in fact in better agreement with the experimental value of  $3.25 \pm 1.25$ .

Similarly, for  $\pi\eta \rightarrow \pi\eta$ , described in Sec. IIE, the charge-conservation sum rule may be used. The resulting change in the Chan-Paton factor is  $\frac{1}{6}$ , thus yielding a value for  $G_{A_2 \rho \pi^2}/4\pi$  given by

$$\frac{G_{A_2 \rho \pi^2}}{4\pi} = 610 \text{ GeV}^{-4}$$

(instead of  $101.68 \text{ GeV}^{-4}$ ), which is to be compared with the experimental number of  $313.4 \text{ GeV}^{-4}$ .

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