Light quark mass ratio from the Dalitz plot of $\eta \rightarrow \pi^+ \pi^- \pi^0$ decay

B. V. Martemyanov and V. S. Sopov

Institute for Theoretical and Experimental Physics, B. Cheremushkinskaya 25, 117259 Moscow, Russia (Received 26 November 2004; published 13 January 2005)

High statistics Dalitz-plot distribution of $\eta \to \pi^+ \pi^- \pi^0$ decay obtained recently by KLOE collaboration is fitted to the results of corresponding theoretical calculations in Chiral Perturbation Theory (ChPT) with unitarity corrections taken into account. The quark mass ratio $Q = \sqrt{[m_s^2 - (m_d + m_u)^2/4]/(m_d^2 - m_u^2)}$ can be obtained from this analysis. We get $Q = 22.8 \pm 0.4$ which differs from the value $Q_{DT} = 24.2$ that follows from Dashen's theorem and agrees with recently calculated electromagnetic kaon mass difference.

DOI: 10.1103/PhysRevD.71.017501

The possibility to extract light quark mass difference from $\eta \rightarrow \pi^+ \pi^- \pi^0$ decay is known for a long time [1]. In Chiral Perturbation Theory (ChPT) the decay width Γ depends on quark mass ratios and theoretically calculable factor $\overline{\Gamma}$ [2]:

$$\Gamma = \left(\frac{Q_{DT}}{Q}\right)^4 \bar{\Gamma},\tag{1}$$

where

$$Q^{-2} = \frac{m_d^2 - m_u^2}{m_s^2 - \hat{m}^2}, \quad \hat{m} = \frac{m_d + m_u}{2}, \tag{2}$$

 m_u, m_d, m_s are up, down and strange quark masses,

$$Q_{DT}^{-2} = \frac{\left[(m_{K^0}^2 - m_{K^+}^2) - (m_{\pi^0}^2 - m_{\pi^+}^2)\right]m_{\pi^0}^2}{(m_K^2 - m_{\pi^0}^2)m_K^2} = (24.2)^{-2},$$
(3)

with $m_K^2 = (m_{K^+}^2 + m_{K^0}^2 - m_{\pi^+}^2 + m_{\pi^0}^2)/2$. Note that $Q_{DT} = Q$ if electromagnetic mass differences

of kaons and pions are equal to each other as Dashen's theorem states [3]. Experimental (Particle Data Group) value $\Gamma = 291 \pm 21$ eV [4] is far from one-loop ChPT value $\bar{\Gamma} = 167 \pm 50 \text{ eV}$ [2] and from the values $\bar{\Gamma} =$ $209 \pm 20 \text{ eV}$ [5] and $\overline{\Gamma} = 219 \pm 22 \text{eV}$ [6] obtained with higher order corrections taken into account by dispersion method. In [5] the subtraction polynomial was taken from the decomposition of one-loop order amplitude and had therefore uncertainties connected to higher orders ChPT corrections. These uncertainties were further fixed [7] by fitting the experimental data [8] on Dalitz-plot distribution in the decay considered: $\bar{\Gamma} = 213^{+3}_{-12}$ eV. There was a conjecture in [7] that new experimental data on Dalitzplot distribution will give slightly different value of $\overline{\Gamma}$. Now these new experimental data (contradicting to the old ones [8]) are available [9]. In what follows we will use them to get new value of $\overline{\Gamma}$ and new value of quark mass ratio Q, as a consequence.

We use the method of work [7] and remind it here for completeness. In order to simulate the experimental DalitzPACS numbers: 13.60.Le

plot distribution we take it in a form

$$1 + ay + by^2 + cy^3 + dx^2 \tag{4}$$

with $a = -1.075 \pm 0.008$, $b = 0.118 \pm 0.009$, $c = 0.13 \pm 0.02$, $d = 0.049 \pm 0.008$ [9] and y, x defined in a standard way

$$y = \frac{3T_0}{Q} - 1,$$

$$x = \frac{\sqrt{3}}{Q}(T_+ - T_-), Q = T_+ + T_- + T_0,$$

 T_+, T_-, T_0 are the kinetic energies of pions in the rest frame $\eta \rightarrow \pi^+ \pi^- \pi^0$ decay. We divide the Dalitz-plot in 10×10 bins $(x \times y)$ that have equal number of events for the distribution considered. Then the number of events in each bin (*n*) is simulated by Gaussian distribution with variance equal to n. We used $n = 10\,000$ to get the full statistics $N = 100n = 1\,000\,000$ like in the experiment [9]. From theoretical point of view the amplitude of $\eta \rightarrow \pi^+ \pi^- \pi^0$ decay have an approximate solution from Eq. (5.28) of [5]. It contains the subtraction polynomial

$$P(s) = \alpha + \beta s_a + \gamma s_a^2 + \delta(s_b - s_c)^2, \qquad (5)$$

where s_a, s_b and s_c are invariant masses squared of $\pi^+ \pi^-$, $\pi^+ \pi^0$ and $\pi^- \pi^0$ pairs, respectively. For the values of parameters α , β , γ and δ within the regions

$$\alpha = -1.28 \pm 0.14, \quad \beta = 21.81 \pm 1.52 \text{ GeV}^{-2}$$

$$\gamma = 4.09 \pm 3.18 \text{ GeV}^{-4}, \quad \delta = 4.19 \pm 1.08 \text{ GeV}^{-4}$$
(6)

(the case of zero subtraction points [5]) the "Minuit" fit of above simulated experimental Dalitz-plot distribution has terminated on the values

$$\alpha_0 = -1.14, \quad \beta_0 = 23.33 \text{ GeV}^{-2}$$

$$\gamma_0 = 1.03q \text{ GeV}^{-4}, \quad \delta_0 = 5.27 \text{ GeV}^{-4}$$
(7)

with $\chi^2/Nd.o.f. = 152/(100 - 4)$.

Three from four parameters are at the boundary of allowed region (6). This probably means that the guess [5] of the size of allowed region should be changed.

Equally possible is the fit with the scaled values of parameters α , β , γ and δ because the normalization factor of the amplitude is not defined by the Dalitz-plot distribution. In our case the scaling of parameters α , β , γ and δ puts them outside the allowed region (6) and no freedom in the scaling (no error in $\overline{\Gamma}$) is possible. This way we get $\overline{\Gamma} = 229\text{eV}$ what corresponds according to Eqs. (1)–(3) to the light quark mass ratio $Q = 22.8 \pm 0.4$. The errors here are due to the errors in the experimental value of the width Γ . So, we conclude, high statistics Dalitz-plot distribution gives the value of light quark mass ratio Q slightly lower

than that from the assumption of equality of kaon and pion electromagnetic mass differences ($Q_{DT} = 24.2$). This is in agreement with calculations of electromagnetic mass differences for pions and kaons [10,11] which find large violations to Dashen's theorem ($Q = 22.0 \pm 0.6$). Our result agrees also very well with that of works [5,6], where the values $Q = 22.4 \pm 0.9$ and $Q = 22.7 \pm 0.8$ were obtained, correspondingly.

This work was partially supported by RFBR grant No. 02-02-16957.

- [1] H. Leutwyler, hep-ph/9609467.
- [2] J. Gasser and H. Leutwyler, Nucl. Phys. B **250**, 539 (1985).
- [3] R. Dashen, Phys. Rev. 183, 1245 (1969).
- [4] S. Eidelman et al., Phys. Lett. B 592, 1 (2004).
- [5] J. Kambor, C. Wiesendanger, and D. Wyler, Nucl. Phys. B 465, 215 (1996).
- [6] A. V. Anisovich and H. Leutwyler, Phys. Lett. B 375, 335 (1996).
- [7] B. V. Martemyanov and V. S. Sopov, Phys. At. Nucl. 67, 424 (2004).
- [8] M. Gormley et al., Phys. Rev. D 2, 501 (1970).
- [9] KLOE collaboration, B. Di Micco et al., hep-ex/0410072.
- [10] J. Donoghue and A. Perez, Phys. Rev. D 55, 7075 (1997).
- [11] J. Bijnens and J. Prades, Nucl. Phys. B 490, 239 (1997).